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1                           **Oblique rifting at oceanic ridges:**  
2                           **Relationship between spreading and stretching directions**  
3                           **from earthquake focal mechanisms**

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5  
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9  
10 **Abstract.** The relationship between spreading and stretching directions is investigated at  
11 oblique-spreading oceanic ridges using earthquake focal mechanisms. The stretching direction  
12 at ridge axes corresponds to the direction of the greatest principal strain  $\epsilon_1$  taken as the mean  
13 trend of the seismic T-axes of extensional earthquake focal mechanisms. It is compared with  
14 the spreading direction provided by global plate-motion models. We find that the stretching  
15 direction trends approximately halfway between the spreading direction and the normal to the  
16 ridge trend, a result in line with analogue experiments of oblique rifting. This result is  
17 satisfactorily accounted for with an analytical model of oblique rifting, for which the direction  
18 of  $\epsilon_1$  is calculated with respect to rifting obliquity for different amounts of stretching using  
19 continuum mechanics. For low stretching factors, typical of incremental seismic  
20 deformations,  $\epsilon_1$  obliquity is two times lower than rifting obliquity. For higher stretching  
21 factors, the stretching and spreading directions become parallel.

22  
23 **Keywords:** oblique rifting, oblique-spreading ridges, stretching, extension

## 24 1. Introduction

25 Determining the direction of relative motion between two rigid plates on either side of a  
26 deformation zone can be achieved by analysing the strain within the deformation zone. In  
27 oblique deformation settings, i.e., when the direction of displacement between the two rigid  
28 plates is oblique to the deformation zone, the direction of relative motion is generally not  
29 parallel to the principal strain directions (e.g., Sanderson and Marchini, 1984; Tikoff and  
30 Teyssier, 1994; Dewey et al., 1998; Fossen and Tikoff, 1998). This result is for example the  
31 case at the axial rifts of oblique-spreading mid-oceanic ridges (Taylor et al., 1994; Tuckwell  
32 et al., 1996), which are investigated in this paper.

33 The process of oblique divergence between two tectonic plates often involves the  
34 formation of an oblique rift. Oblique rifting occurs in the continental domain (e.g. Lake  
35 Baikal; Petit et al, 1996) as well as in the oceanic domain at the axis of slow-spreading ridges  
36 (e.g., Southwest Indian Ridge). The faulting and strain patterns associated with oblique rifting  
37 have been investigated for both oceanic and continental rifts (Dauteuil and Brun, 1993, 1996;  
38 Murton and Parson, 1993; Shaw and Lin, 1993; Taylor et al., 1994; Applegate and Shor,  
39 1994; Carbotte and Mac Donald, 1994; McAllister et al., 1995; Dauteuil et al., 2001; Acocella  
40 and Korme, 2002; Clifton and Schlische, 2003; Fournier et al., 2004a), and by means of  
41 experimental (Withjack and Jamison, 1986, Tron and Brun, 1991; Dauteuil and Brun, 1993;  
42 McClay and White, 1995; Bonini et al., 1997; Clifton et al., 2000; Mart and Dauteuil, 2000;  
43 Clifton and Schlische, 2001; Venkat-Ramani and Tikoff, 2002), analytical (Elliot, 1972;  
44 Sanderson and Marchini, 1984; McCoss, 1986; Withjack and Jamison, 1986; Fossen and  
45 Tikoff, 1993; Tikoff and Fossen, 1993, 1998; Krantz, 1995; Tuckwell et al., 1996; Abelson  
46 and Agnon, 1997), and numerical (Tuckwell et al., 1998) models. These studies show that  
47 oblique rifting is accommodated by both normal and strike-slip faults, whose relative  
48 proportions and orientations depend on rifting obliquity defined as the angle between the  
49 normal to the rift trend and the direction of displacement. Oblique rifting typically produces  
50 en echelon fault patterns that are not perpendicular to the direction of relative motion.

51 Withjack and Jamison (1986) demonstrated, with analogue clay models marked at their  
52 surface by deformed circles, that three structural directions are linked in the process of  
53 oblique rifting: the rift trend (or its perpendicular), the direction of relative motion between  
54 the two plates, and the trend of the greatest principal strain axis  $\epsilon_1$  of the finite strain ellipsoid  
55 (Figure 1). When the direction of relative motion is perpendicular to the rift trend, the rift  
56 formation involves pure shear extension without simple shear and the deformation is  
57 accommodated by dip-slip normal faults parallel to the rift. The  $\epsilon_1$  axis is then horizontal,  
58 perpendicular to the normal faults, and parallel to the direction of divergence. When the  
59 relative motion is oblique to the rift trend, i.e., in transtensional settings, the rift formation  
60 involves a combination of pure shear extension and simple shear. The deformation is  
61 accommodated by a combination of normal faults parallel and oblique to the rift trend, and  
62 also by strike-slip faults when the rifting obliquity increases. In this case,  $\epsilon_1$  is approximately  
63 bisector of the angle between the displacement vector and the normal to the rift (Withjack and  
64 Jamison, 1986). The analytical solution to the problem of oblique rifting, based on the general  
65 theory of transpression-transtension developed by Sanderson and Marchini (1984) and Tikoff  
66 and Teyssier (1994), confirms that the infinitesimal extension direction is exactly the bisector  
67 of the angle between the displacement vector and the normal to the rift (see also McCoss,  
68 1986).

69 Tron and Brun (1991) and Clifton et al. (2000) showed with laboratory experiments that  
70 the fault strike distribution in oblique rifts depended on the rifting obliquity. Consequently, a  
71 statistical analysis of fault strikes in natural rifts may provide an accurate estimate of the  
72 direction of divergence. This rule has been applied successfully to determine the direction of  
73 spreading along two slow-spreading ridges, the Mohns Ridge in the North Atlantic Ocean  
74 (Dauteuil and Brun, 1993) and the West Sheba Ridge in the Gulf of Aden (Dauteuil et al.,  
75 2001), and the kinematic evolution of the Okinawa Trough (Sibuet et al., 1995, Fournier et  
76 al., 2001a). Taylor et al. (1994) and Tuckwell et al. (1996) examined the relationship between  
77 the orientation of extensional fractures and the plate motion vector at oblique spreading ridges  
78 and at so-called “extensional transform zones” (ETZ) characterized by an obliquity between

79 45° and 75° (Taylor et al., 1994). They observed that, at oblique spreading ridges, most  
80 normal faults form at an angle with the ridge axis approximately equal to the half of the plate  
81 motion obliquity, a result in line with the experiments of Withjack and Jamison (1986), Tron  
82 and Brun (1991), and Clifton et al. (2000).

83         However, with the exception of the work of Withjack and Jamison (1986), these studies  
84 mainly focused on fault strikes and did not regard the implications in terms of strain. In  
85 experimental models as well as in the offshore domain, statistical analysis of fault  
86 distributions does not allow estimation of strain axes directions because slip vectors on fault  
87 planes cannot be directly observed. In seismically active rifts, however, the direction of  
88 maximum stretching can be inferred from earthquake focal mechanisms. In the following, we  
89 investigate the relationship between spreading and stretching directions as determined from  
90 earthquake focal mechanisms at six oblique spreading ridges.

91

## 92 **2. Stretching direction determined from earthquake focal mechanisms**

93         In a homogeneous and isotropic material, rupture occurs on two conjugate planes of  
94 maximum shear stress oriented with respect to the maximum and minimum stresses  $\sigma_1$  and  $\sigma_3$ .  
95 Because most earthquakes occur on pre-existing faults, earthquakes do not provide direct  
96 evidence for the orientation of principal stresses, but instead provide evidence for the  
97 orientation of the strain axes (e.g., Twiss and Unruh, 1998). The compression (P) and tension  
98 (T) axes of the double-couple focal mechanism solutions are defined kinematically by fault  
99 slip and correspond to the principal strain axes  $\epsilon_3$  and  $\epsilon_1$ , respectively. They represent the  
100 principal axes of the incremental (or instantaneous) strain tensor for fault movements (e.g.,  
101 McKenzie, 1969; Marrett and Allmendinger, 1990). Thus, in extensional settings, T-axes of  
102 normal faulting earthquakes can be used to determine the direction of stretching. This method  
103 is applicable in regions of homogeneous deformation, i.e. when focal mechanisms are all of  
104 the same type, which is the case at spreading centres of oceanic ridges.

105

## 106 **3. Stretching vs spreading directions at oblique spreading ridges**

107           In the oceanic domain, rifting occurs at the crest of slow-spreading mid-oceanic ridges  
108 characterised by high seismic activity. Fast spreading centres are devoid of an axial rift and  
109 seismicity, and are characterized by orthogonal spreading except in a few back-arc basins  
110 where ETZ have been described, such as the Manus and Lau basins (Taylor et al., 1994). At  
111 fast spreading ridges, the obliquity between the spreading direction and the plate boundary is  
112 taken up by transform faults (e.g., Pacific-Antarctic Ridge). The main oblique-spreading  
113 ridges on Earth are the Southwest Indian Ridge (SWIR) in the Indian Ocean (Figure 2; Ewing  
114 and Heezen, 1960; Fisher and Sclater, 1983; Patriat, 1987), the Sheba Ridge in the Gulf of  
115 Aden (Figure 3; Matthews et al., 1967; Laughton et al., 1970), and the Reykjanes (Figure 4;  
116 Vine, 1966), Mohs (Figure 5; Talwani and Eldholm, 1977), and Knipovich (Figure 5; Vogt  
117 et al., 1979; Okino et al., 2002) ridges in the North Atlantic Ocean. These five ridges have  
118 been surveyed together with the Carlsberg Ridge in the northwest Indian Ocean (Figure 3;  
119 Schmidt, 1932; Vine and Matthews, 1963), which is generally considered as a type example  
120 of orthogonal-spreading ridge.

121           We have selected in the Harvard centroid moment tensor (CMT) catalog all focal  
122 mechanisms of earthquakes shallower than 50 km which occurred between 1976 and 2000 (25  
123 years) along these six ridges (Dziewonski et al., 1981). 271 mechanisms of extensional or  
124 strike-slip type have been obtained and are plotted in Figures 2 to 5. For each ridge or ridge  
125 segment, we determined its mean trend, the mean spreading direction, and the mean stretching  
126 direction (Table 1). If necessary, the ridges have been divided in roughly rectilinear segments.  
127 For example, the SWIR has been divided into two parts: the northeastern part strikes N54°E  
128 on average and the southwestern part N105°E (Figure 2). The ridge mean trend has been  
129 directly measured on bathymetric and seismic maps. The mean spreading direction  
130 corresponds to the average of the spreading directions calculated at the ridge segment  
131 extremities from the NUVEL-1A plate motion model (DeMets et al., 1990; 1994), except for  
132 Sheba and Carlsberg ridges for which we used Fournier et al. (2001b) solution (Table 1). The  
133 mean stretching direction is computed from the normal faulting solutions (inserts in Figures 2  
134 to 5). From these data, the spreading and stretching obliquities have been calculated for each

135 ridge (Table 1). Strike-slip focal mechanisms along transform faults are also plotted in  
136 Figures 2 to 5 to show the consistency between slip vectors of strike-slip mechanisms and  
137 spreading directions provided by plate motion models.

138 The stretching obliquity ( $S_{obl}$ ) is plotted against spreading (or rifting) obliquity ( $R_{obl}$ ) for  
139 the selected ridges in Figure 6. Spreading obliquities greater than  $45^\circ$  are never observed  
140 along slow-spreading ridges. The points plot along the  $S_{obl} = R_{obl} / 2$  line for spreading  
141 obliquities less than  $30^\circ$  (Carlsberg, southwestern SWIR, Reykjanes, and Knipovich ridges),  
142 and slightly depart from this line for obliquities between  $30^\circ$  and  $45^\circ$  (Mohns, northeastern  
143 SWIR, and Sheba ridges).

144

#### 145 **4. Analytical model of oblique rifting**

146 A horizontal plane-strain model of oblique rifting is presented in Figure 6a. A unit  
147 length of lithosphere (initial rift) is obliquely extended to a length  $\beta$  measured perpendicularly  
148 to the rift axis.  $\beta$  thus defines a stretching factor corresponding to the ratio of the final versus  
149 initial length (e.g., McKenzie, 1978). The stretching obliquity, defined as the angle between  
150 the normal to the rift trend and greatest principal strain axis of the strain ellipse ( $\epsilon_1$ ), is  
151 calculated as a function of the rifting obliquity and  $\beta$ .

152 The finite strain ellipse is calculated from continuum mechanics by decomposing the  
153 deformation matrix (deformation gradient tensor) in finite strain (shape and orientation of the  
154 strain ellipse in 2D) and finite rotation of the principal strain axes (e.g., Elliot, 1972; Jaeger  
155 and Cook, 1979; McKenzie and Jackson, 1983; Fournier et al., 2004b). The eigenvalues and  
156 eigenvectors of the finite strain matrix provide the length and orientation of the principal axes  
157 of the finite strain ellipse. Exactly the same result is obtained by factorization of the  
158 deformation matrix into pure shear and simple shear components (e.g., Sanderson and  
159 Marchini, 1986; Tikoff and Fossen, 1993; Fossen and Tikoff, 1993; Tikoff and Teyssier,  
160 1994; Krantz, 1995; Fossen and Tikoff, 1998).

161 The strain ellipse resulting from oblique rifting is shown as a function of the rifting  
162 obliquity for various values of stretching factor  $\beta$  in Figure 6b. For a given rifting obliquity,

163 the principal strain axes progressively rotate as  $\beta$  increases. For a rifting obliquity of  $45^\circ$ , the  
 164 stretching obliquity increases from  $24^\circ$  for  $\beta = 1.1$  to  $36^\circ$  for  $\beta = 3$ . Furthermore, for a  
 165 given  $\beta$ , the stretching obliquity increases as the rifting obliquity increases. For example, for  
 166  $\beta = 2$ , the stretching obliquity increases to  $10^\circ$  to  $20^\circ$ ,  $32^\circ$ ,  $45^\circ$ , and  $63^\circ$  for rifting obliquity  
 167 of  $15^\circ$  to  $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ , and  $75^\circ$ , respectively.

168 In Figure 6c, the stretching obliquity is plotted against rifting obliquity for various  
 169 values of  $\beta$ . When  $\beta$  is small ( $\beta < 1.1$ ), the stretching obliquity is equal to the half of the  
 170 rifting obliquity ( $S_{obl} = R_{obl} / 2$ ). With increasing strain ( $\beta > 5$ ), the stretching obliquity  
 171 becomes almost equal to the rifting obliquity ( $S_{obl} = R_{obl}$ ).

172

## 173 5. Discussion

174 These predictions can be compared with the results obtained for the selected oblique-  
 175 spreading ridges (Figure 6). For most ridges, the  $\epsilon_1$  direction ranges along the  $S_{obl} = R_{obl} / 2$   
 176 line, which corresponds to a low amount of extension in the model. A simple interpretation is  
 177 that rocks of the Earth's upper crust undergo small strains of a few per cent before brittle  
 178 failure occurs and relieves the accumulated strain. The principal strain directions deduced  
 179 from earthquake focal mechanisms thus represent the infinitesimal (or instantaneous) strain  
 180 ellipsoid.

181 Our results can also be compared with those of Taylor et al. (1994) and Tuckwell et al.  
 182 (1996) for the Reykjanes, Mohns, Southwest Indian (NE), and Sheba ridges, provided one  
 183 converts their  $\alpha$  and  $\phi$  angles into rifting and stretching obliquities:

$$184 \quad \begin{aligned} R_{obl} &= 90 - \alpha \\ S_{obl} &= \phi - \alpha \end{aligned}$$

185 In contrast with us, Taylor et al. (1994) and Tuckwell et al. (1996) defined the stretching  
 186 direction as the perpendicular to the mean trend of normal faults in extension zones.

187 We find a very good agreement for the Reykjanes Ridge, where our estimates of  
 188 spreading and stretching obliquities differ only by  $1^\circ$  and  $4^\circ$ , respectively, which is smaller  
 189 than the uncertainties. For the Mohns Ridge, our results compare well with those of Taylor et



190 al. (1994) but slightly differ from Tuckwell et al. (1996) estimates of spreading obliquity  
 191 ( $34\pm 8^\circ$  vs  $40\pm 6^\circ$ ), mainly because we (and Taylor et al., 1994) use a different azimuth of  
 192 spreading (N119°E vs N110°E). Despite this, we find no large discrepancies between our  
 193 estimates of stretching obliquity and theirs. Much larger differences are found for the  
 194 Southwest Indian and Sheba ridges: for the former, whereas Taylor et al. (1994) and Tuckwell  
 195 et al. (1996) give comparable values of  $14^\circ$  and  $23\text{--}27^\circ$  for stretching and spreading  
 196 obliquities, we find  $29\pm 7^\circ$  and  $42\pm 11^\circ$ , respectively. These differences are entirely  
 197 attributable to different estimates of ridge trend and spreading directions, due to the fact that  
 198 Taylor et al. (1994) and Tuckwell et al. (1996) took into account only a small part of the  
 199 SWIR located near the Rodrigues triple junction ( $26^\circ\text{N}$ ; Mitchell, 1991), whereas we have  
 200 taken into account all the northeastern part of the SWIR over several thousands kilometres  
 201 (Figure 2). However, here again, the determination of stretching directions from earthquake  
 202 focal mechanisms gives results comparable to the analysis of normal fault trends. Concerning  
 203 the Gulf of Aden (Sheba Ridge), a difference up to  $5\text{--}10^\circ$  exists between our values and those  
 204 of Taylor et al. (1994) and Tuckwell et al. (1996). Once again, these differences come from  
 205 the selection of different study areas. The results of Taylor et al. (1994) and Tuckwell et al.  
 206 (1996) concern the westernmost part of the Sheba Ridge near the Gulf of Tadjura ( $45^\circ\text{E}$ ;  
 207 Tamsett and Searle, 1988), whereas our results encompass the entire ridge from  $46^\circ\text{E}$  to  $56^\circ\text{E}$   
 208 (Figure 3; Table 1). Hence, the differences between our results and those of Taylor et al.  
 209 (1994) and Tuckwell et al. (1996) come from different scales of study. Studying normal fault  
 210 strikes at ridge axes requires detailed mapping of fault fabrics. Working with focal  
 211 mechanisms from the world seismicity catalogs allows surveying of larger areas.

212 In general, the results of Tuckwell et al. (1996) show that most values of stretching vs  
 213 rifting obliquities range along the  $S_{obl} = R_{obl} / 2$  line, like in the present study. Surprisingly, the  
 214 direction of  $\epsilon_1$  deduced from infinitesimal strain (earthquakes) does not differ from the  
 215 perpendicular to the normal faults, which are markers of finite strain and can have  
 216 accommodated a significant amount of deformation. This result suggests that normal faults  
 217 initially form perpendicular to the direction of  $\epsilon_1$  of the infinitesimal strain ellipsoid, keep this

218 orientation during ongoing extension, and do not significantly rotate as the strain increases.  
219 As oblique slip (characterized by oblique focal mechanisms) is seldom observed, this implies  
220 that normal faults at ridge axes only accommodate a small amount of deformation during the  
221 time when they are located in the seismically active part of the rift (about 2 Ma for a ridge  
222 with a half-spreading rate of 5 mm/yr and a 20 km large axial rift).

223

## 224 **6. Conclusion**

225 Plate-motion models such as RM2 and NUVEL-1 (Minster and Jordan, 1978; DeMets et  
226 al., 1990) did not account for slip vectors of extensional focal mechanisms along oceanic  
227 ridges. The first reason was of course that, for extensional mechanisms, it is not possible to  
228 determine which of the two nodal planes is the fault plane and which is the actual slip vector.  
229 The second reason was that at oblique-spreading ridges, slip vectors are not parallel but  
230 oblique to the plate relative motion. Here, we demonstrate that, at slow-spreading oblique  
231 ridges, the maximum strain axis determined from earthquake focal mechanisms trends  
232 halfway between the direction of spreading and the normal to the ridge. Hence, the kinematics  
233 of oblique ridges and rifts can possibly be determined from a set of extensional focal  
234 mechanisms, without assumption on the fault plane and the slip vector. This result could be  
235 useful in continental rifts where transform faults are not developed and plate kinematics  
236 difficult to assess. The comparison with an analytical model of oblique rifting shows that  
237 these features correspond to small deformations at ridge axes, which is consistent with the  
238 fact that earthquakes represent infinitesimal strains. Furthermore, the analysis of normal faults  
239 directions (Taylor et al., 1994; Tuckwell et al., 1996) yields similar conclusions, though  
240 normal fault heaves represent thousands of co-seismic slips. Yet, compared to the rift width  
241 (~10 to 20 km), the cumulated stretching factor on each fault must remain low.

242

243

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402 **Figure captions**

403 Figure 1. Geometrical relationship between the main structural directions at oblique rifts.

404

405 Figure 2. Bathymetric map (Sandwell and Smith, 1997), shallow seismicity between 1964  
406 and 1995 (focal depth < 50 km; magnitude > 2; Engdahl et al., 1998), and all available  
407 earthquake focal mechanisms (Harvard CMT for the period 1976-2000; Dziewonski et al.,  
408 1981) for the Southwest Indian Ridge (SWIR). Inserted stereoplots are equal-area projections  
409 of the P and T axes of the extensional focal mechanisms and the mean direction of extension  
410 ( $\epsilon_1$ ). The SWIR has been divided into two parts with different trends: the northeastern part  
411 between the Rodrigues triple junction and the Prince Edward-Marion-Andrew Bain fracture  
412 zone (PEMABFZ; Grindlay et al., 1998) trends N054°E  $\pm$ 2°, and the southwestern part  
413 between PEMABFZ and 53°S, 14°E trends N105°E  $\pm$ 2°. Bathymetric contour interval is  
414 1000m. Strike-slip focal mechanisms along fracture zones show the consistency between slip  
415 vector azimuths and directions of relative motion (solid arrows) calculated from plate motion  
416 models.

417

418 Figure 3. Same legend as Figure 2 for the Sheba and Carlsberg ridges. OTF is Owen  
419 transform faults. Bathymetric contour interval is 500m.

420

421 Figure 4. Same legend as Figure 2 for the Reykjanes Ridge. Between 55.5°N, 35.5°W and  
422 63.5°N, 24°W, the ridge strikes N037°E  $\pm$  3°. Bathymetric contour interval is 500m.

423

424 Figure 5. Same legend as Figure 2 for the Mohns and Knipovich ridges (location in Figure 4).  
425 The Mohns Ridge strikes N063°E  $\pm$  2° on average between 71°N, 7.5°W and 73.5°N, 8°E  
426 The mean trend of the Knipovich Ridge between 73.7°N, 9°E and 78°N, 8°E is N178°E  $\pm$  2°.  
427 Bathymetric contour interval is 200m.

428

429 Figure 6. Stretching obliquity  $S_{obl}$  (maximum principal strain  $\epsilon_1$ ) as a function of spreading  
430 obliquity  $R_{obl}$  in degrees for seven oblique-spreading ridges. Data sources and abbreviations  
431 are given in Table 1. See text for additional explanation. Error bars for spreading obliquity  
432 represent the sum of the uncertainties in the measurement of the ridge mean trend and in the  
433 azimuth of spreading calculated along the ridge. Error bars for stretching obliquity represent  
434 the standard deviation of the T-axes azimuth.

435

436 Figure 7. a. Plane-strain analytical model of oblique rifting. See text for additional  
437 explanation. b. Strain ellipse for various stretching factors and rifting obliquities. c.  
438 Stretching obliquity  $S_{obl}$  as a function of rifting obliquity  $R_{obl}$  in degrees. The curves are  
439 calculated from the analytical model and the straight lines correspond to  $S_{obl} = R_{obl}$  and  
440  $S_{obl} = R_{obl} / 2$ .

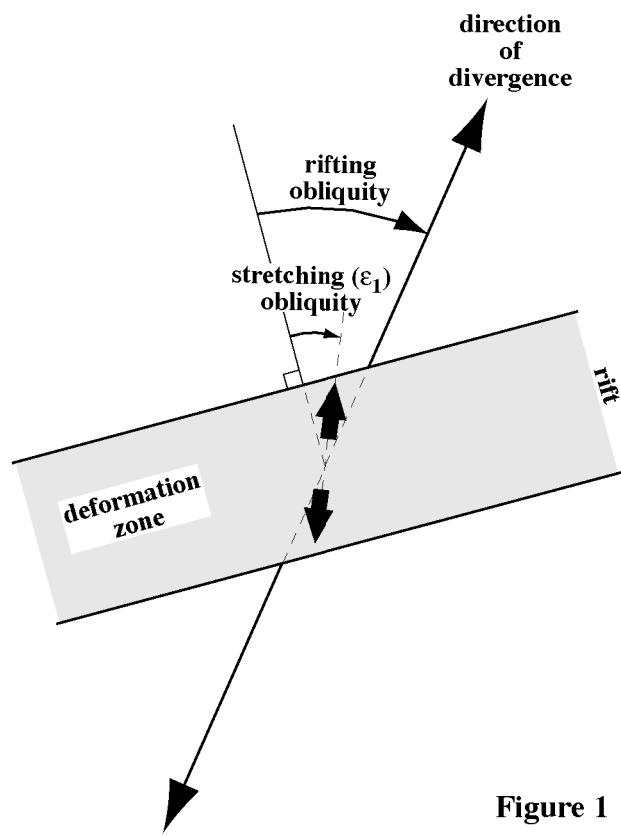


Figure 1

Figure

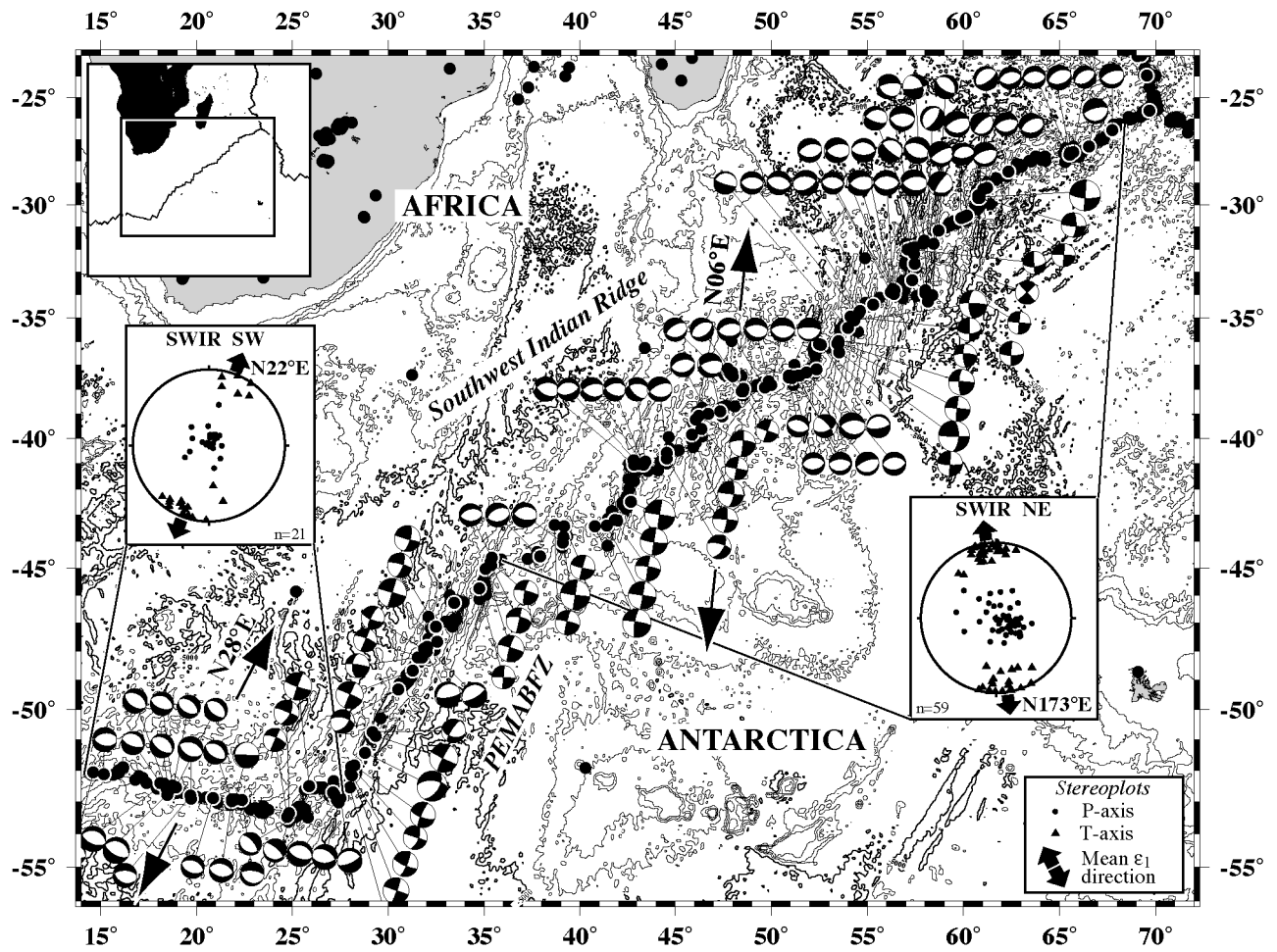


Figure 2

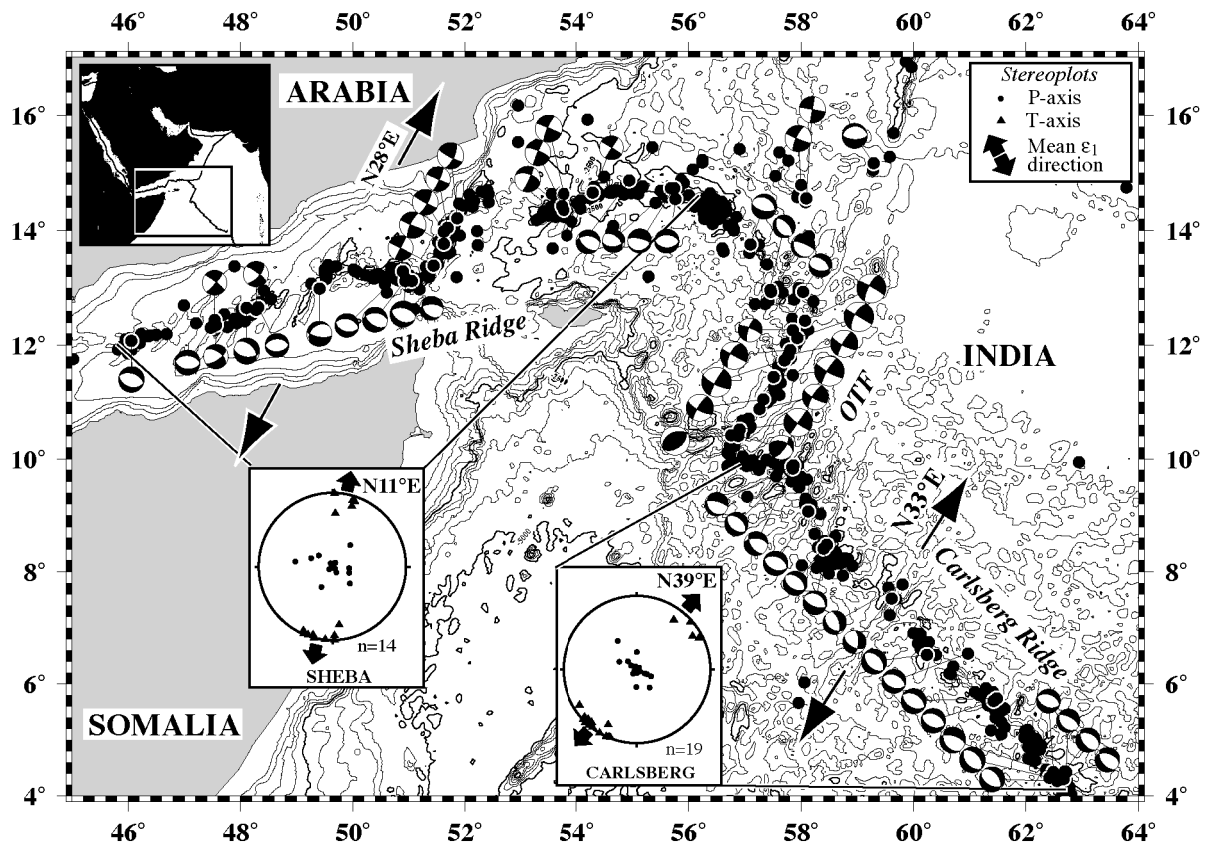


Figure 3

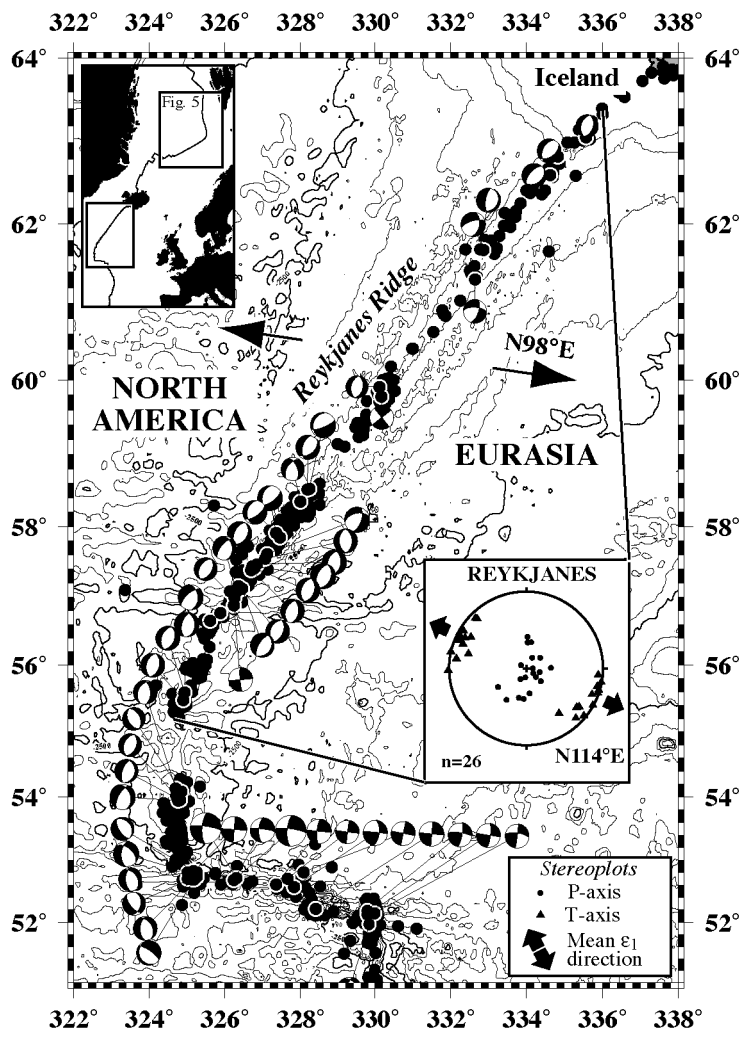


Figure 4

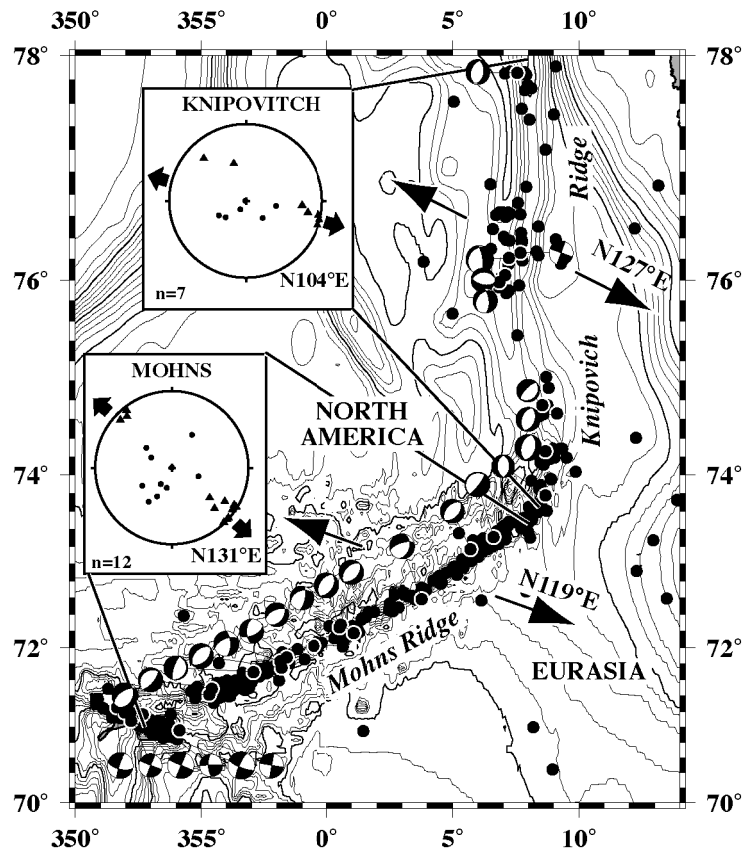


Figure 5



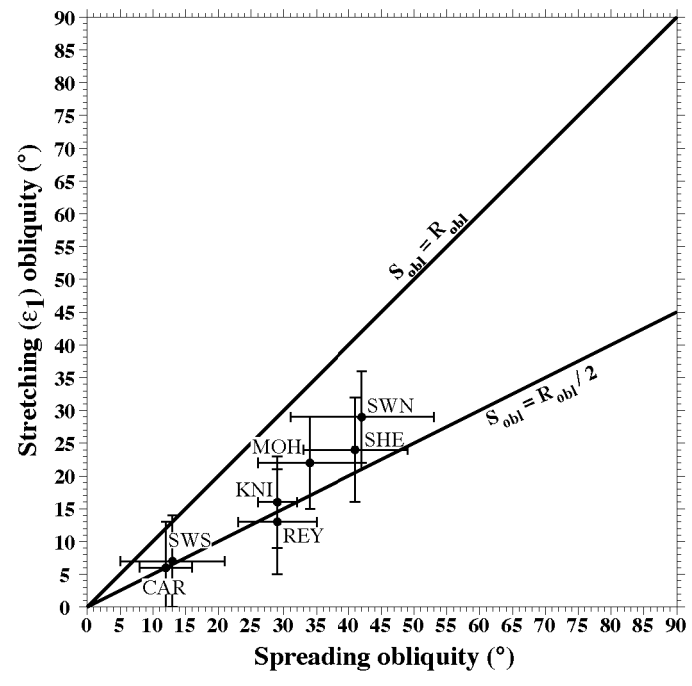


Figure 6

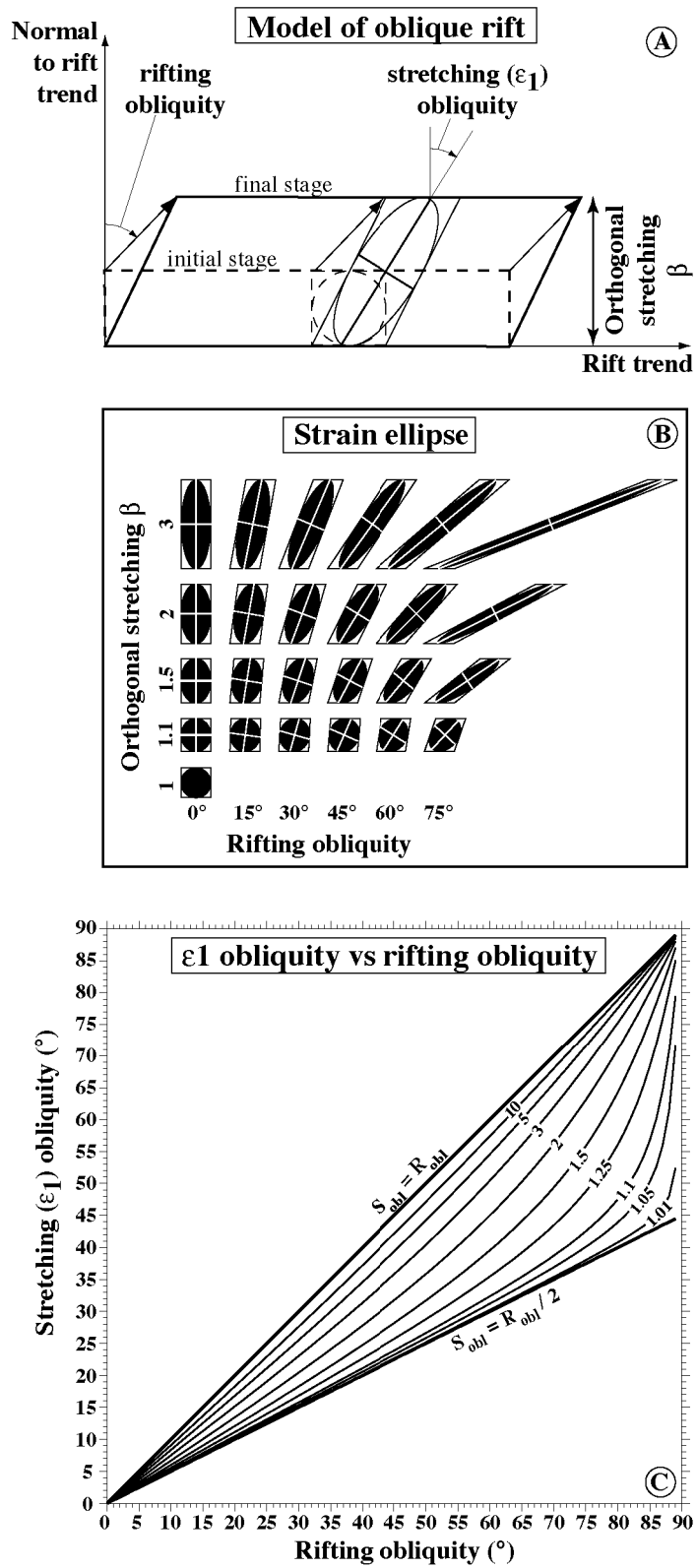


Figure 7

Table 1. Mean trend, azimuth of spreading, spreading obliquity, and principal strain  $\epsilon_1$  obliquity for oblique spreading ridges

Ridge	Ridge mean trend	Ridge extremities			Mean azimuth of spreading	Mean T-axis Strike, Dip	Spreading obliquity	Principal strain $\epsilon_1$ obliquity	Labels <sup>2</sup>
		Latitude	Longitude	Azimuth of spreading <sup>1</sup>					
		(°E)	(°N)	(°E)					
Aden - Sheba	N077°E $\pm 3^\circ$	12	46	33	028 $\pm 5$	011, 1 (n=14)	41 $\pm 8$	24 $\pm 8$	SHE
		14.5	56	23					
Carlsberg	N135°E $\pm 2^\circ$	10	57	31	033 $\pm 2$	219, 0 (n=18)	12 $\pm 4$	6 $\pm 7$	CAR
		4	63	35					
SWIR NE	N054°E $\pm 2^\circ$	-45	35	15	006 $\pm 9$	353, 3 (n=59)	42 $\pm 11$	29 $\pm 7$	SWN
		-26	69	177					
SWIR SW	N105°E $\pm 2^\circ$	-52	14	34	028 $\pm 6$	202, 4 (n=21)	13 $\pm 8$	7 $\pm 7$	SWS
		-53	28	22					
Reykjanes	N037°E $\pm 3^\circ$	55.5	-35.5	95	098 $\pm 3$	294, 2 (n=26)	29 $\pm 6$	13 $\pm 8$	REY
		63.5	-24	102					
Mohn	N063°E $\pm 2^\circ$	71	-7.5	113	119 $\pm 6$	131, 4 (n=12)	34 $\pm 8$	22 $\pm 7$	MOH
		73.5	8	125					
Knipovitch	N178°E $\pm 2^\circ$	73.7	9	126	127 $\pm 1$	104, 4 (n=7)	29 $\pm 3$	16 $\pm 7$	KNI
		78	8	128					

n is the number of extensional earthquake focal mechanisms used to determined the mean T-axes azimuth.

<sup>1</sup>Azimuths of spreading after DeMets et al. (1990), except for SHE and CAR after Fournier et al. (2001).

<sup>2</sup>Labels are for data plotted in Figure 5.