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Line intensities of $^{12}\text{C}_2\text{H}_2$ in the 7.7 μm spectral region

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Abstract

Absolute intensities of 414 lines are measured in eight bands of the 7.7 μm spectral region of the $^{12}\text{C}_2\text{H}_2$ molecule, with an average accuracy of 5%. Vibrational transition dipole moment squared values and empirical Herman-Wallis coefficients are obtained in order to model the rotational dependence of the transition dipole moment squared, except for some forbidden bands for which smoothed values are given. These data are used to calculate a line list for atmospheric or astrophysical applications.

Keywords: Acetylene; Infrared; Vibro-rotational transitions; Line intensities; Databases

1. Introduction

The 7.7 μm spectral region of acetylene $^{12}\text{C}_2\text{H}_2$ corresponds to the $\Delta P = 2$ sequence of vibrational transitions [1], P being a pseudo-quantum number equals to $5\nu_1 + 3\nu_2 + 5\nu_3 + \nu_4 + \nu_5$, where $\nu_1, \nu_2, \nu_3, \nu_4$, and ν_5 are the quantum numbers associated with the normal modes of vibration of the molecule in the ground electronic state. The more extensive work on this spectral region is this of Kabbadj et al [2] who assigned vibro-rotational lines in 16 bands around the $(\nu_4 + \nu_5)_+^0$ strong cold band at 1328 cm^{-1} . These 16 bands are hot bands, except for $(\nu_4 + \nu_5)_+^2$ at 1343 cm^{-1} . Among them, four “forbidden” bands were observed by Kabbadj et al, i.e., two $\Delta \leftarrow \Sigma$ and one $\Sigma \leftarrow \Delta$ bands, as also one $\Phi \leftarrow \Pi$ band. Recently, Robert et al [3] performed a global vibration-rotation analysis on states up to 3000 cm^{-1} , thus including the 7.7 μm spectral region.

The integrated intensity of the band at 1328 cm^{-1} was first measured by Kelly et al [4], and later more accurately by Varanasi and Bangaru [5], and by Koops et al [6]. Theoretical calculations of this band intensity was performed by Abbouti Temsamani et al [7]. Intensities of a few lines of the R -branch of $(\nu_4 + \nu_5)_+^0$ were obtained by Podolske et al [8]. A few years ago, an accurate set of absolute line intensities was obtained by Vander Auwera [9] for the cold bands $(\nu_4 + \nu_5)_+^0$ and $(\nu_4 + \nu_5)_+^2$, the accuracy being estimated better than 2%. Intensities of some lines of $(\nu_4 + \nu_5)_+^0$ were measured again by Jacquemart et al [10] and Lepère et al [11], confirming the announced accuracy. Furthermore, collisional widths of $(\nu_4 + \nu_5)_+^0$ were studied by Varanasi et al [12] (N_2 -broadening at 153 and 200 K), then by Podolske et al [8] (N_2 - and He-broadening at room temperature) and by Devi et al [13] (N_2 - and air-broadening at room temperature), whereas self-broadening coefficients were obtained recently by Lepère et al [11], as well as some N_2 -broadening coefficients by Fissiaux et al [14] and Dyne et al [15].

The aim of this work is to initiate absolute line intensity measurements in bands for which no result has been reported in the literature. Line intensities are known for the $(\nu_4 + \nu_5)_+^0$ and $(\nu_4 + \nu_5)_+^2$ bands [9], but they are unknown for the numerous other bands observed and assigned by Kabbadj et al [2]. These bands are gathered in Table 1. In this table, a given value of P is assigned to a given set of interacting vibrational states, named polyad or cluster. Then, polyads are noted $\{P\nu_5\}$. Vibrational levels are noted $\nu_1 \nu_2 \nu_3 (\nu_4 \nu_5)_{\pm}^{\ell} r$, with $\ell = |\ell_4 + \ell_5|$, ℓ_t being the vibrational angular momentum quantum number associated with the degenerate bending mode t , \pm being the symmetry type for Σ vibrational states ($\ell = 0$), and r a roman numeral indicating the rank of the level, by decreasing energy value ($r = \text{I}$ for the highest energy level), inside the set of states having the same vibrational symmetry, and coupled by ℓ -type resonances.

Intensity measurements were undertaken because the knowledge of C_2H_2 line intensities in the 7.7 μm spectral region is important for several applications, especially for astrophysical ones. For example, the acetylene molecule has been observed in the circumstellar envelopes of carbon-rich stars. Using the Infrared Spectrograph (IRS) on board the Spitzer Space Telescope (SST), Matsuura et al [16] detected acetylene bands at 7 and 14 μm in carbon-rich asymptotic giant branch stars in the Large Magellanic Cloud. Around 7 μm , spectroscopic databases as HITRAN [17] and GEISA [18] only contain line positions and intensities that Vander Auwera calculated from his absolute measurements in the $(\nu_4 + \nu_5)_+^0$ band [9], for the rotational quantum number J up to 35. But intensities measured in Ref. [9] for some lines of the $(\nu_4 + \nu_5)_+^2$ band are not reported in the databases. The temperature of interest for applications being around

500 K [16], the knowledge of intensities in the remaining hot bands is also important. In the quoted paper [16], Matsuura et al could not reproduce the shapes that they observed in their IRS-SST spectra around 7 μm because of the lack of data in databases.

The acetylene spectrum being very crowded around 1300 cm^{-1} , we first focused our work on 8 bands, that include the strongest allowed bands and the forbidden ones for their interest. To perform intensity measurements, several spectra have been recorded with a Fourier transform spectrometer. A set of 414 absolute line intensities have been obtained with an average accuracy of 5%. Section 2 of the paper will be devoted to the experimental conditions and to the measurements. The results and their analysis will be given in Section 3. The last section explains how we have proceeded to set up a line list for databases.

2. Experimental details and measurements

Six spectra have been recorded with the rapid scan Bruker IFS 120 HR interferometer of the LADIR (LABoratoire de Dynamique, Interactions et Réactivité, Paris). Experimental conditions are gathered in Table 2. This interferometer was equipped with a Globar source, Ge/KBr beam splitter, and a photovoltaic MCT (Mercury-Cadmium-Telluride) detector. An optical filter covering the spectral region between 800 to 2000 cm^{-1} has been used. The whole optical path was under vacuum and a multipass cell of 1-m base length was used. The cell was equipped with KCl windows. The temperature of the gas in the cell was recorded via four platinum probes at different places inside the cell. The uncertainty on the temperature measurements has been estimated to be ± 0.5 K. Pressures were measured using two full scale ranges MKS Baratron (10- and 100-Torr manometers) with an accuracy of 0.5%. Each scan among about 400 recorded for every spectrum has then been individually transformed to a spectrum using the Fourier transform procedure included in the Bruker software OPUS package [19], selecting a Mertz phase error correction [20,21].

To deduce line intensities from the spectra, a multispectrum procedure was used following a method already described [22]. In this method, experimental data are concatenated in a single vector, and the adjusted parameters are constrained to be the same for all the spectra simultaneously adjusted. Then, line parameters are obtained in a single run of the procedure and their accuracy is improved. Before to run the procedure, one has first to adjust an effective value of the iris radius for each spectrum, to check that there is no residual due to a phase error, and to calibrate the absolute wavenumber scale of each spectrum. For that, the wavenumbers given by Vander Auwera around 7.7 μm [9] have been used. The accuracy of measured line positions of not too weak and well isolated lines is estimated to be 0.0002 cm^{-1} on the mean. Then, the multispectrum procedure was run in the following conditions. A Voigt profile was used to calculate the absorption coefficient of the lines. The Doppler-width is kept fixed at its theoretical value and the baseline is adjusted as a polynomial of the second degree around each studied line. Because of the relatively low pressures, the self-broadening coefficients were fixed at the values calculated according to Ref. [23], and the self-shifting coefficients were fixed at zero. On the whole, 414 line intensities have been measured in 8 bands with an average accuracy estimated around 5%, but that can be degraded up to 10% for very weak lines or overlapped ones. The results obtained for selected bands are given in Tables 3 and 4. Results obtained for the other bands can be found in the supplementary material attached to this paper.

3. Data reduction and results

For each line intensity $S(T_0)$ obtained from the multispectrum fitting procedure, in cm molecule^{-1} for pure $^{12}\text{C}_2\text{H}_2$ (i.e., for a sample containing 100% of $^{12}\text{C}_2\text{H}_2$) at the standard temperature $T_0 = 296$ K, we used the following formula to deduce the transition dipole moment squared $|R|^2$, in D^2 (1 debye = 3.33546×10^{-30} C m)

$$S(T_0) = (1/4\pi\epsilon_0) (8\pi^3/3hc) [g''\nu_0 / g_V Q(T_0)] |R|^2 L(J,\ell) \exp(-hcE''/kT_0) [1-\exp(-hc\nu_0/kT_0)], \quad (1)$$

where $1/4\pi\epsilon_0 = 10^{-36} \text{ erg cm}^3 \text{ D}^{-2}$; h is Planck's constant equal to $6.6260755 \times 10^{-27} \text{ erg s}$ (1 erg = 10^{-7} J); c is the speed of light in vacuum equal to $2.99792458 \times 10^{10} \text{ cm s}^{-1}$; g'' is the statistical weight due to nuclear spin of the lower level (1 for s -type levels and 3 for a -type levels); ν_0 is the transition wavenumber in cm^{-1} ; g_V depends on the degeneracy of the levels involved, with the convention g_V equals 2 if both upper and lower vibrational states are degenerate and equals 1 otherwise; $Q(T_0)$ is the total partition function at temperature T_0 , calculated from Fischer et al [24]; $L(J,\ell)$ is the Hönl-London factor; E'' , in cm^{-1} , is the energy of the lower level taken from HITRAN [17]; k is Boltzmann's constant equal to $1.380658 \times 10^{-16} \text{ erg K}^{-1}$.

Among the 8 studied bands, 4 are parallel ones (3 of them being ℓ -type doubled) and 4 are forbidden ones. Only one Q -branch could be observed. For P - and R -branches of parallel bands ($\Delta\ell = 0$) and of forbidden ones ($\Delta\ell = \pm 2$), the Hönl-London factor is taken as in Ref. [9]

$$L(J,\ell) = (J+1+\ell) (J+1-\ell) / (J+1) \text{ (for } R\text{-branch)}, \quad (2)$$

$$L(J,\ell) = (J+\ell) (J-\ell) / J \text{ (for } P\text{-branch)}, \quad (3)$$

whereas for the Q -branch of the forbidden band with $\Delta\ell = +2$ and $\ell = 0$, it is taken as

$$L(J,\ell) = (2J+1) / 2. \quad (4)$$

To reduce the data, effective parameters can be deduced expanding $|R|^2$ to take into account the rotational dependence and resonances. For all the bands studied in the present work, the following empirical Herman-Wallis-type factors have been tried

$$|R|^2 = |R_0|^2 (1 + A_1^{PR} m + A_2^{PR} m^2)^2 F' \text{ (for } P\text{- and } R\text{-branches)}, \quad (5)$$

$$|R|^2 = |R_0|^2 (1 + A_2^Q m^2)^2 F' \text{ (for the } Q\text{-branch)}, \quad (6)$$

m being equal to $-J$ in the P -branch, $J+1$ in the R -branch, and J in the Q -branch. $|R_0|^2$ is the vibrational transition dipole moment squared, and A_1^{PR} , A_2^{PR} , and A_2^Q , are Herman-Wallis coefficients. Note that the terms between parentheses in Eqs. (5,6) are squared in the present work. Moreover,

$$F' = 1 \text{ for } \Delta\ell = 0. \quad (7)$$

For $\Delta\ell = \pm 2$, F' is taken as

$$F' = [m(m+1)]^2. \quad (8)$$

The case of the two strong bands $(\nu_4 + \nu_5)^0_+$ and $(\nu_4 + \nu_5)^2$, for which line intensities were already measured by Vander Auwera [9], deserves a discussion. For $(\nu_4 + \nu_5)^0_+$, line assignments and measurements of positions were performed by Kabbadj et al [2] up to $P43$ and $R41$, and line intensities were measured in Ref. [9] up to $P35$ and $R35$. As we needed relatively large values of the product “pressure \times absorbing path length” in order to study weak bands, most of the lines measured by Vander Auwera are saturated in our spectra, so that we could measure only 10 lines that overlap the set of lines studied in Ref. [9] and reported in HITRAN [17]. However, the advantage of spectra with large pressure \times absorbing path length product is that we could enlarge the set of intensity data up to $P43$ and $R45$, increasing J by about 10 in each branch (see Table 3). Then, we have calculated $|R|^2$ values from the experimental intensities reported in Ref. [9], averaging them when, for a line, several values coming from different spectra had been quoted, and converting them to pure $^{12}\text{C}_2\text{H}_2$ in cm molecule^{-1} at 296 K. (The obtained values are equal to those given in Ref. [9] \pm a few 0.1%, because of roundings and averaging of line intensities, and because of a slight difference between the total partition functions used in Ref. [9] and in the present work. Note that in Table 2 of Ref. [9], the listed $|R|^2$ values are given for natural abundances.) $|R|^2$ values such obtained are plotted on Fig. 1 together with ours. Experimental values deduced from the line intensities of Lepère et al [11] have also been plotted (we have chosen those obtained from a Rautian profile in Ref. [11]). Though the presence of resonances calls for a rigorous theoretical model such these of Refs. [9,25,26], Fig. 1 shows that an effective treatment through an Herman-Wallis factor is powerful and allows to calculate the spectrum with enough precision for most applications (RMS of the fit smaller than 2.7%). To get improved values of $|R_0|^2$ and of Herman-Wallis coefficients, $|R|^2$ values from this work (26 values) and from Ref. [9] (40 values) have been fitted together. The obtained constants are listed in Table 5. A fit merging data from this work and from Refs. [9] and [11] has also been performed. It gave very similar results but degraded slightly the quality of the residuals. Table 5 shows that the retrieved $|R_0|^2$ value is very close to this given in Ref. [9], however keeping in mind the fact that these two values are not strictly comparable since coming from different theoretical models.

As far as the $(\nu_4 + \nu_5)^2$ band is concerned, line intensities could be measured for 83 lines (see Table 4), compared with 24 in Ref. [9]. Particularly, P - and R -branches are now known from $P4$ to $P45$ and from $R6$ to $R41$. Let us recall that in Ref. [10] a good agreement was found between intensities of 28 lines of $(\nu_4 + \nu_5)^0_+$ measured from Fourier transform spectra recorded at GSMA (Groupe de Spectrométrie Moléculaire et Atmosphérique, Reims) and in Ref. [9], namely $-1.5(30)\%$, with 1 standard deviation (noted 1SD in the following) between parenthesis in unit of the last digit. Such a comparison is not so easy in the present work since the orders of magnitude of line intensities to compare are very different in the two works. For example, intensities of Q -lines of $(\nu_4 + \nu_5)^2$, for which large systematic negative discrepancies up to about -40% are observed with Ref. [9] (see Fig. 2), cannot lead to a significant comparison since it involves lines having intensities smaller than about $5 \times 10^{-23} \text{ cm molecule}^{-1}$, and then weak and difficult to measure in Ref. [9]. If one compares common results between this work and Ref. [9], disregarding Q -lines, i.e., considering only 22 P - and R -lines of the $(\nu_4 + \nu_5)^0_+$ and $(\nu_4 + \nu_5)^2$

bands, one finds a mean difference of $-1.9(30)\%$ between our results and those of Ref. [9]. If one compares the results of Refs. [11] and [9], for 14 common lines, one finds $-1.6(30)\%$, that is very close to the above differences. These are very acceptable agreements, with a small dispersion of the differences (1 SD 3%), if one takes into account the accuracies announced in the concerned works.

One should note that the R^2 values of the $(\nu_4 + \nu_5)^2$ forbidden bands cannot be modelled by Eqs. (5,6,8), because of the intensity asymmetry observed between the P - and R -branches, due to the ℓ -type resonance and Coriolis coupling already pointed out in Ref. [9]. Values in the Q -branch can be smoothed with enough precision for applications by using a polynomial of the second degree, but the coefficients of this polynomial have no physical meaning. Conversely, values in the P - and R -branches cannot be adjusted by a polynomial, even of high degree, without losing the precision of the experimental data. So, in order to generate a line list for applications, we thought preferable to keep the P - and R -experimental $|R|^2$ values and to simply smooth, by linear interpolations, the values of missing lines and of a few lines for which $|R|^2$ slightly departs from the general tendency of variation of neighboring lines. $|R|^2$ values of a few low J missing lines have also been obtained by this way. However, no extrapolation of line intensities toward high J should be performed for this band.

In Ref. [2], Kabbadj et al assigned 14 lines of the $(2\nu_4 + 2\nu_5)^0 - (\nu_4 + \nu_5)^2$ forbidden band, i.e., 11 R -lines (from $R13$ to $R23$), and 3 P -lines ($P11$, $P12$, and $P13$). All P -lines fall in very dense bunches of lines, so that we could neither assign new lines nor measure intensities in this branch. Conversely, it has been possible to measure the intensities of 8 R -lines already known from Kabbadj et al, plus 3 new R -lines ($R10$, $R11$, and $R12$). Experimental $|R|^2$ values have been obtained from the measured intensities using Eqs. (1-3) and could be reduced using Eqs. (5,8). The obtained constants are given in Table 5. These constants should not be considered as true vibrational transition moment squared and Herman-Wallis coefficients, since our set of experimental data is somewhat poor: intensities could be measured in only one branch, and the set of J values is limited and narrow (from 10 to 23). Therefore, the very small value of $|R_0|^2$ has a large absolute uncertainty, at least 20%. However, these constants are efficient to reproduce the data. For that, one needs to take into account 4 digits in these constants because of terms of high degree of m in Eqs. (5,8). Consequently, no extrapolation of line intensities toward high J should be performed. All results concerning this band are given tentatively.

The $(2\nu_4 + 2\nu_5)^2 \Pi - (\nu_4 + \nu_5)^0$ forbidden band was known by 28 lines assigned by Kabbadj et al [2], i.e., $P13$ to $P25$ and $R11$ to $R26$. We could measure line intensities for 43 lines in a more extended set of J values, from $P5$ to $P33$ and from $R3$ to $R32$. Experimental $|R|^2$ values have been obtained using Eqs. (1-3) but could not be fitted by Eqs. (5,8). However, as shown on Fig. 3, the precision of these values is good enough to allow to smooth them as performed for the $(\nu_4 + \nu_5)^2$ band.

The last forbidden band is $(2\nu_4 + \nu_5)^3 - \nu_4^1$, that is of the $\Phi \leftarrow \Pi$ type, and for which only f -levels are involved. Kabbadj et al [2] assigned 28 lines in this band, from $P16$ to $P29$, and from $R12$ to $R28$. We have been able to measure line intensities in a slightly more extended domain of J values, namely from $P10$ to $P32$, and from $R12$ to $R28$. Furthermore, the positions of some lines could be noticeably improved, e.g., $R13$, $R16$, or $R17$. However, numerous lines of this band are hidden by stronger ones of other bands. This is unfortunately the case for several lines of low J values (J smaller than 10). As for other forbidden bands, the experimental $|R|^2$ values could not be fitted by Eqs. (5,8), so that we needed to smooth the values by linear

interpolations. As overlappings and blends are particularly numerous for the lines of this band, we had to remove those of them for which some difficulty in the measurement process would had as consequence to lead to too dispersed $|R|^2$ values. Then, interpolations were easier.

The $(2\nu_4 + \nu_5)^1 \text{II} - \nu_4^1$ band is one of the strongest bands after $(\nu_4 + \nu_5)^0_+$ and $(\nu_4 + \nu_5)^2$. Its strongest lines (having J between about 4 and 18) are saturated in all our spectra and therefore could not be measured. However, a few lines with low J values ($J = 1, 2, \text{ or } 3$) could be measured. Added to numerous lines having J greater than about 19 and up to $Pe37$ and $Rf42$ (see Table 3), they allow a significant determination of a vibrational transition dipole moment squared and Herman-Wallis coefficients. Indeed, the transition dipole moment squared has no important rotational dependence for this band (see Fig. 4), so that the lack of data in the middle of P - and R -branches is not a drawback. The two e and f sub-bands have been fitted simultaneously, giving rise to a common value for the vibrational transition dipole moment squared but to distinct Herman-Wallis coefficients (see Table 5). Note that A_2^{PR} of the e sub-band could not be determined significantly and was fixed at zero. Such a data reduction is typical of the other bands with ℓ -type doubling. Figure 5 is an interesting example showing the strong and opposite rotational dependences of the sub-bands of $(2\nu_4 + \nu_5)^1 \text{I} - \nu_4^1$.

The remaining bands presented in this paper have been studied in the same manner. The spectra being very crowded, one should note that numerous lines could not be measured, e.g., because of unresolved overlappings, but that it was not a drawback to account for the rotational dependence of $|R|^2$ (see Figs. 4,5) for most of the bands. The complete list of line intensities is given in the supplementary material.

4. Line list for applications

The obtained results show that a global model, taking into account the numerous resonances, will be necessary to generate a precise synthetic spectrum, especially to account for the case of the forbidden bands. However, such a model cannot operate so long as line intensities of the other remaining weak bands are not measured. That is why, to benefit from the results of the present work, we choose to set up a line list in the HITRAN format that could be used even now for applications. Presently, HITRAN only contains data for the $(\nu_4 + \nu_5)^0_+$ band from Ref. [9].

The line list has been established in the same way as in Refs. [27,28]. As the present work was not specifically dedicated to line position measurements, we have relied on the best values existing in the literature. For the $(\nu_4 + \nu_5)^0_+$ band, we have picked up the line positions calculated by Vander Auwera [9] using the constants of Ref. [2], as they are given in HITRAN [17] from $P35$ to $R35$. For the other bands, the experimental line positions of Kabbadj et al [2] have been retained. For lines not observed in Refs. [2,9], as for a few lines that could be treated in a better way in our spectra, the positions obtained by the multispectrum procedure have been taken. For the missing lines and for a few extrapolated ones, positions have been obtained using an effective model based on a polynomial adjustment of the experimental line positions of Refs. [2,9] together with additional lines of this work if any.

To generate line intensities, vibrational transition dipole moments squared and Herman-Wallis coefficients (see Table 5) have been used to calculate $|R|^2$ values and then the intensities. Table 3 and the supplementary material show that the residuals are smaller than 3% on the average, and therefore are acceptable for usual applications. Line intensities corresponding to a

few higher values of J (about 5 additional values) have been extrapolated for some bands, their uncertainty codes being degraded consequently. To calculate the intensities of these additional lines, $|R|^2$ has been fixed at the value calculated for the highest observed J line in each band. As far as the $(\nu_4 + \nu_5)_+^0$ band is concerned, intensities calculated by Vander Auwera [9] from his absolute measurements, through a rigorous treatment, have been picked up. Line intensities of additional lines measured in this work for this band have been calculated using the constants of Table 5. Note that the statistical weights and the Einstein-A coefficients already present in HITRAN have also been picked up for these lines (Einstein-A coefficients have been put in place of $|R|^2$ values in this case). For the $(\nu_4 + \nu_5)^2$ forbidden band, smoothed values of $|R|^2$ have been used to calculate line intensities. These smoothed values can be found in the line list. No extrapolation toward high J has been performed for this band. For the $(2\nu_4 + 2\nu_5)_-^0 - (\nu_4 + \nu_5)^2$ forbidden band, the constants of Table 5 have been used, without extrapolation toward high J . Low J missing lines have been calculated in the R -branch, as well as P -lines from $P2$ to the last line, $P13$, observed in Ref. [2]. Calculated intensities of these lines have an absolute uncertainty of at least 20% and are given tentatively in the line list. For the $(2\nu_4 + \nu_5)^3 - \nu_4^1$ forbidden band, we have proceeded similarly. However, because of the lack of measurements between $P10$ and $R12$, the extrapolation of $|R|^2$ is difficult for the missing low J lines. As is usually done when one wants to extrapolate the intensities of weak high J lines, we chose not to take into account the rotational dependence of $|R|^2$ between $P10$ and $R12$: in the P -branch and for J smaller than 10, $|R|^2$ is fixed at the value measured for $P10$, whereas in the R -branch and for J smaller than 13, $|R|^2$ is fixed at the value measured for $R13$. This has no great consequence for applications since the concerned lines are weak, but the uncertainty of their intensity given in the line list has been increased (possibly more than 20%).

Other spectroscopic data are the same as those already put in the last updates of the databases: air- and self-broadening coefficients, default value for the temperature exponent of air-broadening coefficients, constant value for the air-pressure shifting coefficient, and their accuracies [17,29]. Table 6 summarizes the data now available for the $^{12}\text{C}_2\text{H}_2$ molecule in the studied spectral region. To facilitate comparisons, this table has been set up in the same format as in Refs. [27,28]. This line list contains 784 transitions and is included in the supplementary material.

5. Conclusion

Absolute intensities have been measured for more than 400 lines in 8 bands of the $^{12}\text{C}_2\text{H}_2$ spectrum around $7.7\ \mu\text{m}$, including the strongest bands. For allowed bands, an effective model of Herman-Wallis factor has been applied to model the spectrum with a satisfactory precision (better than 5% on the mean). For some forbidden bands, experimental values had to be smoothed by linear interpolations. A line list containing 784 transitions has been set up in the HITRAN format, which can be helpful for applications even now. However, an important work has still to be performed to study line intensities in some ten weaker hot bands observable in the recorded spectra. Once acquired, such a set of line intensities can be treated by a global model, as this of Perevalov et al (see, e.g., Refs. [25,26]), able to take into account the numerous resonances between levels involved in the $7.7\ \mu\text{m}$ spectral region and to account for the forbidden bands.

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Captions of Tables

Table 1

List of the bands observed by Kabbadj et al [2] in the $\Delta P = 2$ series of transitions of $^{12}\text{C}_2\text{H}_2$ around $7.7\ \mu\text{m}$

Table 2

Experimental conditions and characteristics of the spectra recorded around $7.7\ \mu\text{m}$ using the rapid-scan interferometer in Paris (LADIR)

Table 3

Line positions and intensities for selected allowed bands of the $^{12}\text{C}_2\text{H}_2$ molecule in the $7.7\ \mu\text{m}$ region.

Table 4

Line positions and intensities for the forbidden $(\nu_4 + \nu_5)^2$ band of the $^{12}\text{C}_2\text{H}_2$ molecule in the $7.7\ \mu\text{m}$ region.

Table 5

Summary of $^{12}\text{C}_2\text{H}_2$ experimental vibrational transition dipole moments squared $|R_0|^2$ in D^2 ($1\ \text{D} = 3.33546 \times 10^{-30}\ \text{C m}$), and Herman-Wallis coefficients, see Eqs. (2-8), in the $7.7\ \mu\text{m}$ spectral region

Table 6

Summary of new bands and transitions available for the $^{12}\text{C}_2\text{H}_2$ molecule at $7.7\ \mu\text{m}$

Captions of Figures

Fig. 1. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^0_+$ band. Solid triangles are experimental values obtained in this work. Solid squares are experimental values deduced from intensities measured by Vander Auwera [9] (when several values, coming from different spectra, are given for a line in Ref. [9], they have been averaged). Open triangles are experimental values deduced from intensities measured by Lepère et al [11]. (Experimental intensities quoted in Refs. [9,11] were preliminarily converted for a sample of 100% of $^{12}\text{C}_2\text{H}_2$.) The solid line represents the values calculated using the constants reported in Table 5. They result from an unweighted fit including the experimental data of this work and those of Vander Auwera [9] (see Section 3).

Fig. 2. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^2$ band. Solid triangles are experimental values obtained in this work. Open triangles are experimental values deduced from intensities measured by Vander Auwera [9]. The upper plot shows P - and R -branches whereas the lower plot shows the Q -branch.

Fig. 3. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + 2\nu_5)^2 \text{ II} - (\nu_4 + \nu_5)^0_-$ band. Solid triangles are experimental values obtained in this work.

Fig. 4. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + \nu_5)^1 \text{ II} - \nu_4^1$ band. Solid triangles are experimental values obtained in the e sub-band, and open ones in the f sub-band. The lines represent the values calculated using the constants reported in Table 5. The solid line is for the e sub-band, and the dashed one for the f sub-band.

Fig. 5. Variation of the transition dipole moment squared $|R|^2$, vs. m , for the $(2\nu_4 + \nu_5)^1 \text{ I} - \nu_4^1$ band. See caption of Fig. 4.

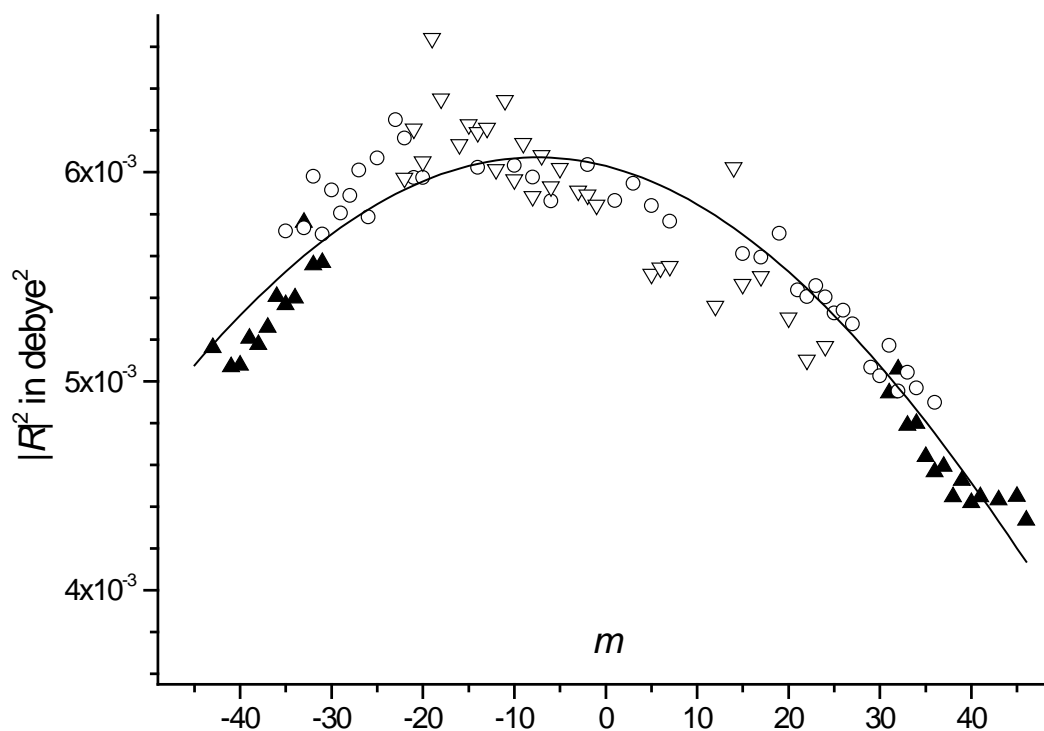


Fig. 1. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^0_+$ band. Solid triangles are experimental values obtained in this work. Open circles are experimental values deduced from intensities measured by Vander Auwera [9] (when several values, coming from different spectra, are given for a line in Ref. [9], they have been averaged). Open triangles are experimental values deduced from intensities measured by Lepère et al [11]. (Experimental intensities quoted in Refs. [9,11] were preliminarily converted for a sample of 100% of $^{12}\text{C}_2\text{H}_2$.) The solid line represents the values calculated using the constants reported in Table 5. They result from an unweighted fit including the experimental data of this work and those of Vander Auwera [9] (see Section 3).

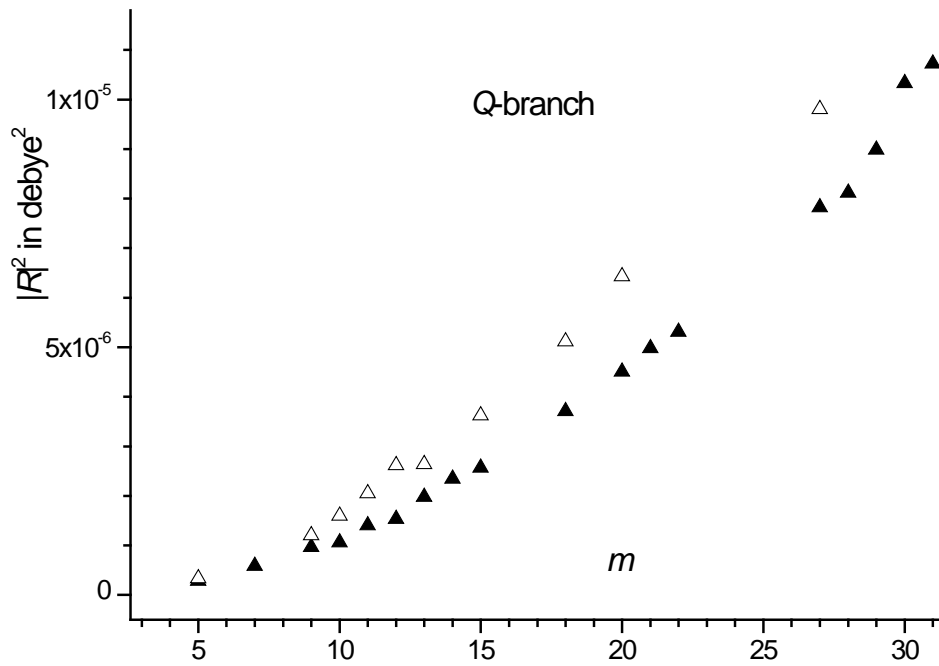
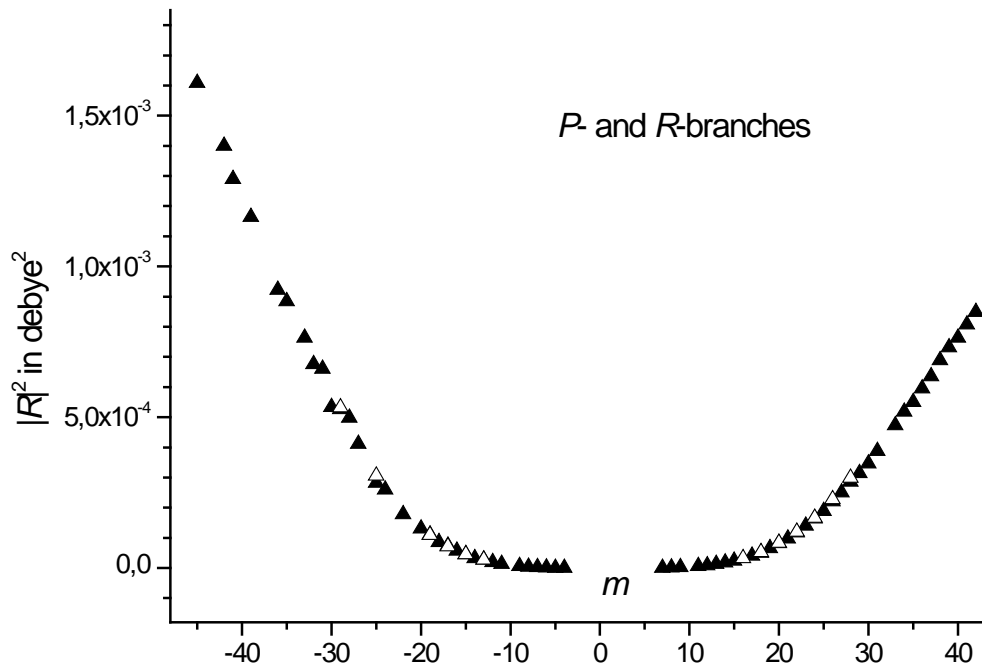


Fig. 2. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^2$ band. Solid triangles are experimental values obtained in this work. Open triangles are experimental values deduced from intensities measured by Vander Auwera [9]. The upper plot shows *P*- and *R*-branches whereas the lower plot shows the *Q*-branch.

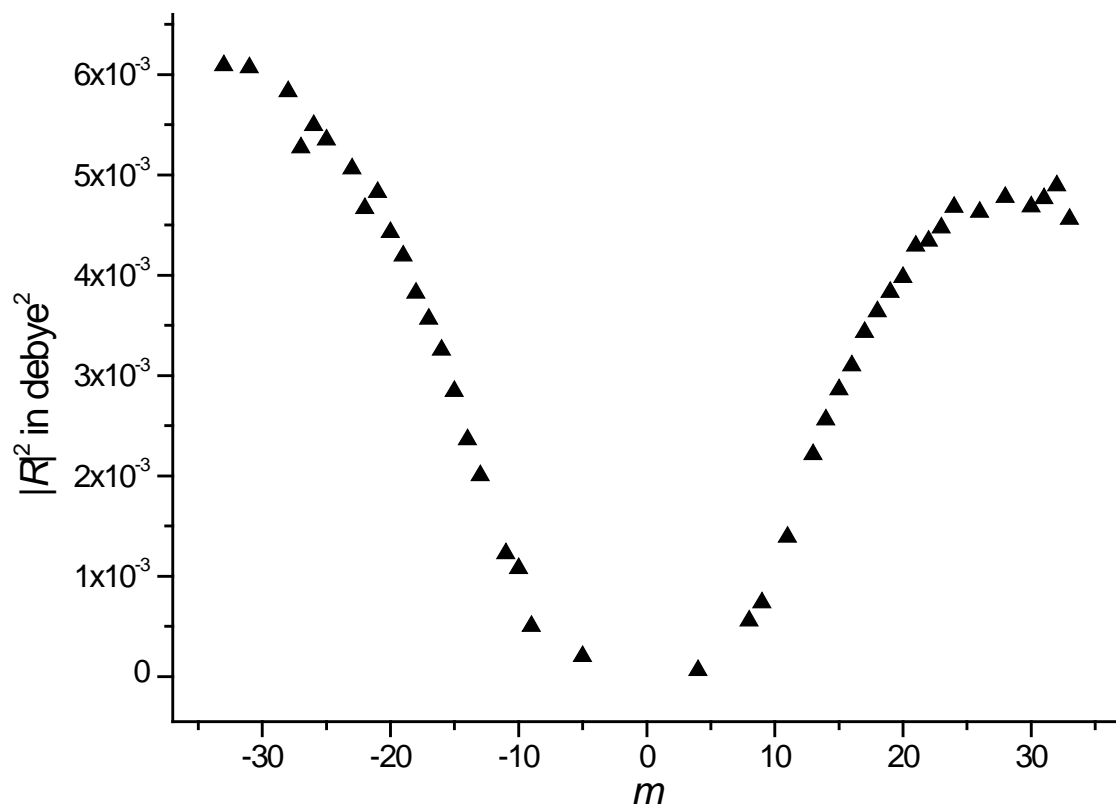


Fig. 3. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + 2\nu_5)^2 \Pi - (\nu_4 + \nu_5)^0_-$ band. Solid triangles are experimental values obtained in this work.

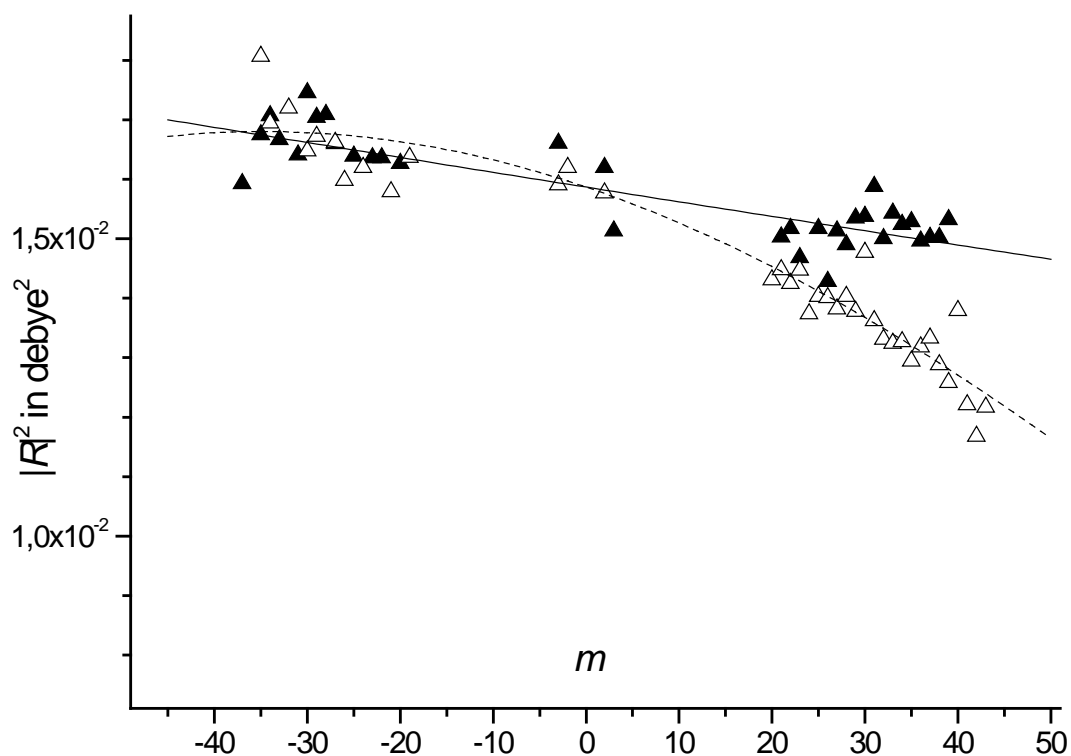


Fig. 4. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + \nu_5)^1 \text{ II} - \nu_4^1$ band. Solid triangles are experimental values obtained in the e sub-band, and open ones in the f sub-band. The lines represent the values calculated using the constants reported in Table 5. The solid line is for the e sub-band, and the dashed one for the f sub-band.

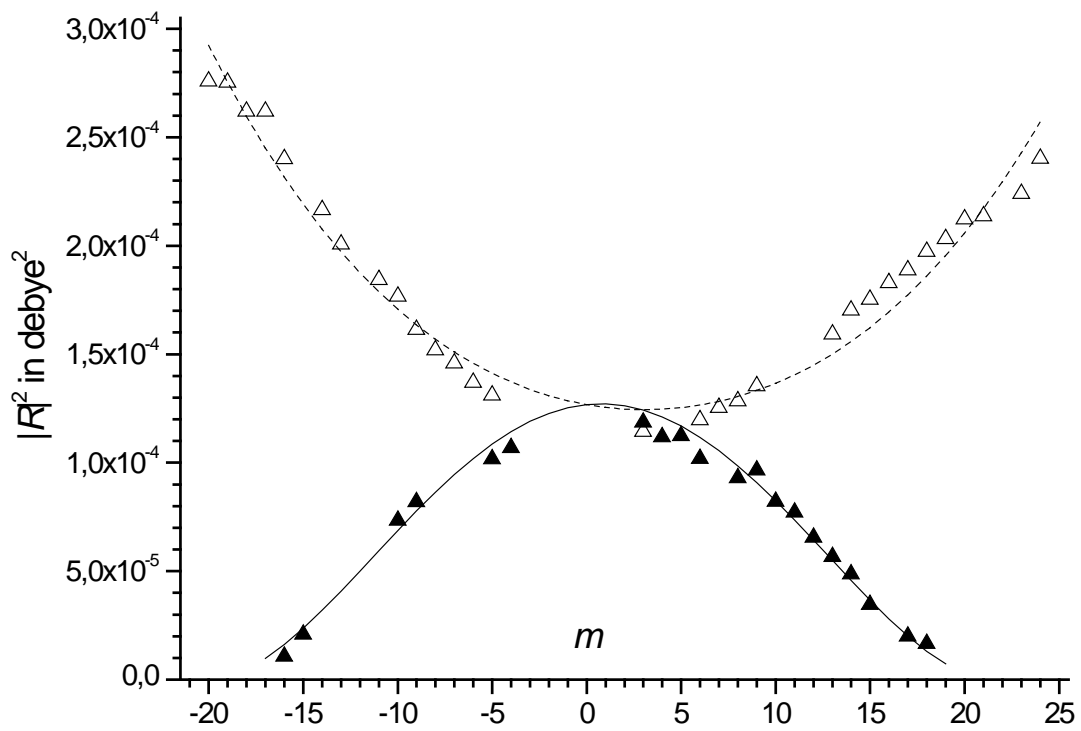


Fig. 5. Variation of the transition dipole moment squared $|R|^2$, vs. m , for the $(2\nu_4 + \nu_5)^1 I - \nu_4^1$ band. See caption of Fig. 4.

Table 1

List of the bands observed by Kabbadj et al [2] in the $\Delta P = 2$ series of transitions of $^{12}\text{C}_2\text{H}_2$ around $7.7 \mu\text{m}$

Band	Center ^a	Upper level ^b	Polyad ^b	Symmetry
$\nu_2 - \nu_5^1$	1245.140	$010(00)_+^0$	$\{3\nu_5\}$	$\Sigma_g^+ \leftarrow \Pi_u$
$(\nu_4 + 3\nu_5)_+^0 - 2\nu_5^0$	1308.686	$000(13)_+^0$	$\{4\nu_5\}$	$\Sigma_u^+ \leftarrow \Sigma_g^+$
$(\nu_4 + 3\nu_5)^2 - 2\nu_5^2$	1310.182	$000(13)^2$	$\{4\nu_5\}$	$\Delta_u \leftarrow \Delta_g$
$(2\nu_4 + 2\nu_5)_-^0 - (\nu_4 + \nu_5)^2$	1311.308	$000(22)_-^0 f$	$\{4\nu_5\}$	$\Sigma_g^- \leftarrow \Delta_u$
$(2\nu_4 + 2\nu_5)^2 \Pi - (\nu_4 + \nu_5)^2$	1318.652	$000(22)^2 \Pi$	$\{4\nu_5\}$	$\Delta_g \leftarrow \Delta_u$
$(\nu_4 + 2\nu_5)^1 \Pi - \nu_5^1$	1318.730	$000(12)^1 \Pi$	$\{3\nu_5\}$	$\Pi_g \leftarrow \Pi_u$
$(2\nu_4 + 2\nu_5)_+^0 - (\nu_4 + \nu_5)_+^0$	1319.942	$000(22)_+^0$	$\{4\nu_5\}$	$\Sigma_g^+ \leftarrow \Sigma_u^+$
$(2\nu_4 + 2\nu_5)_-^0 - (\nu_4 + \nu_5)_-^0$	1320.638	$000(22)_-^0$	$\{4\nu_5\}$	$\Sigma_g^- \leftarrow \Sigma_u^-$
$(2\nu_4 + 2\nu_5)^2 \Pi - (\nu_4 + \nu_5)_-^0$	1320.908	$000(22)^2 \Pi f$	$\{4\nu_5\}$	$\Delta_g \leftarrow \Sigma_u^-$
$(3\nu_4 + \nu_5)^2 - 2\nu_4^2$	1328.019	$000(31)^2$	$\{4\nu_5\}$	$\Delta_u \leftarrow \Delta_g$
$(\nu_4 + \nu_5)_+^0$	1328.081	$000(11)_+^0$	$\{2\nu_5\}$	$\Sigma_u^+ \leftarrow \Sigma_g^+$
$(2\nu_4 + \nu_5)^1 \Pi - \nu_4^1$	1328.314	$000(21)^1 \Pi$	$\{3\nu_5\}$	$\Pi_u \leftarrow \Pi_g$
$(3\nu_4 + \nu_5)_+^0 - 2\nu_4^0$	1330.206	$000(31)_+^0$	$\{4\nu_5\}$	$\Sigma_u^+ \leftarrow \Sigma_g^+$
$(\nu_4 + 2\nu_5)^1 \Pi - \nu_5^1$	1336.644	$000(12)^1 \Pi$	$\{3\nu_5\}$	$\Pi_g \leftarrow \Pi_u$
$(\nu_4 + \nu_5)^2$	1342.821	$000(11)^2 e$	$\{2\nu_5\}$	$\Delta_u \leftarrow \Sigma_g^+$
$(2\nu_4 + \nu_5)^1 \Pi - \nu_4^1$	1348.006	$000(21)^1 \Pi$	$\{3\nu_5\}$	$\Pi_u \leftarrow \Pi_g$
$(2\nu_4 + \nu_5)^3 - \nu_4^1$	1350.273	$000(21)^3 f$	$\{3\nu_5\}$	$\Phi_u \leftarrow \Pi_g$

^a Band centers, in cm^{-1} , have been compiled from Ref. [2].

^b For each band, the upper vibrational level and the polyad to which it belongs have been quoted. For the 4 forbidden bands, the sub-level (e or f) observed in Ref. [2] is mentioned. Levels are noted as explained in Section 1.

Table 2

Experimental conditions and characteristics of the spectra recorded around 7.7 μm using the rapid-scan interferometer in Paris (LADIR)

Commercial sample (Air Liquide Alphagaz)

Natural C ₂ H ₂	97.760 % of ¹² C ₂ H ₂
Stated purity	99.55 %
Maximum path difference	180 cm
Unapodized FWHM resolution	$\approx 2.8 \times 10^{-3} \text{ cm}^{-1}$
Spectral step after post-zero filling	$6.3 \times 10^{-4} \text{ cm}^{-1}$
Number of coadded scans	400
SNR in the involved spectral domain	≈ 350
Collimator focal length	418 mm
Nominal iris radius	0.85 mm
Adjusted effective iris radius	1.00(4) mm ^a

Spectrum number	Total pressure (hPa) $\pm 0.5\%$ ^b	Absorbing path (cm) $\pm 1 \text{ cm}$ ^b	Temperature (K) $\pm 0.5 \text{ K}$ ^b
1	7.585	2015	297.15
2	2.283	2015	296.45
3	1.254	2015	297.15
4	0.8185	2015	297.85
5	0.4283	2015	298.35
6	0.1971	2015	298.95

^a 1 SD between parentheses in unit of the last quoted digit.

^b Absolute uncertainty (excess digits are given as a guide).

Table 3

Line positions and intensities for selected allowed bands of the $^{12}\text{C}_2\text{H}_2$ molecule in the $7.7\ \mu\text{m}$ region.^a

Line	Position	S_{obs}	S_{calc}	%	$ R ^2_{\text{obs}}$
$(\nu_4 + \nu_5)_+$					
Pee43	1229.63445	1.687E-23	1.692E-23	-0.30	5.161E-03
Pee41	1234.33025	4.174E-23	4.339E-23	-3.95	5.069E-03
Pee40	1236.66863	2.172E-23	2.274E-23	-4.70	5.076E-03
Pee39	1239.00039	1.029E-22	1.060E-22	-3.01	5.205E-03
Pee38	1241.32569	5.194E-23	5.422E-23	-4.39	5.176E-03
Pee37	1243.64441	2.380E-22	2.465E-22	-3.57	5.258E-03
Pee36	1245.95659	1.212E-22	1.230E-22	-1.49	5.405E-03
Pee35	1248.26226	5.300E-22	5.457E-22	-2.96	5.366E-03
Pee34	1250.56156	2.578E-22	2.657E-22	-3.06	5.398E-03
Pee33	1252.85460	1.183E-21	1.150E-21	2.79	5.760E-03
Pee32	1255.14166	5.382E-22	5.458E-22	-1.41	5.557E-03
Pee31	1257.42288	2.262E-21	2.304E-21	-1.86	5.567E-03
Ree30	1402.99255	1.065E-21	1.082E-21	-1.60	4.944E-03
Ree31	1405.37220	2.373E-21	2.332E-21	1.73	5.058E-03
Ree32	1407.74330	5.370E-22	5.515E-22	-2.70	4.788E-03
Ree33	1410.10556	1.143E-21	1.159E-21	-1.40	4.797E-03
Ree34	1412.45877	2.579E-22	2.672E-22	-3.61	4.639E-03
Ree35	1414.80270	5.265E-22	5.476E-22	-4.01	4.567E-03
Ree36	1417.13728	1.205E-22	1.231E-22	-2.16	4.592E-03
Ree37	1419.46238	2.363E-22	2.463E-22	-4.23	4.447E-03
Ree38	1421.77794	5.346E-23	5.404E-23	-1.08	4.526E-03
Ree39	1424.08394	1.031E-22	1.054E-22	-2.23	4.418E-03
Ree40	1426.38037	2.252E-23	2.256E-23	-0.18	4.446E-03
Ree42	1430.94472	9.184E-24	8.970E-24	2.33	4.432E-03
Ree44	1435.47138	3.595E-24	3.395E-24	5.56	4.448E-03
Ree45	1437.72060	6.449E-24	6.154E-24	4.57	4.334E-03
$(2\nu_4 + \nu_5)^1 \Pi - \nu_4^1$					
Pee37	1243.41767	1.859E-23	1.961E-23	-5.49	1.593E-02
Pee35	1248.01886	4.260E-23	4.260E-23	0.00	1.675E-02
Pee34	1250.31348	2.099E-23	2.056E-23	2.05	1.707E-02
Pee33	1252.60446	8.802E-23	8.816E-23	-0.16	1.667E-02
Pee31	1257.17618	1.713E-22	1.738E-22	-1.46	1.641E-02
Pee30	1259.45762	8.384E-23	7.982E-23	4.79	1.746E-02
Pee29	1261.73641	3.347E-22	3.260E-22	2.60	1.704E-02
Pee28	1264.01299	1.506E-22	1.460E-22	3.05	1.709E-02
Pee27	1266.28762	5.836E-22	5.817E-22	0.33	1.660E-02
Pee25	1270.83268	9.797E-22	9.860E-22	-0.64	1.639E-02
Pee23	1275.37468	1.579E-21	1.587E-21	-0.51	1.636E-02
Pee22	1277.64554	6.551E-22	6.575E-22	-0.37	1.636E-02
Pee20	1282.18879	9.706E-22	9.765E-22	-0.61	1.627E-02
Pee 3	1321.27042	4.209E-21	4.040E-21	4.02	1.661E-02
Ree 1	1333.02342	2.466E-21	2.408E-21	2.35	1.620E-02
Ree 2	1335.38979	1.337E-21	1.396E-21	-4.41	1.513E-02

Table 3 (continued)

Line	Position	S_{obs}	S_{calc}	%	$ R ^2_{\text{obs}}$
Ree20	1378.59278	1.013E-21	1.035E-21	-2.17	1.503E-02
Ree21	1381.00425	2.535E-21	2.562E-21	-1.07	1.517E-02
Ree22	1383.41398	6.664E-22	6.948E-22	-4.26	1.468E-02
Ree24	1388.22619	4.395E-22	4.420E-22	-0.57	1.517E-02
Ree25	1390.62765	9.722E-22	1.037E-21	-6.67	1.428E-02
Ree26	1393.02574	2.655E-22	2.669E-22	-0.53	1.513E-02
Ree27	1395.41954	5.992E-22	6.106E-22	-1.90	1.490E-02
Ree28	1397.80900	1.551E-22	1.532E-22	1.23	1.535E-02
Ree29	1400.19348	3.473E-22	3.418E-22	1.58	1.538E-02
Ree30	1402.57259	8.790E-23	8.365E-23	4.84	1.588E-02
Ree31	1404.94633	1.808E-22	1.818E-22	-0.55	1.500E-02
Ree32	1407.31374	4.450E-23	4.344E-23	2.38	1.543E-02
Ree33	1409.67516	9.342E-23	9.219E-23	1.32	1.524E-02
Ree34	1412.02985	2.187E-23	2.148E-23	1.78	1.529E-02
Ree35	1414.37759	4.441E-23	4.450E-23	-0.20	1.496E-02
Ree36	1416.71806	1.016E-23	1.012E-23	0.39	1.502E-02
Ree37	1419.05136	2.057E-23	2.047E-23	0.49	1.502E-02
Ree38	1421.37702	4.667E-24	4.545E-24	2.61	1.532E-02
Pff35	1250.12963	1.487E-23	1.383E-23	6.99	1.807E-02
Pff34	1252.33119	6.076E-23	6.024E-23	0.86	1.695E-02
Pff32	1256.71873	1.253E-22	1.224E-22	2.31	1.720E-02
Pff30	1261.08877	2.323E-22	2.366E-22	-1.85	1.648E-02
Pff29	1263.26912	1.072E-22	1.076E-22	-0.37	1.672E-02
Pff27	1267.62446	1.912E-22	1.928E-22	-0.84	1.662E-02
Pff26	1269.80108	7.247E-22	7.594E-22	-4.79	1.598E-02
Pff24	1274.15638	1.220E-21	1.259E-21	-3.20	1.620E-02
Pff21	1280.70596	7.687E-22	8.109E-22	-5.49	1.579E-02
Pff19	1285.09156	1.157E-21	1.174E-21	-1.47	1.637E-02
Pff 3	1321.25387	1.343E-21	1.352E-21	-0.67	1.591E-02
Pff 2	1323.59436	2.393E-21	2.359E-21	1.42	1.620E-02
Rff 1	1333.05919	8.000E-22	7.995E-22	0.06	1.577E-02
Rff19	1377.39542	1.142E-21	1.160E-21	-1.58	1.431E-02
Rff20	1379.91348	2.900E-21	2.895E-21	0.17	1.448E-02
Rff21	1382.43330	7.850E-22	7.917E-22	-0.85	1.425E-02
Rff22	1384.95401	1.947E-21	1.923E-21	1.23	1.447E-02
Rff23	1387.47488	4.949E-22	5.117E-22	-3.39	1.374E-02
Rff24	1389.99477	1.202E-21	1.210E-21	-0.67	1.403E-02
Rff25	1392.51314	3.133E-22	3.138E-22	-0.16	1.401E-02
Rff26	1395.02890	7.160E-22	7.225E-22	-0.91	1.382E-02
Rff27	1397.54132	1.848E-22	1.825E-22	1.24	1.403E-02
Rff28	1400.04937	4.099E-22	4.095E-22	0.10	1.378E-02
Rff29	1402.55207	1.089E-22	1.008E-22	7.44	1.477E-02
Rff30	1405.04898	2.212E-22	2.206E-22	0.27	1.362E-02
Rff31	1407.53887	5.229E-23	5.300E-23	-1.36	1.331E-02
Rff32	1410.02106	1.118E-22	1.131E-22	-1.16	1.324E-02
Rff33	1412.49475	2.642E-23	2.648E-23	-0.23	1.327E-02
Rff34	1414.95929	5.404E-23	5.514E-23	-2.04	1.294E-02
Rff35	1417.41389	1.266E-23	1.259E-23	0.55	1.318E-02
Rff36	1419.85815	2.621E-23	2.558E-23	2.40	1.333E-02
Rff37	1422.29128	5.689E-24	5.702E-24	-0.23	1.288E-02
Rff38	1424.71292	1.110E-23	1.130E-23	-1.80	1.258E-02

Table 3 (continued)

Line	Position	S_{obs}	S_{calc}	%	$ R _{\text{obs}}^2$
Rff39	1427.12270	2.669E-24	2.459E-24	7.87	1.379E-02
Rff40	1429.52048	4.608E-24	4.757E-24	-3.23	1.221E-02
Rff41	1431.90577	9.444E-25	1.011E-24	-7.05	1.168E-02
Rff42	1434.27815	1.874E-24	1.909E-24	-1.87	1.217E-02

^a The quoted line position is this measured in this work, in cm^{-1} . S_{obs} and S_{calc} are measured and calculated intensities, respectively, for pure $^{12}\text{C}_2\text{H}_2$ (i.e., for a sample containing 100% of $^{12}\text{C}_2\text{H}_2$), in cm molecule^{-1} at 296 K. % is the ratio $100 \times (S_{\text{obs}} - S_{\text{calc}}) / S_{\text{obs}}$. $|R|_{\text{obs}}^2$ is the experimental transition dipole moment squared value, in D^2 ($1 \text{ D} = 3.33546 \times 10^{-30} \text{ C m}$), deduced from S_{obs} using Eqs. (1-3).

Table 4

Line positions and intensities for the forbidden $(\nu_4 + \nu_5)^2$ band of the $^{12}\text{C}_2\text{H}_2$ molecule in the 7.7 μm region.^a

Line	Position	S_{obs}	$ R ^2_{\text{obs}}$	$ R ^2_{\text{smooth}}$	%
Pee45	1248.78310	2.035E-24	1.608E-03	1.608E-03	0.00
Pee42	1253.91204	2.480E-24	1.400E-03	1.369E-03	2.21
Pee41	1255.66410	1.084E-23	1.289E-03	1.289E-03	0.00
Pee39	1259.23186	2.337E-23	1.164E-03	1.164E-03	0.00
Pee36	1264.74491	2.099E-23	9.223E-04	9.223E-04	0.00
Pee35	1266.62567	8.859E-23	8.838E-04	8.730E-04	1.22
Pee33	1270.45154	1.591E-22	7.636E-04	7.417E-04	2.87
Pee32	1272.39624	6.640E-23	6.761E-04	6.761E-04	0.00
Pee31	1274.36192	2.715E-22	6.605E-04	6.046E-04	8.46
Pee30	1276.34806	1.009E-22	5.330E-04	5.330E-04	0.00
Pee29	1278.35433	4.079E-22	5.264E-04	4.838E-04	8.09
Pee28	1280.38115	1.728E-22	4.973E-04	4.346E-04	12.61
Pee27	1282.42648	5.706E-22	4.115E-04	3.855E-04	6.32
Pee25	1286.57394	6.655E-22	2.817E-04	2.978E-04	-5.72
Pee24	1288.67440	2.612E-22	2.593E-04	2.593E-04	0.00
Pee22	1292.92634	2.822E-22	1.781E-04	1.781E-04	0.00
Pee20	1297.24109	3.075E-22	1.300E-04	1.300E-04	0.00
Pee19	1299.42044	9.323E-22	1.100E-04	1.100E-04	0.00
Pee18	1301.61362	2.861E-22	8.591E-05	8.591E-05	0.00
Pee17	1303.81966	8.572E-22	7.384E-05	7.384E-05	0.00
Pee16	1306.03816	2.538E-22	5.734E-05	5.734E-05	0.00
Pee15	1308.26857	6.702E-22	4.471E-05	4.471E-05	0.00
Pee14	1310.51011	1.807E-22	3.265E-05	3.265E-05	0.00
Pee13	1312.76231	5.009E-22	2.760E-05	2.760E-05	0.00
Pee12	1315.02439	1.216E-22	1.879E-05	1.879E-05	0.00
Pee11	1317.29587	2.607E-22	1.272E-05	1.272E-05	0.00
Pee 9	1321.86562	1.309E-22	6.117E-06	6.117E-06	0.00
Pee 8	1324.16302	2.630E-23	3.733E-06	3.733E-06	0.00
Pee 7	1326.46831	4.587E-23	2.261E-06	2.261E-06	0.00
Pee 6	1328.78108	7.072E-24	1.123E-06	1.123E-06	0.00
Pee 5	1331.10092	1.105E-23	6.514E-07	6.514E-07	0.00
Pee 4	1333.42772	8.858E-25	1.855E-07	1.855E-07	0.00
Qfe 5	1342.89810	5.326E-24	2.844E-07	3.146E-07	-10.62
Qfe 7	1342.98093	1.283E-23	5.811E-07	5.865E-07	-0.93
Qfe 9	1343.08942	2.207E-23	9.620E-07	9.439E-07	1.88
Qfe10	1343.15298	7.978E-24	1.057E-06	1.155E-06	-9.24
Qfe11	1343.22331	3.069E-23	1.404E-06	1.387E-06	1.23
Qfe12	1343.29997	1.062E-23	1.534E-06	1.640E-06	-6.93
Qfe13	1343.38257	3.806E-23	1.974E-06	1.915E-06	2.98
Qfe14	1343.47178	1.379E-23	2.344E-06	2.211E-06	5.65
Qfe15	1343.56714	4.082E-23	2.565E-06	2.529E-06	1.40
Qfe18	1343.89142	1.305E-23	3.707E-06	3.610E-06	2.61
Qfe20	1344.13921	1.127E-23	4.505E-06	4.438E-06	1.49
Qfe21	1344.27247	3.092E-23	4.977E-06	4.884E-06	1.87
Qfe22	1344.41219	8.953E-24	5.312E-06	5.351E-06	-0.74
Qfe27	1345.20371	1.164E-23	7.821E-06	8.008E-06	-2.39
Qfe28	1345.38070	3.017E-24	8.113E-06	8.604E-06	-6.05

Table 4 (continued)

Line	Position	S_{obs}	$ R _{\text{obs}}^2$	$ R _{\text{smooth}}^2$	%
Qfe29	1345.56395	7.464E-24	8.984E-06	9.220E-06	-2.63
Qfe30	1345.75320	2.096E-24	1.033E-05	9.859E-06	4.56
Qfe31	1345.98592	4.739E-24	1.072E-05	1.052E-05	1.88
Ree 6	1359.45680	4.732E-24	6.294E-07	6.294E-07	0.00
Ree 7	1361.86361	3.411E-23	1.431E-06	1.431E-06	0.00
Ree 8	1364.27775	1.961E-23	2.399E-06	2.399E-06	0.00
Ree10	1369.13001	4.938E-23	6.127E-06	6.127E-06	0.00
Ree11	1371.56944	2.074E-22	8.880E-06	8.880E-06	0.00
Ree12	1374.01690	9.446E-23	1.285E-05	1.285E-05	0.00
Ree13	1376.47437	3.649E-22	1.780E-05	1.780E-05	0.00
Ree14	1378.94166	1.486E-22	2.383E-05	2.383E-05	0.00
Ree16	1383.90889	2.037E-22	4.088E-05	4.088E-05	0.00
Ree17	1386.40773	6.574E-22	5.025E-05	5.025E-05	0.00
Ree18	1388.92278	2.463E-22	6.557E-05	6.557E-05	0.00
Ree19	1391.44885	7.775E-22	8.144E-05	8.144E-05	0.00
Ree20	1393.98831	2.600E-22	9.755E-05	9.755E-05	0.00
Ree21	1396.54189	7.762E-22	1.175E-04	1.175E-04	0.00
Ree22	1399.11015	2.496E-22	1.394E-04	1.394E-04	0.00
Ree23	1401.69371	7.102E-22	1.642E-04	1.642E-04	0.00
Ree24	1404.29302	2.156E-22	1.889E-04	1.889E-04	0.00
Ree25	1406.90872	5.911E-22	2.200E-04	2.200E-04	0.00
Ree26	1409.54126	1.733E-22	2.498E-04	2.498E-04	0.00
Ree27	1412.19099	4.534E-22	2.856E-04	2.856E-04	0.00
Ree28	1414.85824	1.250E-22	3.137E-04	3.137E-04	0.00
Ree29	1417.54325	3.080E-22	3.461E-04	3.461E-04	0.00
Ree30	1420.24715	8.452E-23	3.874E-04	3.874E-04	0.00
Ree32	1425.70969	5.366E-23	4.727E-04	4.727E-04	0.00
Ree33	1428.46892	1.250E-22	5.177E-04	5.177E-04	0.00
Ree34	1431.24710	3.097E-23	5.502E-04	5.502E-04	0.00
Ree35	1434.04395	6.952E-23	5.946E-04	5.946E-04	0.00
Ree36	1436.85956	1.694E-23	6.349E-04	6.349E-04	0.00
Ree37	1439.69380	3.707E-23	6.882E-04	6.882E-04	0.00
Ree38	1442.54669	8.757E-24	7.309E-04	7.309E-04	0.00
Ree39	1445.41781	1.812E-23	7.638E-04	7.638E-04	0.00
Ree40	1448.30716	4.153E-24	8.067E-04	8.067E-04	0.00
Ree41	1451.21437	8.433E-24	8.487E-04	8.487E-04	0.00

^a The quoted line position is this measured in this work, in cm^{-1} . S_{obs} is the measured intensity for pure $^{12}\text{C}_2\text{H}_2$ (i.e., for a sample containing 100% of $^{12}\text{C}_2\text{H}_2$), in cm molecule^{-1} at 296 K. $|R|_{\text{obs}}^2$ is the experimental transition dipole moment squared value deduced from S_{obs} using Eqs. (1-4), and $|R|_{\text{smooth}}^2$ its smoothed value, in D^2 ($1 \text{ D} = 3.33546 \times 10^{-30} \text{ C m}$). % is the ratio $100 \times (|R|_{\text{obs}}^2 - |R|_{\text{smooth}}^2) / |R|_{\text{obs}}^2$. Except for the Q-branch, this ratio should not be understood as an observed – calculated residual. When non zero, it means that the experimental value has been changed to set up the line list, and indicates how much.

Table 5

Summary of $^{12}\text{C}_2\text{H}_2$ experimental vibrational transition dipole moments squared $|R_0|^2$ in D^2 ($1 \text{ D} = 3.33546 \times 10^{-30} \text{ C m}$), and Herman-Wallis coefficients, see Eqs. (2-8), in the $7.7 \mu\text{m}$ spectral region ^a

Band	Origin	$ R_0 ^2$	A_1^{PR}	A_2^{PR}
$(2\nu_4 + 2\nu_5)^0_- - (\nu_4 + \nu_5)^2$	1311.308	$3.15(57) \times 10^{-7}$ ^b	$-5.377(460) \times 10^{-2}$	$+8.377(1700) \times 10^{-4}$
$(\nu_4 + 2\nu_5)^1 \text{ II} - \nu_5^1$	1318.730	$1.673(12) \times 10^{-2}$	$e -6.9(18) \times 10^{-4}$	$-1.19(70) \times 10^{-5}$
			$f -1.71(16) \times 10^{-3}$	
$(2\nu_4 + 2\nu_5)^2 \text{ II} - (\nu_4 + \nu_5)^0_-$ factors		1320.908	Impossible to adjust with Herman-Wallis factors	
$(\nu_4 + \nu_5)^0_+$	1328.081	$6.030(73) \times 10^{-3}$ ^c	$-9.2(11) \times 10^{-4}$	$-6.12(56) \times 10^{-5}$
$(2\nu_4 + \nu_5)^1 \text{ II} - \nu_4^1$	1328.314	$1.587(14) \times 10^{-2}$	$e -7.8(17) \times 10^{-4}$	
			$f -1.67(19) \times 10^{-3}$	$-2.40(69) \times 10^{-5}$
$(\nu_4 + \nu_5)^2$	1342.821	Impossible to adjust with Herman-Wallis factors		
$(2\nu_4 + \nu_5)^1 \text{ I} - \nu_4^1$	1348.006	$1.268(42) \times 10^{-4}$	$e +3.5(27) \times 10^{-3}$	$-2.29(23) \times 10^{-3}$
			$f -6.11(72) \times 10^{-3}$	$+9.91(74) \times 10^{-4}$
$(2\nu_4 + \nu_5)^3 - \nu_4^1$	1350.273	Impossible to adjust with Herman-Wallis factors		

^a 95% confidence intervals (2 SD, in unit of the last quoted digit) are given between parenthesis. For $|R_0|^2$ values, the overall accuracy is 5%, except for the $(2\nu_4 + 2\nu_5)^0_- - (\nu_4 + \nu_5)^2$ band for which it is estimated to be at least 20%. Non given Herman-Wallis coefficients have been fixed at zero.

^b For this band, $|R_0|^2$, A_1^{PR} and A_2^{PR} have a poor physical meaning. Excess digits in A_1^{PR} and A_2^{PR} are non significant but necessary to reproduce the experimental data (see Section 3). Note that Eq. (8) has been used for this forbidden band, contrary to the allowed bands for which Eq. (7) has been used.

^c $|R_0|^2 = 5.8229 \times 10^{-3} \text{ D}^2$ in Ref. [9], strictly speaking, not comparable with our value since obtained from a different theoretical model.

Table 6

Summary of new bands and transitions available for the $^{12}\text{C}_2\text{H}_2$ molecule at $7.7\ \mu\text{m}^{\text{a}}$

Band	Origin	$\nu_{\text{min}} - \nu_{\text{max}}$	ΣS	$S_{\text{min}} - S_{\text{max}}$	$J_{\text{max}\nu}/J_{\text{max}S}/J_{\text{max}}$	Cdv	CdS
000220--000112	1311.31	1288-1377	1.9E-21	1.0E-25-6.4E-22	13-23 ^b -23 13-23	4/2	5/3
000121_2-000011	1318.73	1225-1423	2.1E-19	1.0E-24-5.5E-21	37-38 37-38 40-40	4/3	6/5
000222_2-000110-	1320.91	1246-1403	1.2E-21	3.4E-25-7.4E-23	33-32 33-32 33-32	4/3	6/4
000110+_--000000+	1328.08	1224-1449	2.7E-18	1.5E-25-1.4E-19	43-45 43-45 45-50	4/3	7/5
000211_2-000101	1328.31	1224-1442	3.6E-19	1.5E-25-9.4E-21	39-42 37-42 45-45	4/3	6/5
000112__-000000+	1342.82	1248-1452	1.8E-20	2.1E-25-9.1E-22	45-41 45-41 45-41 31 31 31	4/3	6/5
000211_1-000101	1348.01	1288-1420	2.7E-21	1.8E-24-9.1E-23	22-26 20-23 25-30	4/3	6/5
000213__-000101	1350.27	1277-1421	9.4E-22	2.8E-24-4.8E-23	32-28 32-28 32-28	4/2	6/4

^a Explanation of the column headings:

Band: vibrational assignment used in the line list, according to Section 1: $\nu_1 \nu_2 \nu_3 \nu_4 \nu_5 \ell \pm r$ for the upper and lower states. When \pm or r does not occur for the upper state, it is replaced by an underscore. Note that r is mentioned only if necessary to avoid ambiguities.

Origin: approximate value of the band center, in cm^{-1} .

$\nu_{\text{min}} - \nu_{\text{max}}$: limiting values of line positions, in cm^{-1} .

ΣS : sum of line intensities, in cm molecule^{-1} at 296 K.

$S_{\text{min}} - S_{\text{max}}$: limiting values of line intensities, in cm molecule^{-1} at 296 K.

$J_{\text{max}\nu}$: maximum value of J for which a line position has been measured.

$J_{\text{max}S}$: maximum value of J for which a line intensity has been measured.

J_{max} : maximum value of J present in the line list.

(The first value is for the P -branch and the second for the R -branch. When a value is on a separate line, it concerns the Q -branch of the above band.)

Cdv: uncertainty code for line positions [17]. Code 2: 10^{-2} - $10^{-1}\ \text{cm}^{-1}$. Code 3: 10^{-3} - $10^{-2}\ \text{cm}^{-1}$. Code 4: 10^{-4} - $10^{-3}\ \text{cm}^{-1}$.

CdS: uncertainty code for line intensities [17]. Code 4: 10-20%. Code 5: 5-10%. Code 6: 2-5%. Code 7: 1-2%.

(The second value is for some interpolated or extrapolated lines.)

Other spectroscopic data are the same as those already put in the last updates of the databases: air- and self-broadening coefficients, default value for the temperature exponent of air-broadening coefficients, constant value for the air-pressure shifting coefficient, and their accuracies [17,28].

^b No intensity measurement in the P -branch of this band.

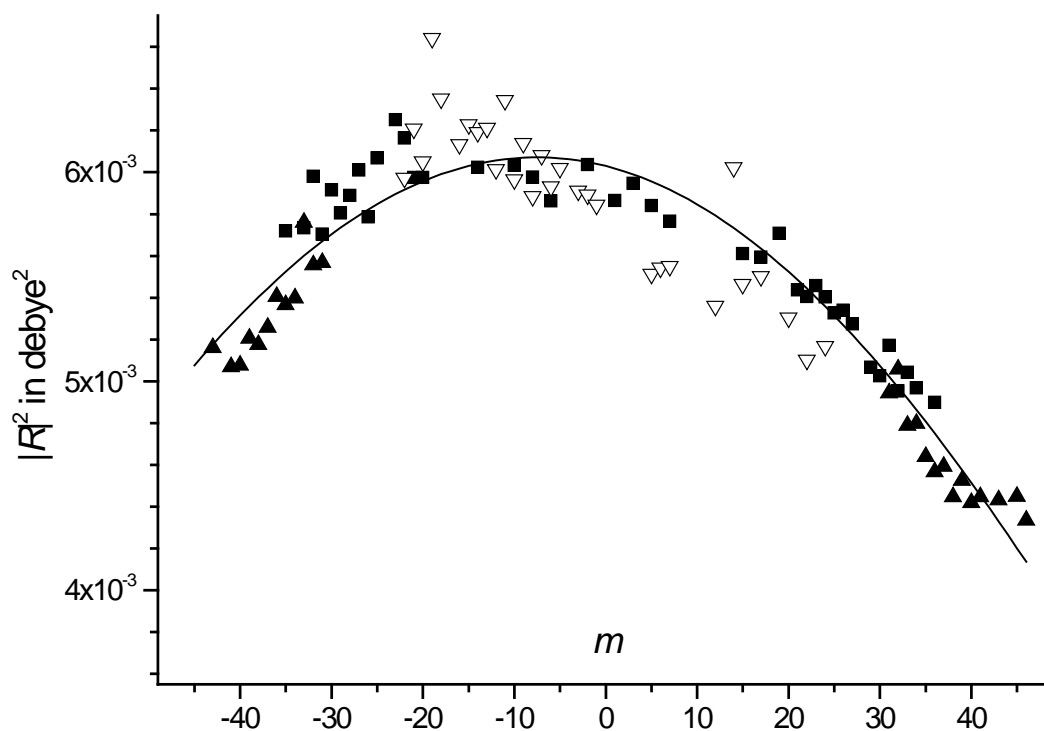


Fig. 1. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^0_+$ band. Solid triangles are experimental values obtained in this work. Solid squares are experimental values deduced from intensities measured by Vander Auwera [9] (when several values, coming from different spectra, are given for a line in Ref. [9], they have been averaged). Open triangles are experimental values deduced from intensities measured by Lepère et al [11]. (Experimental intensities quoted in Refs. [9,11] were preliminarily converted for a sample of 100% of $^{12}\text{C}_2\text{H}_2$.) The solid line represents the values calculated using the constants reported in Table 5. They result from an unweighted fit including the experimental data of this work and those of Vander Auwera [9] (see Section 3).

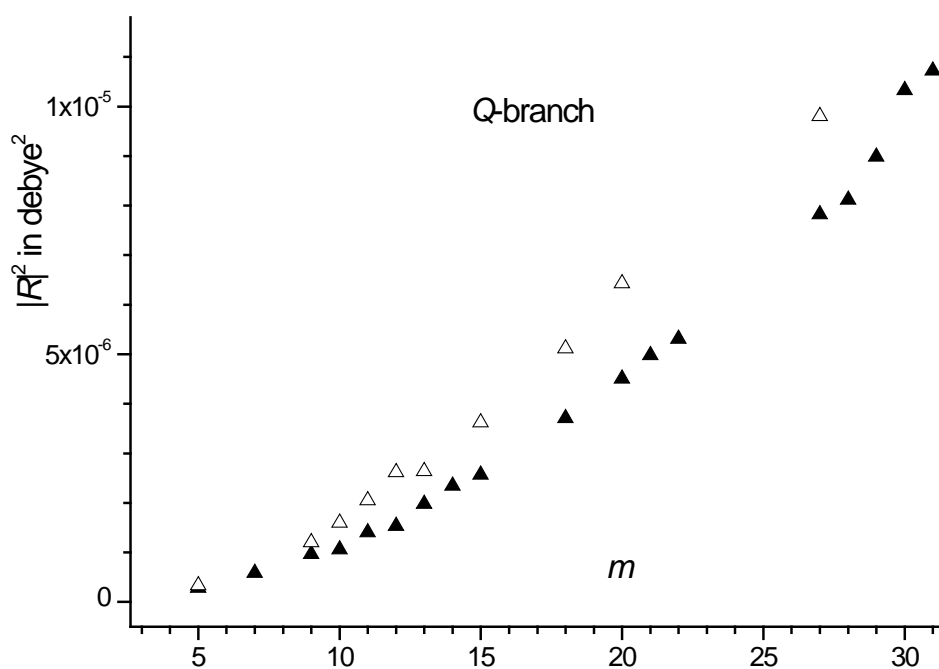
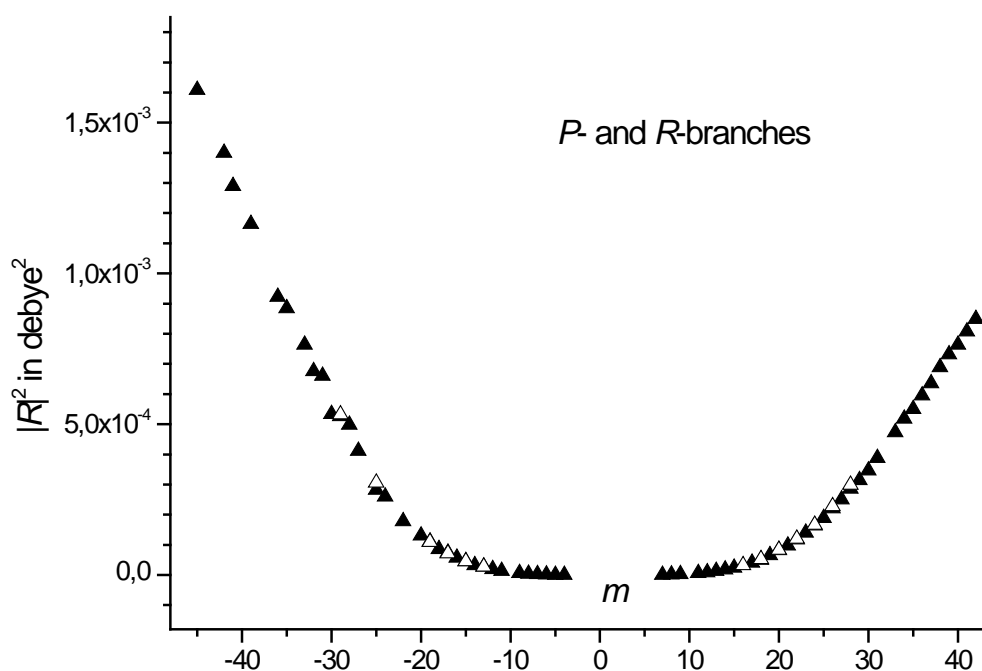


Fig. 2. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(\nu_4 + \nu_5)^2$ band. Solid triangles are experimental values obtained in this work. Open triangles are experimental values deduced from intensities measured by Vander Auwera [9]. The upper plot shows *P*- and *R*-branches whereas the lower plot shows the *Q*-branch.

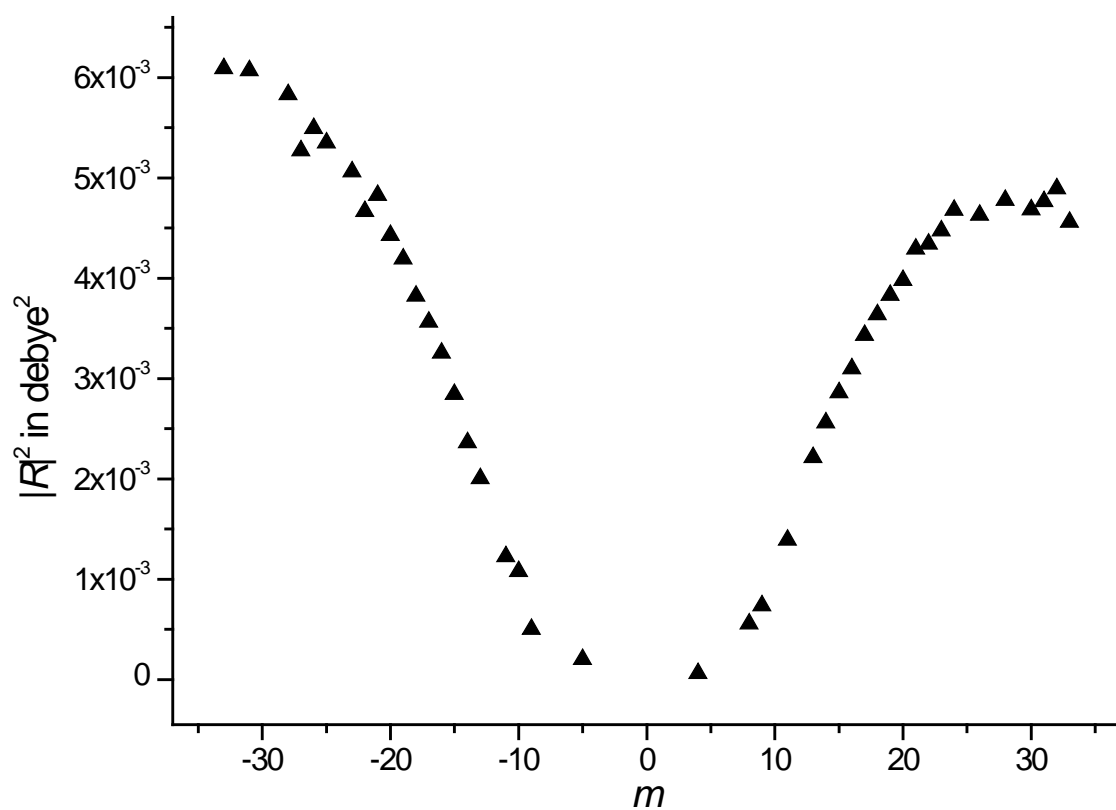


Fig. 3. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + 2\nu_5)^2 \Pi - (\nu_4 + \nu_5)^0_-$ band. Solid triangles are experimental values obtained in this work.

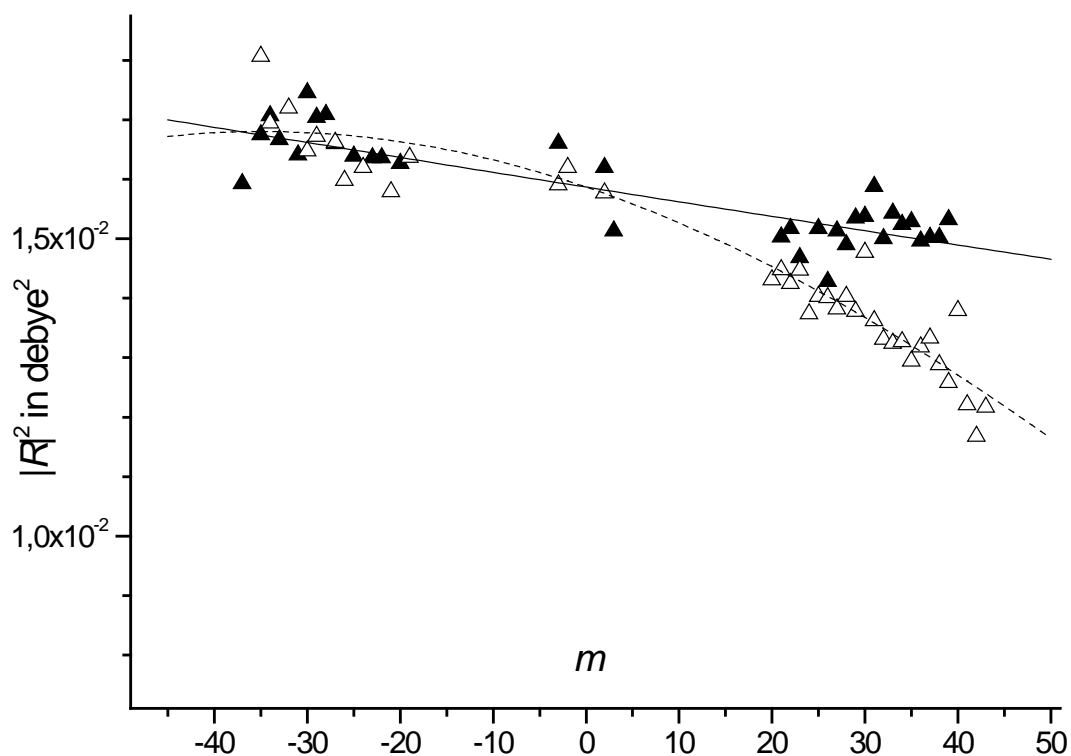


Fig. 4. Variation of the transition dipole moment squared $|R|^2$, in D^2 ($1 D = 3.33546 \times 10^{-30} \text{ C m}$), vs. m , for the $(2\nu_4 + \nu_5)^1 \text{ II} - \nu_4^1$ band. Solid triangles are experimental values obtained in the e sub-band, and open ones in the f sub-band. The lines represent the values calculated using the constants reported in Table 5. The solid line is for the e sub-band, and the dashed one for the f sub-band.

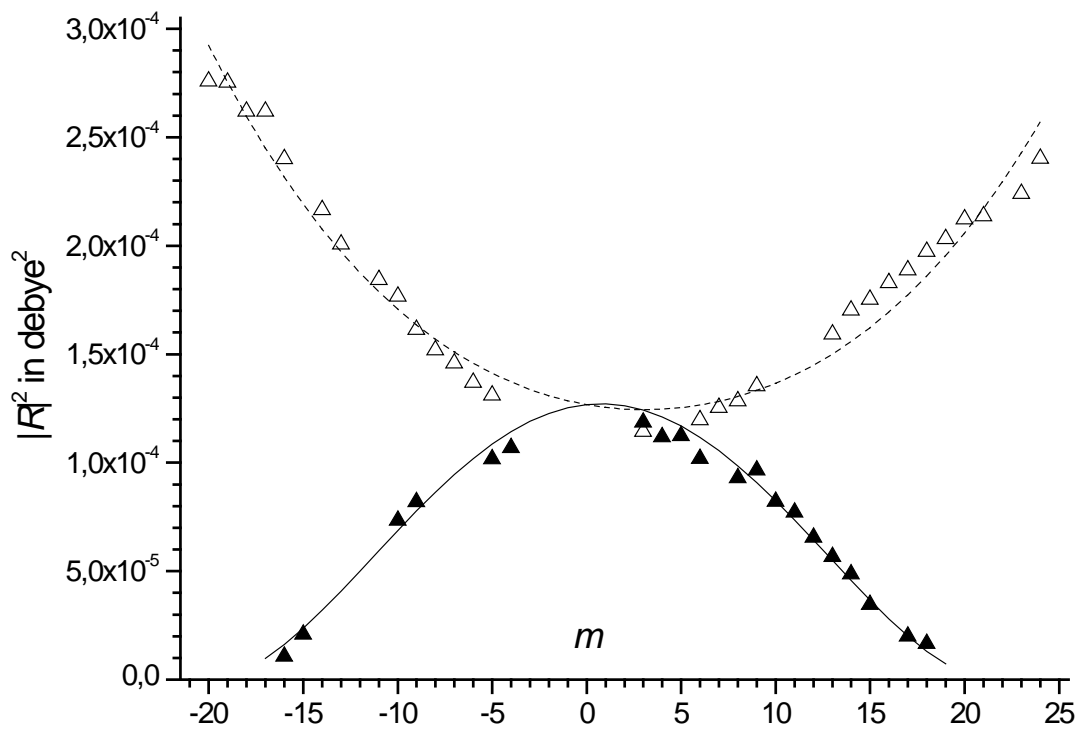


Fig. 5. Variation of the transition dipole moment squared $|R|^2$, vs. m , for the $(2\nu_4 + \nu_5)^1 I - \nu_4^1$ band. See caption of Fig. 4.