

# Cubic three-dimensional hybrid silica solids for nuclear hyperpolarization

D Baudouin, Henri A van Kalkeren, A Bornet, B Vuichoud, L Veyre, M. Cavaillès, Martin Schwarzwalder, W.-C Liao, D Gajan, G Bodenhausen, et al.

## ▶ To cite this version:

D Baudouin, Henri A van Kalkeren, A Bornet, B Vuichoud, L Veyre, et al.. Cubic three-dimensional hybrid silica solids for nuclear hyperpolarization . Chemical Science, 2016, 7 (11), pp.6846-6850 10.1039/c6sc02055k . hal-01360905

# HAL Id: hal-01360905 https://hal.sorbonne-universite.fr/hal-01360905v1

Submitted on 6 Sep 2016

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



Distributed under a Creative Commons Attribution 4.0 International License

# Chemical Science

# EDGE ARTICLE

CrossMark

Cite this: DOI: 10.1039/c6sc02055k

Received 11th May 2016 Accepted 15th July 2016 DOI: 10.1039/c6sc02055k

www.rsc.org/chemicalscience

# Cubic three-dimensional hybrid silica solids for nuclear hyperpolarization<sup>†</sup>

D. Baudouin,<sup>\*a</sup> H. A. van Kalkeren,<sup>a</sup> A. Bornet,<sup>b</sup> B. Vuichoud,<sup>b</sup> L. Veyre,<sup>a</sup> M. Cavaillès,<sup>a</sup> M. Schwarzwälder,<sup>c</sup> W.-C. Liao,<sup>c</sup> D. Gajan,<sup>d</sup> G. Bodenhausen,<sup>befg</sup> L. Emsley,<sup>b</sup> A. Lesage,<sup>d</sup> S. Jannin,<sup>b</sup> C. Copéret<sup>\*c</sup> and C. Thieuleux<sup>\*a</sup>

Hyperpolarization of metabolites by dissolution dynamic nuclear polarization (D-DNP) for MRI applications often requires fast and efficient removal of the radicals (polarizing agents). Ordered mesoporous SBA-15 silica materials containing homogeneously dispersed radicals, referred to as HYperPolarizing SOlids (HYPSOs), enable high polarization  $-P(^{1}H) = 50\%$  at 1.2 K - and straightforward separation of the polarizing HYPSO material from the hyperpolarized solution by filtration. However, the one-dimensional tubular pores of SBA-15 type materials are not ideal for nuclear spin diffusion, which may limit efficient polarization. Here, we develop a generation of hyperpolarizing solids based on a SBA-16 structure with a network of pores interconnected in three dimensions, which allows a significant increase of polarization, *i.e.*  $P(^{1}H) = 63\%$  at 1.2 K. This result illustrates how one can improve materials by combining a control of the incorporation of radicals with a better design of the porous network structures.

One of the major limitations of nuclear magnetic resonance spectroscopy (NMR) and magnetic resonance imaging (MRI) is their intrinsic low sensitivity, which arises from low thermal equilibrium (Boltzmann) polarization of the nuclear spins at room temperature, even with the highest field instruments. This low sensitivity is particularly challenging for low-gamma nuclear spins. For carbon-13, polarization is as low as  $P(^{13}C) \approx$  $2 \times 10^{-5}$  at T = 300 K and  $B_0 = 23.5$  T. This weakness can be circumvented by dissolution dynamic nuclear polarization (D-DNP).<sup>1,2</sup> D-DNP takes advantage of the high electron spin polarization at low temperatures to enhance the nuclear spin polarization well beyond thermodynamic equilibrium (>10 000 times) via microwave irradiation. Such huge gains in sensitivity allow metabolic imaging and have for example enabled the detection of anomalous metabolic rates in prostate tumors in living patients.<sup>3</sup>

<sup>a</sup>Université de Lyon, Institut de Chimie de Lyon, LC2P2, UMR 5265 CNRS-CPE Lyon-UCBL, CPE Lyon, 43 Bvd du 11 Novembre 1918, 69100 Villeurbanne, France. E-mail: david.baudouin@univ-lyon1.fr; chloe.thieuleux@univ-lyon1.fr

<sup>8</sup>UMR 7203, CNRS/UPMC/ENS, Paris, France

† Electronic supplementary information (ESI) available. See DOI: 10.1039/c6sc02055k

Hyperpolarization by D-DNP involves microwave irradiation at low temperatures in moderate magnetic fields (typically T = 1.2 K and  $B_0 = 3.35$  or 6.7 T) of frozen glassy solutions doped with stable free radicals and molecules of interest (e.g. metabolites or tracers). The preparation of a glassy frozen sample is important because it ensures that the radicals are statistically distributed, without the formation of ice crystals, which leads to optimal DNP efficiency. For that purpose, glass-forming agents such as glycerol, DMSO or ethanol are usually included in high concentrations. In a typical D-DNP experiment, following polarization, the polarized solution is rapidly brought to room temperature using superheated water and quickly transferred to the MRI or NMR machine for further studies. Once the solution is hyperpolarized, both radicals and glass-forming agents are unwanted and should obviously not be injected into patients. Furthermore, radicals act as paramagnetic relaxing agents, inducing faster depolarization.4 Therefore, rapid removal of radicals after polarization and before use is essential. In the case of trityl and BDPA radicals, the removal can be achieved by precipitation followed by filtration or by ion exchange.<sup>5,6</sup> However both methods are limited to specific sample formulations. Alternatively, nitroxide based radicals can be scavenged by ascorbate (vitamin C),<sup>7</sup> which attenuates paramagnetic relaxation but leads to contamination of the samples by hydroxylamines. Radicals can be incorporated into polymers such as polystyrene particles8 or hydrogels,9 allowing physical separation of the polarizing agent from the solution, but the efficiency for D-DNP system is limited and the filtration not straightforward.



View Article Online

<sup>&</sup>lt;sup>b</sup>Institut des Sciences et Ingénierie Chimiques, Ecole Polytechnique Fédérale de Lausanne (EPFL), CH-1015 Lausanne, Switzerland

<sup>&</sup>lt;sup>c</sup>ETH Zürich, Department of Chemistry and Applied Biosciences, Vladimir-Prelog-Weg 1-5/10, 8093 Zürich, Switzerland. E-mail: ccoperet@inorg.chem.ethz.ch

<sup>&</sup>lt;sup>d</sup>Université de Lyon, Institut des Sciences Analytiques, UMR 5280, CNRS, Université Lyon 1, ENS Lyon 5 rue de la Doua, F-69100 Villeurbanne, France

<sup>&</sup>lt;sup>e</sup>Département de Chimie, Ecole Normale Supérieure, 24 Rue Lhomond, 75231 Paris Cedex 05, France

<sup>&</sup>lt;sup>1</sup>Université Pierre-et-Marie Curie, Paris, France

In this context, we have recently developed a family of solid polarizing matrices based on hybrid materials containing covalently bound radicals, coined HYperPolarizing SOlids (HYPSO). These materials provide in principle a universal solution to the above-mentioned issues: *i.e.* fast and easy removal of radicals by filtration, and the absence of glassforming agents. HYPSO are porous and robust silica-based solids on which any radicals (i.e. TEMPO, trityl...) can be covalently and homogeneously attached to the surface of their pores.<sup>10,11</sup> We showed an efficient direct polarization approaching  $P(^{13}C) = 15\%$  with a build-up time of 2 h in a 3 M 1-13C sodium pyruvate aqueous solution impregnated with trityl-based HYPSO.11 The first generation TEMPO-based materials, HYPSO-1, allowed reaching a <sup>13</sup>C polarization as high as  $P(^{13}C) = 33\%$  in only 20 min using a state-of-the-art polarizer including microwave frequency modulation<sup>12</sup> and <sup>1</sup>H-<sup>13</sup>C crosspolarization.<sup>13</sup> Despite significant research efforts to improve the DNP performances of the first generations of HYPSOs, it was not possible to enhance the proton polarization beyond  $P(^{1}H) =$ 50%, well below the  $P(^{1}H) = 90\%$  that can be obtained under similar conditions in glassy water/glycerol TEMPOL solutions.13

The two first generations of HYPSOs were based on ordered mesoporous SBA-15 type structures, with a skeleton consisting of 8–10 nm diameter 1D-pore channels stacked in a 2D-hexagonal arrangement (Fig. 1 – top right). Two generations of nitroxide-based materials, HYPSO-1 and HYPSO-2, were hence prepared and differ from the linker used to anchor the radical to the solid surface, a propylamido<sup>10</sup> and a 1,2,3-triazole-propyl tethers, respectively.<sup>11</sup> Using a direct synthesis, the radicals were homogeneously incorporated onto the pore surface of HYPSO-1 and HYPSO-2 by peptide coupling or click chemistry, respectively, avoiding radical aggregation, which is important for

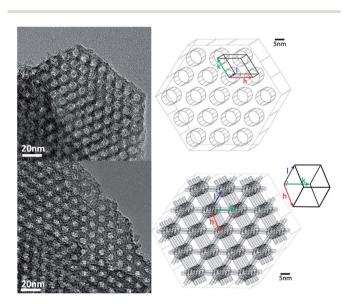


Fig. 1 Left: TEM pictures of hexagonal  $1/100_N_3$ \_SBA-15 in the [001] axis (top) and of cubic  $1/140_N_3$ \_SBA-16 in the [111] axis (bottom). Right: Schematic representations of the 2D pore structure in the [001] axis of SBA-15/HYPSO-2 (top) and the 3D pore structure in the [111] axis of SBA-16/HYPSO-3 (axis). The radicals are distributed uniformly over the surface of the pores. See Fig. SI-1† for SBA-16 in [100] axis.

D-DNP. However, in such a structure, the pores do not communicate with each other; we thus hypothesized that this could be a limiting factor for both nuclear spin diffusion and for the three-dimensional distribution of the radicals, in comparison to frozen glassy solutions. We therefore reasoned that a silica architecture with a 3D cubic porous network (for example using SBA-16 like structures) could improve the DNP performance. Here we describe the development of materials with cubic network arrangement (HYPSO-3) and show that they lead to greater polarization with respect to the one-dimensional porous HYPSO-1/2 materials, yielding proton polarization up to  $P(^{1}H) = 63\%$ . We show how the 3D cubic material can also be efficiently used under Magic Angle Spinning (MAS) DNP conditions.

### Results and discussion

#### Preparation and characterization of materials

First, propylazido functionalized SBA-16 materials were prepared by co-condensation of (3-azidopropyl)triethoxylsilane and tetraethyl orthosilicate in an HCl/NaCl aqueous medium in the presence of pluronic F127 as a structure-directing agent (SDA).14 The ratios between (3-azidopropyl)triethoxylsilane and tetraethyl orthosilicate varied in the range 1/34, 1/60, 1/100, 1/140, 1/320, corresponding to a loading of organic functionalities of 472, 272, 164, 118 and 52  $\mu$ mol<sub> $\equiv$ SiR</sub> g<sup>-1</sup>. The materials obtained were filtered, and the SDA was extracted with a pyridine/HCl solution for 24 h and finally washed and dried under 10<sup>-5</sup> mbar vacuum at 135 °C.<sup>15</sup> The different materials were analyzed by N<sub>2</sub> adsorption, which showed a type IV isotherm characteristic of mesoporous materials. The BET model indicated surface areas ranging from *ca.* 910 to 1080 m<sup>2</sup> g<sup>-1</sup> (see Table 1) with total pore volumes ( $P/P_0 < 0.99$ ) of 0.62–0.82 cm<sup>3</sup>  $g^{-1}$ , of which 37–46% are micropores according to the  $\alpha$ -plot model. The pore sizes have narrow bimodal distributions at 1.4-1.7 (inter-connecting windows, MP model) and 6.2-7.0 nm (cavities, BJH model). The sharp drop in the desorption branch of the isotherm observed at *ca.* 0.45  $P/P_0$  is typical of nitrogen desorption from mesoporous materials through smaller pores (ink bottle effect), in line with the formation of 6-7 nm cages interconnected by 1-2 nm connecting micropores. Small Angle X-Ray Diffraction (SA-XRD, see Table SI-1†) of  $1/34_N_3$ \_SBA-16 presented a strong peak at *ca.*  $2\theta = 0.75^\circ$ assigned to the X-ray diffraction of the {110} family of planes, indicative of well-structured cubic-centered body SBA-16 materials ( $Im\bar{3}m$  space group). SA-XRD and N<sub>2</sub> adsorption/desorption data allow one to calculate a mean micropore length of ca. 10 nm (Table 1 and SI-1<sup>†</sup>). Transmission Electron Micrographs (TEM) of 1/140\_N<sub>3</sub>\_SBA-16 confirm the formation of a cubic ordered mesoporous material (see representative pictures in Fig. 1 and S1<sup>†</sup>).

For comparison, several polarizing matrices with 2D hexagonal arrangements of mesopore tube-like pores (SBA-15 type materials), hereafter named HYPSO-2, were also prepared.<sup>10</sup> These materials are highly porous with a BET surface area, a total pore volume, BJH and MP pore diameters of 770–870 m<sup>2</sup>  $g^{-1}$ , 1.1–1.2 cm<sup>3</sup>  $g^{-1}$  and 8.0–9.2 nm, respectively (Table S2†).

Table 1	Organic function loading	is and textural characteristics of the materials obtained	ed from SA-XRD and $N_2$ -adsorption desorption at 77 K

Materials	$[\equiv SiR]/\mu mol_{\equiv SiR} g^{-1}$	$S_{\rm BET}/{ m m}^2~{ m g}^{-1}$	$V_{\rm p}^{\ a}$ (tot.)/m <sup>2</sup> g <sup>-1</sup>	$V_{\rm p}^{\ b}(\mu)/{\rm m}^2~{\rm g}^{-1}$	<i>D</i> <sub>p</sub> <sup>c</sup> /nm	$L_{\mu pore}^{d}/nm$
1/34_N <sub>3</sub> _SBA-16	472	1012	0.68	0.31	1.7/6.2	9.6
1/34_HYPSO-3	472	729	0.50	0.19	1.6/5.4	10.4
1/60_N <sub>3</sub> _SBA-16	272	1010	0.66	0.29	1.6/6.3	n.d.
1/60_HYPSO-3	272	752	0.52	0.22	1.6/5.4	10.2
1/100_N <sub>3</sub> _SBA-16	164	913	0.62	0.11	1.7/6.2	n.d.
1/100_HYPSO-3	164	893	0.63	0.26	1.3/6.3	8.6
1/140_N <sub>3</sub> _SBA-16	118	1184	0.82	0.33	1.4/7.0	n.d.
1/140_HYPSO-3	118	983	0.69	0.26	1.7/7.1	9.3
1/320_N3_SBA-16	52	1068	0.75	0.48	1.7/7.0	n.d.
1/320_HYPSO-3	52	714	0.48	0.27	1.6/5.4	12.3

<sup>*a*</sup> Total pore volume corresponding to the quantity of N<sub>2</sub> adsorbed at  $P/P_0 = 0.99$ . <sup>*b*</sup> Micropore volume, calculated from the  $\alpha_s$  plot model. <sup>*c*</sup> Micropore mean diameter calculated using MP model/mesopore mean diameter calculated using the BJH model (adsorption branch). <sup>*d*</sup> Micropore mean length, calculated using  $L_{\mu pore} = (d_{(110)}/\cos(\pi/4) - D_{meso})$  using the mesoporous diameter  $D_p$  and the *d*-spacing d(110) obtained from Small Angle XRD analysis.

Importantly, these SBA-15 type materials exhibit non-interconnected 1.8 nm micropores (5–8% of pore volume). In contrast, HYPSO-3 (SBA-16 materials) exhibits spherical mesopores interconnected by micro-channels in all three dimensions. This pore structure leads to different textures: a lower pore volume and a different intra-grain pore volume distribution (Fig. 1).

Post-functionalization of N<sub>3</sub>\_SBA-16 to obtain HYPSO-3 materials was performed using copper-catalyzed azide-alkyne cycloaddition (Cu-AAC)<sup>16</sup> in the presence of *o*-propargyl TEMPO, CuI, dry DMF and Et<sub>3</sub>N (see details in ESI and Fig. S2 and S3†). Diffuse Reflectance Infrared Fourier Transform Spectroscopy (DRIFTS) analysis of the powder allowed evaluation of the efficiency of the cycloaddition (referred to as Cu-AAC yield) on HYPSO-3.

As shown in Table 2, 88–64% of the starting azido  $-N_3$  reacted with the *o*-propargyl TEMPO reactant, the yield decreasing slowly when decreasing the molar concentration of radicals (quoted 1/*xx* ratio that stands for 1 mol of radical per *xx* mol of SiO<sub>2</sub>). The concentration of radical incorporated in HYPSO-3 was quantified by recording X-band CW EPR spectra at room temperature. Nitroxyl radical loadings of 246, 135, 79, 50 and 33 µmol<sub>NO</sub> g<sup>-1</sup> were measured for ratios of 1/34, 1/60, 1/100, 1/140 and 1/320 respectively, corresponding to radical concentrations of 491, 260, 125, 72 and 67 µmol<sub>NO</sub> cm<sup>-3</sup> within the

Table 2         Characteristics of HYPSO-3 materials									
Ratio	$[R]/$ $\mu mol g^{-1}$	[NO']/ $\mu$ mol g <sup>-1</sup>	Cu- <sub>AAC</sub> yield <sup>a</sup> (%)	EPR yield <sup>b</sup> (%)	[NO <sup>•</sup> ] <sup>c</sup> / μmol cm <sup>-3</sup>				
1/24	472	0.4.6	0.0	50	401				
1/34	472	246	88	52	491				
1/60	272	135	81	50	260				
1/100	164	79	77	48	125				
1/140	118	50	74	41	72				
1/320	52	33	64	63	67				

<sup>*a*</sup> Percentage of N<sub>3</sub> reacted after Cu-AAC (obtained by DRIFT). <sup>*b*</sup> Percentage of NO' compared to initial N<sub>3</sub>. <sup>*c*</sup> Concentration per total pore volume ( $P/P_0 = 0.99$ ). total volume of the pores. The radical concentrations show that the yields of post-functionalization are in the range 41–63%, similar to 2D-hexagonal materials characterized by the same method (42–57%). No significant difference between the EPR profiles of HYPSO-2 and HYPSO-3 could be observed at room temperature (Fig. S4†).

#### EPR spectroscopy

The average inter-radical distances in HYPSO-3 were evaluated from the line width of the central EPR signal at 110 K, which is known to be broadened by electron–electron dipolar couplings and by spin exchange.<sup>17</sup> Such an analysis is limited to an average inter-radical distance  $r_{\rm RR} < 2$  nm. For larger distances, the dipolar line width is masked by inhomogeneous broadening (≈12 G). The EPR line-widths measured are 20.9, 15.3, 13.1 and 12.1 G for HYPSO-3 with ratios N<sub>3</sub>/TEOS = 1/34, 1/60, 1/100 and 1/140 respectively (Fig. 2 and Table S3†). The dipolar broadening is almost proportional to the radical concentrations for

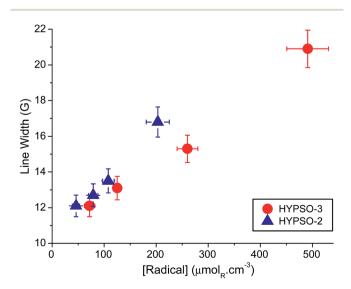


Fig. 2 EPR linewidths of HYPSO-2 and -3 as a function of the molar radical concentration (in  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>).

HYPSO-2 and -3, although the slopes are different. Indeed, the line width of HYPSO-3 (SBA-16) is narrower than that of HYPSO-2 (SBA-15, data from ref. 11) for the same volumetric radical concentration. This narrower line width can be attributed to a more uniform 3D distribution of the radicals. We believe that this improvement likely originates from the difference in pore shape (tubular  $\nu s$ . interconnected cages) and structuration between HYPSO-2 and HYPSO-3.

#### **DNP** performances

The DNP efficiency of HYPSO-3 was first evaluated for samples spinning at the magic angle (MAS) near 100 K and gave results similar to HYPSO-2<sup>10,18,19</sup> (see ESI<sup>†</sup>). The DNP performance of HYPSO-2 and -3 was then determined at 4.2 and 1.2 K. Both materials were impregnated by filling *ca.* 95% of the pore volume with  $D_2O : H_2O(8 : 2)$ . The results obtained at 4.2 K with microwave frequency modulation<sup>12</sup> are presented in Fig. 3.

One can observe a maximum polarization in the vicinity of  $[R] = 50 \ \mu\text{mol cm}^{-3}$ . At this concentration, HYPSO-3 yields a polarization  $P(^{1}\text{H}) = 12.5\%$ , significantly higher than the 7.5% obtained with HYPSO-2. The use of microwave frequency

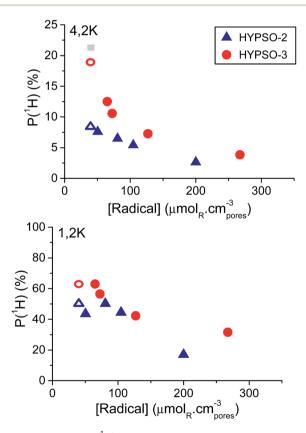


Fig. 3 Polarization  $P({}^{1}H)$  measured using microwave frequency modulation at 4.2 K (top) and 1.2 K (bottom) and 6.7 T as a function of radical concentration (in  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>) for: HYPSO-2 (**A**) and HYPSO-3 (**●**) impregnated with  $D_2O: H_2O$  (8 : 2). The open circles **O** and triangles  $\triangle$  correspond to a solution of 40 mM TEMPOL in  $H_2O: D_2-O:$  glycerol-d<sub>8</sub> (10 : 40 : 50) impregnated in 1/140\_N<sub>3</sub>\_SBA-16 (radical-free HYPSO-3) and 1/140\_N<sub>3</sub>\_SBA-15 (radical-free HYPSO-2), respectively. The grey square **I** corresponds to a TEMPOL reference DNP solution, without HYPSO.

modulation improves the DNP performances of both HYPSO-2 and -3 but only for  $[R] < 100 \ \mu mol \ cm^{-3}$  (see Fig. S6†).

For comparison, an isotropic "glassy"  $H_2O : D_2O : glycerol-d_8$ (10 : 40 : 50) matrix without HYPSO doped with 40 mM TEM-POL (40 µmol<sub>NO</sub> cm<sup>-3</sup>) gave rise to  $P(^{1}H) = 21.5\%$ . Note that a 80 µmol<sub>NO</sub> cm<sup>-3</sup> solution gives comparable results. When impregnating HYPSO-3 type matrices containing surface azidogroups instead of TEMPO units (1/140\_N3\_SBA-16) with a 40 mM TEMPO solution, a polarization  $P(^{1}H) = 19\%$  was obtained. This polarization value is close to that of the isotropic glassy DNP solution ( $P(^{1}H) = 21.5\%$ ). On the contrary, when impregnated in a HYPSO-2 matrix without any radicals (1/140\_N<sub>3</sub>\_SBA-15), the <sup>1</sup>H polarization dropped to  $P(^{1}H) = 9\%$ . We take this as a strong indication that the cubic 3D porous network is advantageous for efficient DNP as compared to a onedimensional network.

The build-up time constant (Fig. S6<sup>†</sup>) was found to be  $\tau_{\text{DNP}} =$ 74 s for the reference DNP solution without HYPSO at 4.2 K. The same solution impregnated in 1/140\_N<sub>3</sub>\_SBA-16 (same as HYPSO-3) gave  $\tau_{\text{DNP}} =$  145 s for a similar polarization, and  $\tau_{\text{DNP}} =$ 77 s when impregnated in 1/140\_N<sub>3</sub>\_SBA-15 (same as HYPSO 2) (Fig. S6<sup>†</sup>). EPR studies confirmed that this lengthening of the DNP build-up in HYPSO-3 was not due to a radical quenching effect.

At 1.2 K, HYPSO-3 yielded a polarization  $P({}^{1}\text{H}) > 40\%$  over a broad range of radical concentrations  $50 < [\text{R}] < 160 \ \mu\text{mol}_{\text{NO}}$ cm<sup>-3</sup>, reaching  $P({}^{1}\text{H}) = 63\%$  for  $[\text{R}] = 67 \ \mu\text{mol}_{\text{NO}}$  cm<sup>-3</sup> (*cf.* Fig. 3, bottom).

By comparison, HYPSO-2 only yielded a maximum polarization  $P({}^{1}\text{H}) = 50\%$  for an optimal [R] = 79  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>. The polarization in the H<sub>2</sub>O : D<sub>2</sub>O : glycerol-d<sub>8</sub> (10 : 40 : 50) mixture containing 40 mM TEMPOL can reach  $P({}^{1}\text{H}) = 90\%$  under the same conditions.<sup>13</sup> At 1.2 K, frequency modulation was found to have a positive effect on DNP for HYPSO-2 and -3 when [R] < 75  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>. When the optimal DNP solution was impregnated in 1/140\_N<sub>3</sub>\_SBA-16 (radical-free HYPSO-3) and 1/140\_N<sub>3</sub>\_SBA-15 (radical-free HYPSO-2), we observed  $P({}^{1}\text{H}) = 63\%$  for HYPSO-3 and 50% for HYPSO-2 at 1.2 K. Note that higher polarization might be reached using HYPSO-3 with lower radical concentration, but build up times ( $\tau_{\text{DNP}}$ ) would become

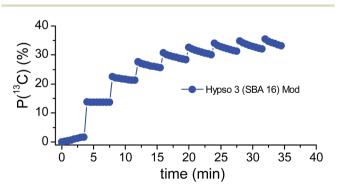


Fig. 4 Left <sup>1</sup>H  $\rightarrow$  <sup>13</sup>C CP-DNP performed on HYPSO-3 material (67  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>) impregnated with a 3 M solution of [1-<sup>13</sup>C]-acetate in D<sub>2</sub>O. *P*(<sup>13</sup>C) 36% is reached in 32.5 min with <sup>1</sup>H  $\rightarrow$  <sup>13</sup>C CP applied at 4 min intervals.

very long (>300 s) (see Fig. S6†). Finally, we prepared a 3 M solution of sodium [1-<sup>13</sup>C]-acetate in H<sub>2</sub>O: D<sub>2</sub>O (1:9) to impregnate HYPSO-3 (with [R] = 67  $\mu$ mol<sub>NO</sub> cm<sup>-3</sup>) and we obtained *P*(<sup>1</sup>H)  $\approx$  50%. Cross polarization<sup>20</sup> was performed with 8 contacts every 4 min, yielding *P*(<sup>1</sup>H  $\rightarrow$  <sup>13</sup>C) = 36% after *ca.* 30 minutes (Fig. 4). After dissolution, the aqueous acetate solution was recovered by filtration and centrifugation and subjected to ESR analysis which confirmed the presence of a negligible quantity of radicals in the liquid (*ca.* 1  $\mu$ mol L<sup>-1</sup> *i.e.* <0.3% of HYPSO) probably arising from the presence of very small grains of HYPSO-3 in the solution.

## Conclusions

In conclusion, we have shown that the architecture of the porous network affects the hyperpolarization properties of HYPSO materials for DNP. Here, a material with 1D tubular pores (HYPSO-2) allows one to reach  $P(^{1}H) = 50\%$  at 1.2 K, while a material with 3D interconnected cage-like pores in cubic symmetry (HYPSO-3) leads to a polarization of  $P(^{1}H) = 63\%$ . We propose that this is due to (i) a more homogeneous 3D distribution of nitroxyl radicals in HYPSO-3, as demonstrated by EPR, and (ii) the interconnection of the pores, which allows for nuclear spin diffusion in all three directions, as in ideal solutions. We are currently working on designing optimal materials based on these underlying principles.

## Acknowledgements

This work was supported by the CNRS, the SATT Lyon (Pulsalys), the European Research Council (ERC) Advanced Grant No. 320860, the Swiss National Science Foundation (SNSF), the Ecole Polytechnique Fédérale de Lausanne (EPFL), ETH Zürich, and the Swiss Commission for Technology & Innovation (CTI). Mathieu Baudin is kindly acknowledged for his contribution to dissolution DNP experiments.

## References

- J. H. Ardenkjær-Larsen, B. Fridlund, A. Gram, G. Hansson, L. Hansson, M. H. Lerche, R. Servin, M. Thaning and K. Golman, *Proc. Natl. Acad. Sci. U. S. A.*, 2003, **100**, 10158– 10163.
- 2 J. H. Ardenkjaer-Larsen, J. Magn. Reson., 2016, 264, 3-12.
- 3 S. J. Nelson, J. Kurhanewicz, D. B. Vigneron, P. E. Z. Larson,
  A. L. Harzstark, M. Ferrone, M. van Criekinge, J. W. Chang,
  R. Bok, I. Park, G. Reed, L. Carvajal, E. J. Small,
  P. Munster, V. K. Weinberg, J. H. Ardenkjaer-Larsen,
  A. P. Chen, R. E. Hurd, L.-I. Odegardstuen, F. J. Robb,
  J. Tropp and J. A. Murray, *Sci. Transl. Med.*, 2013, 5, 198ra108.
- 4 P. Miéville, S. Jannin and G. Bodenhausen, *J. Magn. Reson.*, 2011, **210**, 137–140.

- J. A. Urbahn, J. Ardenkjaer-Larsen, A. M. Leach, E. Telfeyan,
   D. K. Dietrich, D. B. Whitt, P. Miller and E. W. Stautner, US Pat., US 2008/0240998 A1, 2007.
- 6 J. H. Ardenkjaer-Larsen, A. M. Leach, N. Clarke, J. Urbahn, D. Anderson and T. W. Skloss, *NMR Biomed.*, 2011, 24, 927–932.
- 7 P. Miéville, P. Ahuja, R. Sarkar, S. Jannin, P. R. Vasos,
  S. Gerber-Lemaire, M. Mishkovsky, A. Comment,
  R. Gruetter, O. Ouari, P. Tordo and G. Bodenhausen,
  Angew. Chem., Int. Ed., 2010, 49, 6182–6185.
- 8 Y. Zhang, P. J. Baker and L. B. Casabianca, *J. Phys. Chem. B*, 2016, **120**, 18–24.
- 9 B. C. Dollmann, M. J. N. Junk, M. Drechsler, H. W. Spiess, D. Hinderberger and K. Munnemann, *Phys. Chem. Chem. Phys.*, 2010, **12**, 5879–5882; T. Cheng, M. Mishkovsky, M. J. N. Junk, K. Münnemann and A. Comment, *Macromol. Rapid Commun.*, 2016, **37**, 1074–1078.
- D. Gajan, M. Schwarzwälder, M. P. Conley, W. R. Grüning, A. J. Rossini, A. Zagdoun, M. Lelli, M. Yulikov, G. Jeschke, C. Sauvée, O. Ouari, P. Tordo, L. Veyre, A. Lesage, C. Thieuleux, L. Emsley and C. Copéret, *J. Am. Chem. Soc.*, 2013, 135, 15459–15466.
- 11 D. Gajan, A. Bornet, B. Vuichoud, J. Milani, R. Melzi, H. A. van Kalkeren, L. Veyre, C. Thieuleux, M. P. Conley, W. R. Gruning, M. Schwarzwalder, A. Lesage, C. Coperet, G. Bodenhausen, L. Emsley and S. Jannin, *Proc. Natl. Acad. Sci. U. S. A.*, 2014, **111**, 14693–14697.
- 12 A. Bornet, J. Milani, B. Vuichoud, A. J. Perez Linde, G. Bodenhausen and S. Jannin, *Chem. Phys. Lett.*, 2014, 602, 63–67.
- 13 S. Jannin, A. Bornet, R. Melzi and G. Bodenhausen, *Chem. Phys. Lett.*, 2012, **549**, 99–102.
- 14 A. Boullanger, J. Alauzun, A. Mehdi, C. Reye and R. J. P. Corriu, *New J. Chem.*, 2010, 34, 738–743.
- 15 R. J. P. Corriu, A. Mehdi, C. Reye and C. Thieuleux, *New J. Chem.*, 2003, **27**, 905–908.
- 16 H. C. Kolb, M. G. Finn and K. B. Sharpless, *Angew. Chem., Int. Ed.*, 2001, **40**, 2004.
- 17 G. Jeschke, A. Koch, U. Jonas and A. Godt, *J. Magn. Reson.*, 2002, **155**, 72–82.
- 18 A. J. Rossini, A. Zagdoun, M. Lelli, A. Lesage, C. Coperet and L. Emsley, *Acc. Chem. Res.*, 2013, **46**, 1942–1951.
- 19 A. Lesage, M. Lelli, D. Gajan, M. A. Caporini, V. Vitzthum,
  P. Miéville, J. Alauzun, A. Roussey, C. Thieuleux, A. Mehdi,
  G. Bodenhausen, C. Coperet and L. Emsley, *J. Am. Chem.* Soc., 2010, 132, 15459–15461.
- 20 A. Bornet, R. Melzi, A. J. P. Linde, P. Hautle, B. van den Brandt, S. Jannin and G. Bodenhausen, *J. Phys. Chem. Lett.*, 2013, 4, 111–114.