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1                   **UPSCALING ELECTROKINETIC TRANSPORT IN CLAYS**  
2                   **WITH LATTICE BOLTZMANN AND PORE NETWORK MODELS**

3  
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9  
10 **SHORT TITLE:** Upscaling electrokinetic transport in clays

11  
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14  
15 **ABSTRACT**

16 This contribution presents a method for the numerical determination of the steady-state response  
17 of complex charged porous media to pressure, salt concentration and electric potential gradients.  
18 The Pore Network Model (PNM), describing the porosity as a network of pores connected by  
19 channels, is extended to capture electrokinetic couplings which arise at charged solid-liquid  
20 interfaces. This allows us to compute the macroscopic fluxes of solvent, salt and charge across a  
21 numerical sample submitted to macroscopic gradients. On the channel scale, the microscopic  
22 transport coefficients are obtained by solving analytically (in simple cases) or numerically the  
23 Poisson-Nernst-Planck and Stokes equations. The PNM approach then allows us to upscale these  
24 transport properties to the sample scale, accounting for the complex pore structure of the material  
25 via the distribution of channel diameters. The Onsager relations between macroscopic transport  
26 coefficients are preserved, as expected. However, electrokinetic couplings combined with the  
27 sample heterogeneity result for some macroscopic transport coefficients (e.g. permeability or  
28 electro-osmotic coefficient) in qualitative differences with respect to their microscopic  
29 counterparts. This underlines the care that should be taken when accounting for transport  
30 properties based on a single channel of average diameter.

31  
32 **KEYWORDS:** Electrokinetics, Coupled transport, Homogenization, Lattice Boltzmann, Pore  
33 Network Model, Upscaling

34 **INTRODUCTION**

35

36 Electrokinetic effects refer to the dynamic coupling between the solvent and charge flows which  
37 occur at a charged interface. The presence of surface charge in a porous medium has important  
38 practical applications in membrane technology (e.g., ion exchange and water desalination) and in  
39 environmental science, since most rocks and soils contain minerals (such as clays) that bear a  
40 permanent surface charge. As an example, electro-osmosis generates a solvent flow under an  
41 applied electric field, due to the driving of the electrically charged fluid in the vicinity of charged  
42 surfaces. Conversely, a pressure gradient induces the flow of a charged fluid, hence, an electric  
43 current. In geophysics, the electroseismic effect, by which an electro-magnetic wave is generated  
44 from the motion of underground fluids under an applied acoustic wave, is exploited to determine  
45 the properties of geological formations (Thompson, 1936; Pride and Haartsen, 1996; Mizutani *et*  
46 *al.*, 1976). Streaming potentials and electro-osmotic flows can be measured in the laboratory to  
47 characterize the properties of porous media (Luong and Sprik, 2013).

48

49 The modeling and simulation of electrokinetic effects in porous media, and, more generally, of  
50 all coupled transport phenomena, including the osmotic solvent flow due to a salt concentration  
51 gradient, thus have been the subject of a large number of investigations, both on the pore scale  
52 where the couplings originate and on the sample scale corresponding to the experimental  
53 measurements. From the mathematical point of view, this upscaling can be performed rigorously  
54 using the homogenization approach. This provides expressions of the macroscopic transport  
55 coefficients as solutions of coupled partial differential equations on the pore scale, which then  
56 have to be solved using simplifying assumptions or numerically. Some general results, such as  
57 Onsager's relations for the macroscopic transport coefficients, can be demonstrated without even  
58 resorting to the numerical resolution of the mathematical problem (Moyne and Murad, 2006a&b;  
59 Allaire *et al.*, 2010&2014).

60

61 For practical applications, most studies of electrokinetic couplings rely on an oversimplified  
62 idealization of the geometry, with single slit pores or cylinders with dimensions or surface  
63 charge densities estimated from the macroscopic properties of the real system (Bresler, 1973;  
64 Gonçalves *et al.*, 2012). However, the heterogeneity of the material, combined with the

65 electrokinetic couplings, may influence the overall behavior on the sample scale, so that such  
66 idealizations may not reflect the actual response of the medium. Direct numerical resolution of  
67 the coupled Poisson-Nernst-Planck (PNP) and Navier-Stokes (NS) equations in various complex  
68 systems (random packings, reconstructed and fractured porous media) has also been proposed by  
69 Adler and co-workers. Such an approach is usually difficult to implement for macroscopic  
70 samples, due to the lack of experimental data on the fine structure of the material over large  
71 distances (Coelho *et al.*, 1996; Marino *et al.*, 2001; Gupta *et al.*, 2006). The systematic study of a  
72 representative number of samples is also prevented by the computational cost of direct numerical  
73 simulation.

74

75 In the case of clays, an additional difficulty arises due to the complex multiscale porosity of the  
76 material and the lack of experimental data on the intermediate scales, which is at the heart of this  
77 workshop. In the present contribution, we present a numerical homogenization scheme leading to  
78 a description of transport through macroscopic charged porous materials at low computational  
79 cost, thereby enabling the systematic study of the combined effects of electrokinetic couplings  
80 and sample heterogeneity. The algorithm to upscale the electrokinetic couplings is based on the  
81 Pore Network Model (PNM), which relies on the one hand on a simplified description of the  
82 electrokinetic transport on the pore scale and on the other hand on a statistical distribution of the  
83 geometry of the pores. This allows to investigate how the upscaled electrokinetic properties  
84 depend on the heterogeneity of the sample, in addition to the surface charge density and the salt  
85 concentration.

86

87

## 88 **ELECTROKINETICS ON THE SAMPLE SCALE**

89

90 On the macroscopic scale of a clay sample, pressure  $P$ , electric potential  $V$  and salt concentration  
91 gradients (or, equivalently, solvent, cation and anion chemical potentials gradients), induce  
92 macroscopic fluxes of mass, electric charge and salt (or, equivalently, solvent, cation and anion  
93 chemical fluxes). For sufficiently small applied gradients, the response is linear and the fluxes  
94 can be expressed as a function of the applied gradients via a coupling matrix:

95

$$\begin{pmatrix} Q_0 \\ Q_1 \\ Q_2 \end{pmatrix} = -\frac{S}{\eta} \begin{pmatrix} K_0^P & K_0^C & K_0^V \\ K_1^P & K_1^C & K_1^V \\ K_2^P & K_2^C & K_2^V \end{pmatrix} \begin{pmatrix} \nabla P \\ \nabla C \\ \nabla V \end{pmatrix}, \quad (1)$$

96 where the subscripts 0, 1 and 2 refer to solvent, cations and anions, respectively, Q indicate their  
 97 fluxes, C is the logarithm of the salt concentration,  $\eta$  the solvent viscosity and S the cross-section  
 98 of the sample. Along the diagonal one finds for example the sample permeability  $K_0^P$  and the  
 99 electric conductivity  $K_2^V$ . The fundamental question is then: How do these coefficients emerge  
 100 from the microscopic structure of the material, including heterogeneities on intermediate scales,  
 101 and from the surface charge density of the solid matrix?  
 102  
 103  
 104

105

## 106 **ELECTROKINETICS ON THE PORE SCALE**

107

108 Recently, significant progress has been made on the derivation of the macroscopic transport  
 109 equations from the pore-scale ones. These studies usually start from a continuous description of  
 110 the fluid via transport equations, which are then upscaled to derive their average effect on the  
 111 sample scale, which is quantified by a coupling matrix relating the solvent and ionic fluxes to the  
 112 corresponding forces (pressure, potential, and concentration gradients). The solvent flow under  
 113 applied local forces is accounted for via the NS equation (or even the Stokes equation), which  
 114 includes a local force due to electrochemical potential gradients. The solute fluxes are due one  
 115 the one hand to the advection by the fluid and on the other hand to the local electrochemical  
 116 potential gradients; they can be modelled on this scale using the PNP equations. The limitations  
 117 of such continuous descriptions to describe solvent and ion transport in clay nanopores, which  
 118 can be assessed using molecular simulations (Botan *et al.*, 2010&2013), will not be discussed  
 119 here. Rather, the present discussion focusses on how to upscale this to the macroscopic scale,  
 120 since the structure is too complex for a direct resolution on the whole sample.  
 121

122

123 The coupled Navier-Stokes and Poisson-Nernst-Planck equations can be solved numerically  
 124 using finite element or volume methods. For example, Adler and co-workers used this direct  
 125 numerical resolution in various complex systems (random packings, reconstructed and fractured  
 126 porous media) (Coelho *et al.*, 1996; Marino *et al.*, 2001), demonstrating in particular a universal  
 electrokinetic behaviour if appropriate rescaled quantities are introduced (Gupta *et al.*, 2006,

127 2008). Recently, alternative methods have been proposed to simulate electrokinetic effects  
128 starting from a more fundamental description of the fluid than the PNP and NS equation  
129 (Pagonabarraga *et al.*, 2010). For example, Capuani *et al.* proposed a hybrid lattice based  
130 approach (Lattice Boltzmann Electrokinetics, LBE) to capture the coupling of hydrodynamic  
131 flow with ion transport and the simulation of electrokinetic effects in colloidal suspensions  
132 (Capuani *et al.*, 2004; Pagonabarraga *et al.*, 2005). Such Lattice Boltzmann simulations have  
133 already been applied, without accounting for electrokinetic effects, to realistic rock geometries  
134 (Boek and Venturoli, 2010). In the context of the present numerical homogenization, LBE was  
135 recently used in a simple cylindrical geometry, in order to assess the validity range of the  
136 analytical solution of the linearized problem (Obliger *et al.*, 2013). This simpler analytical  
137 solution is then used in the PNM, even though in principle a numerical expression for the  
138 transport coefficient on the pore scale may also be used.

139

140

## 141 **NUMERICAL HOMOGENIZATION VIA A PORE NETWORK MODEL**

142

143 In order to investigate electrokinetic couplings on larger scales, including the effect of the  
144 heterogeneity of the material, we have recently proposed a simplified description based on the  
145 Pore Network Model (PNM). Such a model, originally developed by Fatt (1956) to predict  
146 multiphase flow properties in porous media, describes the porosity as a network of pores  
147 connected by channels. It has been extensively used and extended by petrophysicists in various  
148 situations, such as capillarity and multiphase flow through porous media (Békri *et al.*, 2005;  
149 Blunt, 2001; van Dijke and Sorbie, 2002), or mineral dissolution and precipitation in the context  
150 of CO<sub>2</sub> sequestration (Algive *et al.*, 2010).

151

152 In a nutshell, the PNM approach amounts to solving a set of conservation equations on the nodes  
153 of the network (in analogy with Kirchhoff's law for a network of resistors), on the basis of local  
154 fluxes through the channels connecting the nodes, under the effect of an external, macroscopic  
155 gradient. For electrokinetics, the pressure, salt concentration and electrical potential are  
156 introduced as pore variables on the nodes of the network. The fluxes through each link between  
157 nodes are determined locally using the transport matrix for a cylindrical channel, as determined

158 in the previous section as a function of the channel diameter, the surface charge density of the  
159 solid and the salt concentration inside the channel. The latter is determined via the Donnan  
160 equilibrium with a fictitious reservoir corresponding to the properties of the pores at both ends of  
161 the channel (Obliger *et al.*, 2014). Therefore, the macroscopic problem to be solved numerically  
162 has a non-linear structure, contrary to most previous applications of the PNM approach. This can  
163 be achieved numerically using a non-linear Newton solver.

164

165

## 166 **TRANSPORT COEFFICIENTS ON THE SAMPLE SCALE**

167

168 In addition to the transport coefficients on the channel scale, the crucial ingredient of the PNM is  
169 the distribution of pore/channel sizes and their spatial arrangement describing in a very  
170 simplified manner the complex structure of the porous network. In order to demonstrate the  
171 feasibility of the approach and to investigate systematically the effect of heterogeneity, a model  
172 distribution (of the Weibull type) was considered first. However, it is also possible to introduce a  
173 distribution deduced from experimental data, if a reliable one can be provided.

174

175 For a given pore/channel diameter distribution, a sufficient number of networks must be  
176 generated. For each of numerical sample, the macroscopic coefficients are determined by solving  
177 the conservation equations in the presence of applied gradients and by computing the  
178 macroscopic steady-state flux through the sample. In practice, three calculations must be done  
179 (one for each applied gradient) for which the three fluxes (mass, charge, salt concentration) are  
180 computed. This provides the nine macroscopic coefficients, which must then be averaged over  
181 the networks corresponding to the same diameter distribution. This general approach will be  
182 illustrated during the workshop on a number of test cases.

183

184

## 185 **CONCLUSION**

186

187

188 During the workshop, the various steps of the proposed PNM approach will be presented and its  
189 interest illustrated for charged porous materials. The influence of the surface charge density, the

190 salt concentration in the reservoirs and of the channel diameter distribution will be analyzed. The  
191 symmetry of the transport matrix is preserved by the present upscaling method, as required from  
192 Onsager's theory. In general, the coefficients of this matrix qualitatively behave as their  
193 microscopic counterpart for a channel with the average diameter. However, the combined effects  
194 of electrokinetic couplings on the local scale and of heterogeneity result in a decrease of the  
195 overall transport coefficients, in accordance with Le Châtelier's principle. Overall, the coupling  
196 between the complex pore structure of porous media and electrokinetic effects underlines the  
197 limitations of approaches based on idealized geometries (single slit pore or cylindrical channel)  
198 parametrized directly from the experimental macroscopic properties.

199

200 The relevance and limitations of this new strategy to the case of clay minerals will be discussed.  
201 In that respect, experimental information on the pore network and its size distribution on  
202 intermediate (10-100 nm) scales is highly desirable for the present method to provide more  
203 quantitative predictions in this case. In the future, one should benefit from recent numerical  
204 (Tyagi *et al.*, 2013) and experimental (Brisard *et al.*, 2012; Levitz, 2007) developments for the  
205 generation of realistic numerical samples for the description of real materials. As a recent  
206 example, Robinet *et al.* recently simulated the diffusion of solutes in 3D-images of a Callovo-  
207 Oxfordian clay-rich rock obtained by SEM and micro-CT experiments to investigate the effect of  
208 mineral distribution (Robinet *et al.*, 2012). Multiscale experiments using NMR also provide an  
209 ideal tool to investigate the multiscale dynamics of mobile species in such complex materials  
210 (Porion *et al.*, 2013).

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