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1 **UPSCALING ELECTROKINETIC TRANSPORT IN CLAYS**
2 **WITH LATTICE BOLTZMANN AND PORE NETWORK MODELS**

3
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10 **SHORT TITLE:** Upscaling electrokinetic transport in clays

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14
15 **ABSTRACT**

16 This contribution presents a method for the numerical determination of the steady-state response
17 of complex charged porous media to pressure, salt concentration and electric potential gradients.
18 The Pore Network Model (PNM), describing the porosity as a network of pores connected by
19 channels, is extended to capture electrokinetic couplings which arise at charged solid-liquid
20 interfaces. This allows us to compute the macroscopic fluxes of solvent, salt and charge across a
21 numerical sample submitted to macroscopic gradients. On the channel scale, the microscopic
22 transport coefficients are obtained by solving analytically (in simple cases) or numerically the
23 Poisson-Nernst-Planck and Stokes equations. The PNM approach then allows us to upscale these
24 transport properties to the sample scale, accounting for the complex pore structure of the material
25 via the distribution of channel diameters. The Onsager relations between macroscopic transport
26 coefficients are preserved, as expected. However, electrokinetic couplings combined with the
27 sample heterogeneity result for some macroscopic transport coefficients (e.g. permeability or
28 electro-osmotic coefficient) in qualitative differences with respect to their microscopic
29 counterparts. This underlines the care that should be taken when accounting for transport
30 properties based on a single channel of average diameter.

31
32 **KEYWORDS:** Electrokinetics, Coupled transport, Homogenization, Lattice Boltzmann, Pore
33 Network Model, Upscaling

34 **INTRODUCTION**

35

36 Electrokinetic effects refer to the dynamic coupling between the solvent and charge flows which
37 occur at a charged interface. The presence of surface charge in a porous medium has important
38 practical applications in membrane technology (e.g., ion exchange and water desalination) and in
39 environmental science, since most rocks and soils contain minerals (such as clays) that bear a
40 permanent surface charge. As an example, electro-osmosis generates a solvent flow under an
41 applied electric field, due to the driving of the electrically charged fluid in the vicinity of charged
42 surfaces. Conversely, a pressure gradient induces the flow of a charged fluid, hence, an electric
43 current. In geophysics, the electroseismic effect, by which an electro-magnetic wave is generated
44 from the motion of underground fluids under an applied acoustic wave, is exploited to determine
45 the properties of geological formations (Thompson, 1936; Pride and Haartsen, 1996; Mizutani *et*
46 *al.*, 1976). Streaming potentials and electro-osmotic flows can be measured in the laboratory to
47 characterize the properties of porous media (Luong and Sprik, 2013).

48

49 The modeling and simulation of electrokinetic effects in porous media, and, more generally, of
50 all coupled transport phenomena, including the osmotic solvent flow due to a salt concentration
51 gradient, thus have been the subject of a large number of investigations, both on the pore scale
52 where the couplings originate and on the sample scale corresponding to the experimental
53 measurements. From the mathematical point of view, this upscaling can be performed rigorously
54 using the homogenization approach. This provides expressions of the macroscopic transport
55 coefficients as solutions of coupled partial differential equations on the pore scale, which then
56 have to be solved using simplifying assumptions or numerically. Some general results, such as
57 Onsager's relations for the macroscopic transport coefficients, can be demonstrated without even
58 resorting to the numerical resolution of the mathematical problem (Moyne and Murad, 2006a&b;
59 Allaire *et al.*, 2010&2014).

60

61 For practical applications, most studies of electrokinetic couplings rely on an oversimplified
62 idealization of the geometry, with single slit pores or cylinders with dimensions or surface
63 charge densities estimated from the macroscopic properties of the real system (Bresler, 1973;
64 Gonçalves *et al.*, 2012). However, the heterogeneity of the material, combined with the

65 electrokinetic couplings, may influence the overall behavior on the sample scale, so that such
66 idealizations may not reflect the actual response of the medium. Direct numerical resolution of
67 the coupled Poisson-Nernst-Planck (PNP) and Navier-Stokes (NS) equations in various complex
68 systems (random packings, reconstructed and fractured porous media) has also been proposed by
69 Adler and co-workers. Such an approach is usually difficult to implement for macroscopic
70 samples, due to the lack of experimental data on the fine structure of the material over large
71 distances (Coelho *et al.*, 1996; Marino *et al.*, 2001; Gupta *et al.*, 2006). The systematic study of a
72 representative number of samples is also prevented by the computational cost of direct numerical
73 simulation.

74

75 In the case of clays, an additional difficulty arises due to the complex multiscale porosity of the
76 material and the lack of experimental data on the intermediate scales, which is at the heart of this
77 workshop. In the present contribution, we present a numerical homogenization scheme leading to
78 a description of transport through macroscopic charged porous materials at low computational
79 cost, thereby enabling the systematic study of the combined effects of electrokinetic couplings
80 and sample heterogeneity. The algorithm to upscale the electrokinetic couplings is based on the
81 Pore Network Model (PNM), which relies on the one hand on a simplified description of the
82 electrokinetic transport on the pore scale and on the other hand on a statistical distribution of the
83 geometry of the pores. This allows to investigate how the upscaled electrokinetic properties
84 depend on the heterogeneity of the sample, in addition to the surface charge density and the salt
85 concentration.

86

87

88 **ELECTROKINETICS ON THE SAMPLE SCALE**

89

90 On the macroscopic scale of a clay sample, pressure P , electric potential V and salt concentration
91 gradients (or, equivalently, solvent, cation and anion chemical potentials gradients), induce
92 macroscopic fluxes of mass, electric charge and salt (or, equivalently, solvent, cation and anion
93 chemical fluxes). For sufficiently small applied gradients, the response is linear and the fluxes
94 can be expressed as a function of the applied gradients via a coupling matrix:

95

$$\begin{pmatrix} Q_0 \\ Q_1 \\ Q_2 \end{pmatrix} = -\frac{S}{\eta} \begin{pmatrix} K_0^P & K_0^C & K_0^V \\ K_1^P & K_1^C & K_1^V \\ K_2^P & K_2^C & K_2^V \end{pmatrix} \begin{pmatrix} \nabla P \\ \nabla C \\ \nabla V \end{pmatrix}, \quad (1)$$

96 where the subscripts 0, 1 and 2 refer to solvent, cations and anions, respectively, Q indicate their
 97 fluxes, C is the logarithm of the salt concentration, η the solvent viscosity and S the cross-section
 98 of the sample. Along the diagonal one finds for example the sample permeability K_0^P and the
 99 electric conductivity K_2^V . The fundamental question is then: How do these coefficients emerge
 100 from the microscopic structure of the material, including heterogeneities on intermediate scales,
 101 and from the surface charge density of the solid matrix?
 102
 103
 104

105

106 **ELECTROKINETICS ON THE PORE SCALE**

107

108 Recently, significant progress has been made on the derivation of the macroscopic transport
 109 equations from the pore-scale ones. These studies usually start from a continuous description of
 110 the fluid via transport equations, which are then upscaled to derive their average effect on the
 111 sample scale, which is quantified by a coupling matrix relating the solvent and ionic fluxes to the
 112 corresponding forces (pressure, potential, and concentration gradients). The solvent flow under
 113 applied local forces is accounted for via the NS equation (or even the Stokes equation), which
 114 includes a local force due to electrochemical potential gradients. The solute fluxes are due one
 115 the one hand to the advection by the fluid and on the other hand to the local electrochemical
 116 potential gradients; they can be modelled on this scale using the PNP equations. The limitations
 117 of such continuous descriptions to describe solvent and ion transport in clay nanopores, which
 118 can be assessed using molecular simulations (Botan *et al.*, 2010&2013), will not be discussed
 119 here. Rather, the present discussion focusses on how to upscale this to the macroscopic scale,
 120 since the structure is too complex for a direct resolution on the whole sample.
 121

122

123 The coupled Navier-Stokes and Poisson-Nernst-Planck equations can be solved numerically
 124 using finite element or volume methods. For example, Adler and co-workers used this direct
 125 numerical resolution in various complex systems (random packings, reconstructed and fractured
 126 porous media) (Coelho *et al.*, 1996; Marino *et al.*, 2001), demonstrating in particular a universal
 electrokinetic behaviour if appropriate rescaled quantities are introduced (Gupta *et al.*, 2006,

127 2008). Recently, alternative methods have been proposed to simulate electrokinetic effects
128 starting from a more fundamental description of the fluid than the PNP and NS equation
129 (Pagonabarraga *et al.*, 2010). For example, Capuani *et al.* proposed a hybrid lattice based
130 approach (Lattice Boltzmann Electrokinetics, LBE) to capture the coupling of hydrodynamic
131 flow with ion transport and the simulation of electrokinetic effects in colloidal suspensions
132 (Capuani *et al.*, 2004; Pagonabarraga *et al.*, 2005). Such Lattice Boltzmann simulations have
133 already been applied, without accounting for electrokinetic effects, to realistic rock geometries
134 (Boek and Venturoli, 2010). In the context of the present numerical homogenization, LBE was
135 recently used in a simple cylindrical geometry, in order to assess the validity range of the
136 analytical solution of the linearized problem (Obliger *et al.*, 2013). This simpler analytical
137 solution is then used in the PNM, even though in principle a numerical expression for the
138 transport coefficient on the pore scale may also be used.

139

140

141 **NUMERICAL HOMOGENIZATION VIA A PORE NETWORK MODEL**

142

143 In order to investigate electrokinetic couplings on larger scales, including the effect of the
144 heterogeneity of the material, we have recently proposed a simplified description based on the
145 Pore Network Model (PNM). Such a model, originally developed by Fatt (1956) to predict
146 multiphase flow properties in porous media, describes the porosity as a network of pores
147 connected by channels. It has been extensively used and extended by petrophysicists in various
148 situations, such as capillarity and multiphase flow through porous media (Békri *et al.*, 2005;
149 Blunt, 2001; van Dijke and Sorbie, 2002), or mineral dissolution and precipitation in the context
150 of CO₂ sequestration (Algive *et al.*, 2010).

151

152 In a nutshell, the PNM approach amounts to solving a set of conservation equations on the nodes
153 of the network (in analogy with Kirchhoff's law for a network of resistors), on the basis of local
154 fluxes through the channels connecting the nodes, under the effect of an external, macroscopic
155 gradient. For electrokinetics, the pressure, salt concentration and electrical potential are
156 introduced as pore variables on the nodes of the network. The fluxes through each link between
157 nodes are determined locally using the transport matrix for a cylindrical channel, as determined

158 in the previous section as a function of the channel diameter, the surface charge density of the
159 solid and the salt concentration inside the channel. The latter is determined via the Donnan
160 equilibrium with a fictitious reservoir corresponding to the properties of the pores at both ends of
161 the channel (Obliger *et al.*, 2014). Therefore, the macroscopic problem to be solved numerically
162 has a non-linear structure, contrary to most previous applications of the PNM approach. This can
163 be achieved numerically using a non-linear Newton solver.

164

165

166 **TRANSPORT COEFFICIENTS ON THE SAMPLE SCALE**

167

168 In addition to the transport coefficients on the channel scale, the crucial ingredient of the PNM is
169 the distribution of pore/channel sizes and their spatial arrangement describing in a very
170 simplified manner the complex structure of the porous network. In order to demonstrate the
171 feasibility of the approach and to investigate systematically the effect of heterogeneity, a model
172 distribution (of the Weibull type) was considered first. However, it is also possible to introduce a
173 distribution deduced from experimental data, if a reliable one can be provided.

174

175 For a given pore/channel diameter distribution, a sufficient number of networks must be
176 generated. For each of numerical sample, the macroscopic coefficients are determined by solving
177 the conservation equations in the presence of applied gradients and by computing the
178 macroscopic steady-state flux through the sample. In practice, three calculations must be done
179 (one for each applied gradient) for which the three fluxes (mass, charge, salt concentration) are
180 computed. This provides the nine macroscopic coefficients, which must then be averaged over
181 the networks corresponding to the same diameter distribution. This general approach will be
182 illustrated during the workshop on a number of test cases.

183

184

185 **CONCLUSION**

186

187

188 During the workshop, the various steps of the proposed PNM approach will be presented and its
189 interest illustrated for charged porous materials. The influence of the surface charge density, the

190 salt concentration in the reservoirs and of the channel diameter distribution will be analyzed. The
191 symmetry of the transport matrix is preserved by the present upscaling method, as required from
192 Onsager's theory. In general, the coefficients of this matrix qualitatively behave as their
193 microscopic counterpart for a channel with the average diameter. However, the combined effects
194 of electrokinetic couplings on the local scale and of heterogeneity result in a decrease of the
195 overall transport coefficients, in accordance with Le Châtelier's principle. Overall, the coupling
196 between the complex pore structure of porous media and electrokinetic effects underlines the
197 limitations of approaches based on idealized geometries (single slit pore or cylindrical channel)
198 parametrized directly from the experimental macroscopic properties.

199

200 The relevance and limitations of this new strategy to the case of clay minerals will be discussed.
201 In that respect, experimental information on the pore network and its size distribution on
202 intermediate (10-100 nm) scales is highly desirable for the present method to provide more
203 quantitative predictions in this case. In the future, one should benefit from recent numerical
204 (Tyagi *et al.*, 2013) and experimental (Brisard *et al.*, 2012; Levitz, 2007) developments for the
205 generation of realistic numerical samples for the description of real materials. As a recent
206 example, Robinet *et al.* recently simulated the diffusion of solutes in 3D-images of a Callovo-
207 Oxfordian clay-rich rock obtained by SEM and micro-CT experiments to investigate the effect of
208 mineral distribution (Robinet *et al.*, 2012). Multiscale experiments using NMR also provide an
209 ideal tool to investigate the multiscale dynamics of mobile species in such complex materials
210 (Porion *et al.*, 2013).

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