



HAL
open science

On problems in the calculus of variations in increasingly elongated domains

Hervé Le Dret, Amira Mokrane

► **To cite this version:**

Hervé Le Dret, Amira Mokrane. On problems in the calculus of variations in increasingly elongated domains. 2017. hal-01517824v1

HAL Id: hal-01517824

<https://hal.sorbonne-universite.fr/hal-01517824v1>

Preprint submitted on 3 May 2017 (v1), last revised 18 Jan 2018 (v3)

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

On problems in the calculus of variations in increasingly elongated domains

Hervé Le Dret

Sorbonne Universités, UPMC Univ Paris 06, CNRS,
Laboratoire Jacques-Louis Lions, Boîte courrier 187,
75252 Paris Cedex 05, France. Email: herve.le_dret@upmc.fr

Amira Mokrane

Laboratoire d'équations aux dérivées partielles
non linéaires et histoire des mathématiques, ENS, B.P. 92,
Vieux Kouba, 16050 Alger, Algérie
and USTHB, Faculté des mathématiques, Département d'analyse,
Laboratoire d'analyse mathématique et numérique
des équations aux dérivées partielles, Bab Ezzouar, Alger, Algérie.
Email: mokrane_amira3@yahoo.fr

May 3, 2017

Abstract

We consider minimization problems in the calculus of variations set in a sequence of domains the size of which tends to infinity in certain directions and such that the data only depend on the coordinates in the directions that remain constant. We study the asymptotic behavior of minimizers in various situations and show that they converge in an appropriate sense toward minimizers of a related energy functional in the constant directions.

MSC 2010: 35J25, 35J35, 35J62, 35J92, 49J45, 74K99.

Keywords: Calculus of variations, domains becoming unbounded, asymptotic behavior, exponential rate of convergence.

1 Introduction

In this article, we revisit the “ $\ell \rightarrow +\infty$ ” problem in the context of the calculus of variations. This class of problems was introduced by Chipot and Rougirel in 2000, [6], see also the monograph by Chipot [4], and has since then given rise to many works by several authors dealing with various elliptic and parabolic problems up to until recently.

A prototypical $\ell \rightarrow +\infty$ problem is the following. Let $\omega =]-1, 1[$, $\ell > 0$ be a real number and $\Omega_\ell \subset \mathbb{R}^2$ the rectangle $]-\ell, \ell[\times \omega$. We denote by x_1 the first variable in $]-\ell, \ell[$ and x_2 the second variable in ω . Any function $f \in L^2(\omega)$ in the second variable gives rise to a function in two variables still denoted f by

setting $f(x_1, x_2) = f(x_2)$. We thus consider the two boundary value problems: find u_ℓ , a function in (x_1, x_2) , such that

$$\begin{cases} -\Delta u_\ell = f & \text{in } \Omega_\ell, \\ u_\ell = 0 & \text{on } \partial\Omega_\ell, \end{cases}$$

and find u_∞ , a function in x_2 , such that

$$\begin{cases} -\frac{d^2 u_\infty}{dx_2^2} = f & \text{in } \omega, \\ u_\infty = 0 & \text{on } \partial\omega = \{-1, 1\}. \end{cases}$$

Now the function u_∞ can also be considered as a function in two variables that is independent of x_1 . In this case, it can be shown that, for any $\ell_0 > 0$, one has

$$u_\ell \rightarrow u_\infty \text{ in } H^1(\Omega_{\ell_0}) \text{ when } \ell \rightarrow +\infty,$$

hence the name of the problem. In other words, when the data do not depend on the elongated dimension, the solution of the above boundary value problem converges in some sense at finite distance to the solution of the corresponding boundary value problem posed in the non-elongated dimension when the elongation tends to infinity.

The majority of works on the $\ell \rightarrow +\infty$ problem makes use of the boundary value problem itself, *i.e.*, the PDE plus boundary condition. One exception to this rule is the recent paper [5], in which the authors consider instead a sequence of problems in the calculus of variations posed on elongated domains. This is the approach we adopt here as well.

Our main motivation for this is that certain models, such as nonlinear hyperelasticity, are naturally posed as problems in the calculus of variations for which no Euler-Lagrange equation, *i.e.*, non underlying PDE even in a weak form, is available, see [2]. Moreover, questions surrounding the Saint Venant principle in elasticity, see [14, 15], are typically set in elongated domains, albeit in one direction only. Consequently, it makes sense to attempt dealing with some $\ell \rightarrow +\infty$ problems by using only energy minimization properties and no Euler-Lagrange equation whatsoever.

We are however quite far from achieving the goal of treating nonlinear elasticity, since the approach that we develop below relies a lot on convexity, whereas convexity is not an appropriate hypothesis for nonlinear elasticity. We are nonetheless able to encompass a wide range of nonlinear energies, including the p -Laplacian with some technical restrictions on the number of elongated dimensions with respect to the exponent p . Our hypotheses are weaker and our results are sometimes stronger than those of [5]. The techniques are somewhat different too, with an emphasis here on weak convergence and weak lower semi-continuity techniques, and reliance on such classical techniques as the De Giorgi slicing method which are not dependent on convexity. As a general rule, we try to make as little use of convexity as we can at any given point.

Let us describe our results a little more precisely. We consider bounded open subsets Ω_ℓ of \mathbb{R}^n which are Cartesian products of the form $\ell\omega' \times \omega''$, with $\omega' \subset \mathbb{R}^r$ and $\omega'' \subset \mathbb{R}^{n-r}$, with $1 \leq r \leq n-1$. We let $x = (x', x'')$ with $x' \in \mathbb{R}^r$ being the elongated variable and $x'' \in \mathbb{R}^{n-r}$ the non-elongated variable. Likewise, for a scalar-valued function $v: \Omega_\ell \rightarrow \mathbb{R}$, we decompose the gradient $\nabla v = (\nabla' v, \nabla'' v)$ with obvious notation.

We consider an energy density $F: \mathbb{R}^n \rightarrow \mathbb{R}$ and a function f on ω'' , and introduce the minimization problem of finding $u_\ell \in W_0^{1,p}(\Omega_\ell)$ such that $J_\ell(u_\ell) = \inf_{v \in W_0^{1,p}(\Omega_\ell)} J_\ell(v)$ where

$$J_\ell(v) = \int_{\Omega_\ell} (F(\nabla v(x)) - f(x'')v(x)) dx.$$

We assume that F has p -growth, p -coerciveness and is convex. In particular, there is no assumption of strict convexity or uniform strict convexity made on F .

We then introduce $F'': \mathbb{R}^{n-r} \rightarrow \mathbb{R}$ by letting $F''(\xi'') = F(0, \xi'')$, again with obvious notation. Of course, F'' is convex, has p -growth and p -coerciveness and the minimization problem of finding $u_\infty \in W_0^{1,p}(\omega'')$ such that $J_\infty(u_\infty) = \inf_{v \in W_0^{1,p}(\omega'')} J_\infty(v)$ where

$$J_\infty(v) = \int_{\omega''} (F''(\nabla'' v(x'')) - f(x'')v(x'')) dx'',$$

admits solutions. It turns out that, under additional hypotheses, this problem is the “ $\ell \rightarrow +\infty$ ” limit of the family of minimization problems under consideration.

These hypotheses include appropriate growth and coerciveness hypotheses on the function $G: \mathbb{R}^n \rightarrow \mathbb{R}$, $G(\xi) = F(\xi) - F''(\xi'')$, of the form

$$\forall \xi \in \mathbb{R}^n, \lambda(|\xi'|^p + k|\xi''|^{p-k}|\xi'|^k) \leq G(\xi) \leq \Lambda(|\xi'|^p + k|\xi''|^{p-k}|\xi'|^k),$$

for some $0 < \lambda \leq \Lambda$ and $0 \leq k < p$. Depending on the case, there is no more additional hypothesis (for $k = 0$), or a hypothesis of strict convexity of F'' , or a hypothesis of uniform strict convexity of F'' (for $k > 0$).

The results are a “ $\ell \rightarrow +\infty$ ” convergence in the weak sense for $k = 0$ when $r < p$, sharpened to strong sense when F'' is furthermore assumed to be strictly convex, and a strong “ $\ell \rightarrow +\infty$ ” convergence for $k > 0$ when $r \leq kp/(p-k)$. In the case of the p -Laplacian, $p > 2$, we thus obtain strong “ $\ell \rightarrow +\infty$ ” convergence when $r < 2p/(p-2)$.

In addition, in the case $k = 0$, if we assume that F'' is uniformly strictly convex, we obtain strong convergence at an exponential rate without any restriction on r . This includes the known behavior of the 2-Laplacian in the “ $\ell \rightarrow +\infty$ ” context.

We conclude the article with a few comments and perspectives on the vectorial case, in connection with nonlinear elasticity in particular.

2 Statement of the problem

We consider two bounded open sets $\omega' \subset \mathbb{R}^r$ with $0 \in \omega'$ and ω' is starshaped with respect to 0, and $\omega'' \subset \mathbb{R}^{n-r}$ with $n > r \geq 1$. Let $\ell > 0$ and set

$$\omega'_\ell = \ell\omega' \text{ and } \Omega_\ell = \omega'_\ell \times \omega'' \subset \mathbb{R}^n. \quad (1)$$

Points x in Ω_ℓ will be denoted by $x = (x', x'')$ with $x' = (x_1, x_2, \dots, x_r) \in \omega'_\ell$ and $x'' = (x_{r+1}, \dots, x_n) \in \omega''$. Likewise, vectors ξ in \mathbb{R}^n will be decomposed as $\xi = (\xi', \xi'')$, with $\xi' \in \mathbb{R}^r$ and $\xi'' \in \mathbb{R}^{n-r}$.

Note that because of the starshaped assumption, we have $\Omega_\ell \subset \Omega_{\ell'}$ as soon as $\ell \leq \ell'$ and we are thus dealing with a “growing” family of open sets. We make an additional regularity hypothesis on ω' , which is as follows. Define first the gauge function of ω' as

$$g(x') = \inf\{t \in \mathbb{R}_+^*; x'/t \in \omega'\}.$$

Since ω' is starshaped and bounded, this is well defined, $\omega'_\ell = \{x'; g(x') < \ell\}$, and there exists $0 < R_1 < R_2$ such that $R_1|x'| \leq g(x') \leq R_2|x'|$ for all $x' \in \mathbb{R}^r$.

Now we assume that ω' is such that g is a Lipschitz function with Lipschitz constant K . By Rademacher’s theorem, this implies that g is almost everywhere differentiable, with $|\nabla'g(x')| \leq K$ a.e. Moreover, it is known that g then belongs to $W_{\text{loc}}^{1,\infty}(\mathbb{R}^r)$ and that its almost everywhere derivatives equal its distributional derivatives. This is true for example if ω' is convex. This regularity hypothesis is for convenience only: we use g to build cut-off functions inside the domains, and not up to the boundary. It should be quite clear that our results can be rewritten in order to accommodate arbitrary open sets ω' .

We are interested in a sequence of problems in the calculus of variations \mathcal{P}_ℓ of the form

$$J_\ell(u_\ell) = \inf_{v \in W_0^{1,p}(\Omega_\ell)} J_\ell(v), \quad (2)$$

with $u_\ell \in W_0^{1,p}(\Omega_\ell)$ and

$$J_\ell(v) = \int_{\Omega_\ell} [F(\nabla v(x)) - f''(x'')v(x)] dx. \quad (3)$$

where $f'' \in L^{p'}(\omega'')$, $\frac{1}{p} + \frac{1}{p'} = 1$, is a given function. Observe that the term corresponding to the force term for this problem only depends on the “non-elongated” variable x'' so that it is reasonable to expect that u_ℓ behaves as a function mostly in x'' in the limit $\ell \rightarrow +\infty$, in a sense made precise below. We could also consider more general semilinear force terms of the form $h(x'', v)$ satisfying appropriate growth and convexity assumptions, but we stick here with a linear term for simplicity.

We assume that the energy density $F: \mathbb{R}^n \rightarrow \mathbb{R}$ is convex. We let

$$\begin{aligned} F'' : \mathbb{R}^{n-r} &\rightarrow \mathbb{R} & \text{and} & & G : \mathbb{R}^n &\rightarrow \mathbb{R} \\ \xi'' &\mapsto F(0, \xi'') & & & \xi &\mapsto F(\xi) - F''(\xi'') \end{aligned} \quad (4)$$

so that

$$F(\xi', \xi'') = F''(\xi'') + G(\xi', \xi''), \quad (5)$$

and F'' is convex. These functions are assumed to satisfy the following coerciveness and growth hypotheses

$$\forall \xi \in \mathbb{R}^n, \lambda|\xi|^p \leq F(\xi) \leq \Lambda(|\xi|^p + 1), \quad (6)$$

$$\forall \xi \in \mathbb{R}^n, \lambda(|\xi'|^p + k|\xi''|^{p-k}|\xi'|^k) \leq G(\xi) \leq \Lambda(|\xi'|^p + k|\xi''|^{p-k}|\xi'|^k), \quad (7)$$

for some $0 < \lambda \leq \Lambda$, $p > 1$ and $0 \leq k < p$.¹ Here, for $\xi \in \mathbb{R}^d$, $|\xi|$ denotes the canonical Euclidean norm of ξ in \mathbb{R}^d .

¹Note that $k = p$ yields the same hypothesis as $k = 0$.

Clearly, condition (6) implies the similar condition

$$\forall \xi'' \in \mathbb{R}^{n-r}, \lambda |\xi''|^p \leq F''(\xi'') \leq \Lambda (|\xi''|^p + 1), \quad (8)$$

for F'' .

Energy densities of the form above include that associated with the p -Laplacian for $p \geq 2$. Indeed, in this case, $F(\xi) = \frac{1}{p} |\xi|^p = \frac{1}{p} (|\xi'|^2 + |\xi''|^2)^{p/2}$ and we can take $k = 2$ for $p > 2$, or $k = 0$ for $p = 2$. Another simple energy density that is covered by our analysis is $F(\xi) = \frac{1}{p} (|\xi'|^p + |\xi''|^p)$ or more generally energies of the form $F(\xi) = F'(\xi') + F''(\xi'')$, with appropriate hypotheses on F' and F'' . Here, assuming without loss of generality that $F'(0) = 0$, we have $G(\xi', \xi'') = F'(\xi')$ and we can take $k = 0$.

In addition to the above growth and coerciveness hypotheses, which obviously imply that problem \mathcal{P}_ℓ has at least one solution u_ℓ , we assume that F'' is uniformly strictly convex for $k > 0$, in the sense that there exists a constant $\beta > 0$ such that for all $\xi'', \zeta'' \in \mathbb{R}^{n-r}$ and all $\theta, \mu \in [0, 1]$ with $\theta + \mu = 1$, we have

$$F''(\theta \xi'' + \mu \zeta'') \leq \theta F''(\xi'') + \mu F''(\zeta'') - k \beta \theta \mu (\theta^{p-1} + \mu^{p-1}) |\xi'' - \zeta''|^p. \quad (9)$$

see for instance [1, 10, 12]. The p -Laplacian for $p > 2$, $k = 2$, satisfies this hypothesis (the 2-Laplacian satisfies the alternate hypothesis (53) that will be used later on in Section 5). Note that when $k = 0$, the hypothesis becomes redundant, and there is actually no requirement of even strict convexity, let alone uniform strict convexity, of F'' in this case.

We now introduce our candidate limit problem \mathcal{P}_∞ as that of finding $u_\infty \in W_0^{1,p}(\omega'')$ such that

$$J_\infty(u_\infty) = \inf_{v \in W_0^{1,p}(\omega'')} J_\infty(v), \quad (10)$$

with

$$J_\infty(v) = \int_{\omega''} [F''(\nabla'' v(x'')) - f''(x'') v(x'')] dx''. \quad (11)$$

It is also clear that problem \mathcal{P}_∞ has at least one solution u_∞ .

Here and in the sequel, we use the following notational device

$$\nabla' = (\partial_1, \dots, \partial_r), \quad \nabla'' = (\partial_{r+1}, \dots, \partial_n),$$

that we apply indifferently to functions defined either on Ω_ℓ or on ω'' . For brevity, we refer to ∇' as the ‘‘horizontal’’ part of the gradient and to ∇'' as the ‘‘vertical’’ part of the gradient.

We want to study the asymptotic behavior of u_ℓ when $\ell \rightarrow +\infty$ and compare it with a minimizer u_∞ of the $n - r$ dimensional vertical problem \mathcal{P}_∞ . Actually, our goal is to show that the former converges to the latter in a sense that will be explained later on.

3 Preliminary estimates

We first give several estimates that we will use in the proofs of our convergence results. The first estimate follows immediately from Poincaré’s inequality.

Lemma 3.1 *There exists a constant $c_1 = c_1(\omega'')$ independent of ℓ such that for all $v \in W^{1,p}(\Omega_\ell)$ whose trace vanish on $\omega'_\ell \times \partial\omega''$, we have*

$$\|v\|_{L^p(\Omega_\ell)} \leq c_1 \|\nabla'' v\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})}. \quad (12)$$

Let us now give a first, coarse estimate of u_ℓ .

Lemma 3.2 *There exists a constant c_2 independent of ℓ , such that*

$$\int_{\Omega_\ell} |\nabla u_\ell|^p dx \leq c_2 \ell^r. \quad (13)$$

Proof. Let us take $v = 0$ as a test-function in problem (2). It follows that

$$\int_{\Omega_\ell} F(\nabla u_\ell(x)) dx \leq \int_{\Omega_\ell} f''(x'') u_\ell(x) dx + A \ell^r.$$

where $A = F(0) \mathcal{L}^r(\omega') \mathcal{L}^{n-r}(\omega'')$ does not depend on ℓ (\mathcal{L}^d denotes the d -dimensional Lebesgue measure). By Hölder's inequality and the coerciveness assumption (6), it follows that

$$\begin{aligned} \int_{\Omega_\ell} |\nabla u_\ell(x)|^p dx &\leq \frac{1}{\lambda} \left(\int_{\Omega_\ell} |f''(x'')|^{p'} dx \right)^{1/p'} \left(\int_{\Omega_\ell} |u_\ell(x)|^p dx \right)^{1/p} + \frac{A}{\lambda} \ell^r \\ &\leq \frac{B}{\lambda} \ell^{r/p'} \|\nabla'' u_\ell\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})} + \frac{A}{\lambda} \ell^r, \end{aligned}$$

with $B = c_1 \|f''\|_{L^{p'}(\omega'')} \mathcal{L}^r(\omega')^{1/p'}$, which does not depend on ℓ . Consequently, we obtain an estimate of the form

$$\|\nabla u_\ell\|_{L^p(\Omega_\ell; \mathbb{R}^n)}^p \leq C \ell^{r/p'} \|\nabla u_\ell\|_{L^p(\Omega_\ell; \mathbb{R}^n)} + D \ell^r, \quad (14)$$

where C and D are constants that do not depend on ℓ . Let us set $X = \ell^{-r/p} \|\nabla u_\ell\|_{L^p}$. Estimate (14) now reads

$$X^p \leq CX + D,$$

so that there exists c_2 depending only on C and D such that $X \leq c_2^{1/p}$, which completes the proof. \square

We now recall an elementary estimate similar to what can be found in [11] for instance.

Lemma 3.3 *Let $h(t)$ a nonnegative bounded function defined on an interval $[\tau_0, \tau_1]$, $\tau_0 \geq 0$. Suppose that for $\tau_0 \leq t < s \leq \tau_1$, we have*

$$h(t) \leq \theta h(s) + C(s-t)^{-\nu_1} + D(s-t)^{-\nu_2},$$

where $C, D, \nu_1, \nu_2, \theta$ are nonnegative constants with $0 \leq \theta < 1$. Then, for all $\tau_0 \leq t < s \leq \tau_1$, we have

$$h(t) \leq c(C(s-t)^{-\nu_1} + D(s-t)^{-\nu_2}),$$

where c is a constant that only depends on ν_1, ν_2 and θ .

Proof. If we have two sequences of nonnegative numbers a_i and b_i such that $a_i \leq \theta a_{i+1} + b_{i+1}$, it follows by induction that $a_0 \leq \theta^i a_i + \sum_{j=0}^{i-1} \theta^j b_{j+1}$. We apply this remark to the sequences $a_i = h(t_i)$ and $b_{i+1} = C(t_{i+1} - t_i)^{-\nu_1} + D(t_{i+1} - t_i)^{-\nu_2}$, where $t_i = t + (1 - \sigma^i)(s - t)$, $0 < \sigma < 1$ to be chosen later on, is an increasing sequence in $[\tau_0, \tau_1]$ such that $t_0 = t$. This yields the estimate

$$h(t) \leq \theta^i h(t_i) + \frac{C}{(s-t)^{\nu_1}} (1-\sigma)^{-\nu_1} \sum_{j=0}^{i-1} \left(\frac{\theta}{\sigma^{\nu_1}} \right)^j + \frac{D}{(s-t)^{\nu_2}} (1-\sigma)^{-\nu_2} \sum_{j=0}^{i-1} \left(\frac{\theta}{\sigma^{\nu_2}} \right)^j.$$

We now choose $\sigma < 1$ in such a way that $\frac{\theta}{\sigma^{\nu_1}} < 1$ and $\frac{\theta}{\sigma^{\nu_2}} < 1$, and conclude by letting $i \rightarrow +\infty$, remembering that $h(t_i)$ is bounded. \square

Next, we estimate the horizontal part of the gradient of u_ℓ in $L^p(\Omega_{\ell_0})$ in terms of ℓ , ℓ_0 , u_ℓ and a minimizer u_∞ of the vertical problem \mathcal{P}_∞ .

Theorem 3.4 *There exists a constant c_3 independent of all the other quantities such that, for all $0 < t < s \leq \ell$ and all minimizers u_∞ of the vertical problem, we have*

$$\begin{aligned} & \|\nabla' u_\ell\|_{L^p(\Omega_t; \mathbb{R}^r)}^p + k \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_t; \mathbb{R}^{n-r})}^p \\ & \leq \frac{\delta c_3}{(s-t)^p} \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_s \setminus \Omega_t; \mathbb{R}^{n-r})}^p \\ & + \frac{c_3 k}{(s-t)^{kp/(p-k)}} \left\{ \|\nabla'' u_\ell\|_{L^p(\Omega_s \setminus \Omega_t; \mathbb{R}^{n-r})}^p + (1-\delta) \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_s \setminus \Omega_t; \mathbb{R}^{n-r})}^p \right\}. \end{aligned} \quad (15)$$

where $\delta = 1$ if $0 \leq k \leq p/2$, $\delta = 0$ otherwise.

Proof. We first define a family of cut-off functions as follows. For all $0 < t < s \leq \ell$, we set

$$\rho_{s,t}(x') = \frac{1}{s-t} \min\{(s - g(x'))_+, s - t\}.$$

By the definition of the gauge function, we see that $\rho_{s,t} \equiv 0$ on $\omega'_\ell \setminus \omega'_s$, $\rho_{s,t} \equiv 1$ on ω'_t and $0 \leq \rho_{s,t} \leq 1$. By our regularity assumption on ω' , $\rho_{s,t}$ is Lipschitz and such that

$$\nabla' \rho_{s,t}(x') = -\frac{1}{s-t} \nabla' g(x') \mathbf{1}_{\omega'_s \setminus \omega'_t}(x'),$$

so that we can estimate

$$|\nabla' \rho_{s,t}(x')| \leq \frac{K}{s-t} \mathbf{1}_{\omega'_s \setminus \omega'_t}(x'). \quad (16)$$

We pick a number $0 < \alpha < 1$ and then set

$$v_1(x) = (1 - \alpha \rho_{s,t}(x')) u_\ell(x) + \alpha \rho_{s,t}(x') u_\infty(x''), \quad (17)$$

and

$$v_2(x) = (1 - \alpha \rho_{s,t}(x')) u_\infty(x'') + \alpha \rho_{s,t}(x') u_\ell(x). \quad (18)$$

Clearly, v_1 belongs to $W_0^{1,p}(\Omega_\ell)$ and is thus a suitable test-function for problem \mathcal{P}_ℓ , hence

$$\int_{\Omega_\ell} [F(\nabla u_\ell(x)) - f''(x'') u_\ell(x)] dx \leq \int_{\Omega_\ell} [F(\nabla v_1(x)) - f''(x'') v_1(x)] dx. \quad (19)$$

Next we note that, owing to the embedding $W_0^{1,p}(\Omega_\ell) \hookrightarrow L^p(\omega'_\ell; W_0^{1,p}(\omega''))$, v_2 is suitable test-function for problem \mathcal{P}_∞ for almost all x' , hence

$$\begin{aligned} \int_{\omega''} [F''(\nabla'' u_\infty(x'')) - f''(x'')u_\infty(x'')] dx'' \\ \leq \int_{\omega''} [F''(\nabla'' v_2(x', x'')) - f''(x'')v_2(x', x'')] dx''. \end{aligned} \quad (20)$$

Integrating estimate (20) over ω'_ℓ , we obtain

$$\begin{aligned} \int_{\Omega_\ell} [F''(\nabla'' u_\infty(x'')) - f''(x'')u_\infty(x'')] dx \\ \leq \int_{\Omega_\ell} [F''(\nabla'' v_2(x)) - f''(x'')v_2(x)] dx. \end{aligned} \quad (21)$$

We add estimates (19) and (21) together and note that all the terms involving f'' cancel out since $v_1 + v_2 = u_\ell + u_\infty$. Therefore,

$$\int_{\Omega_\ell} [F(\nabla u_\ell(x)) + F''(\nabla'' u_\infty(x''))] dx \leq \int_{\Omega_\ell} [F(\nabla v_1(x)) + F''(\nabla'' v_2(x))] dx. \quad (22)$$

We observe that $v_1 = u_\ell$ and $v_2 = u_\infty$ on $\Omega_\ell \setminus \Omega_s$, so that estimate (22) boils down to

$$\begin{aligned} \int_{\Omega_s} [F(\nabla u_\ell(x)) + F''(\nabla'' u_\infty(x''))] dx \leq \int_{\Omega_s \setminus \Omega_t} [F(\nabla v_1(x)) + F''(\nabla'' v_2(x))] dx \\ + \int_{\Omega_t} [F(\nabla v_1(x)) + F''(\nabla'' v_2(x))] dx. \end{aligned} \quad (23)$$

The left-hand side of (23) can be rewritten as

$$\begin{aligned} \int_{\Omega_s} [F(\nabla u_\ell(x)) + F''(\nabla'' u_\infty(x''))] dx \\ = \int_{\Omega_s} [G(\nabla u_\ell(x)) + F''(\nabla'' u_\ell(x)) + F''(\nabla'' u_\infty(x''))] dx. \end{aligned} \quad (24)$$

Let I_1 and I_2 be the first and second integrals in the right-hand side of (23). To estimate I_1 , we just use the convexity of F'' , since the vertical gradients of v_1 and v_2 are convex combinations of the vertical gradients of u_ℓ and u_∞ ,

$$\begin{aligned} I_1 &= \int_{\Omega_s \setminus \Omega_t} [G(\nabla v_1(x)) + F''(\nabla'' v_1(x)) + F''(\nabla'' v_2(x))] dx \\ &\leq \int_{\Omega_s \setminus \Omega_t} [G(\nabla v_1(x)) + F''(\nabla'' u_\ell(x)) + F''(\nabla'' u_\infty(x))] dx. \end{aligned} \quad (25)$$

To estimate I_2 , we note that $v_1 = (1 - \alpha)u_\ell + \alpha u_\infty$ and $v_2 = \alpha u_\ell + (1 - \alpha)u_\infty$

on Ω_t , thus owing to the convexity of F and the uniform convexity (9) of F'' ,

$$\begin{aligned} I_2 &\leq \int_{\Omega_t} [(1-\alpha)F(\nabla u_\ell) + \alpha F(\nabla u_\infty) + (1-\alpha)F''(\nabla'' u_\infty) + \alpha F''(\nabla'' u_\ell) \\ &\quad - k\gamma|\nabla''(u_\infty - u_\ell)|^p] dx \\ &= \int_{\Omega_t} [(1-\alpha)G(\nabla u_\ell) + F''(\nabla'' u_\ell) + F''(\nabla'' u_\infty) - k\gamma|\nabla''(u_\infty - u_\ell)|^p] dx, \end{aligned} \quad (26)$$

for some $\gamma > 0$. Putting estimates (23), (25), (26) and equation (24) together, we obtain

$$\int_{\Omega_s \setminus \Omega_t} G(\nabla u_\ell) dx + \int_{\Omega_t} [\alpha G(\nabla u_\ell) + k\gamma|\nabla''(u_\infty - u_\ell)|^p] dx \leq \int_{\Omega_s \setminus \Omega_t} G(\nabla v_1(x)) dx,$$

which, upon using the coerciveness hypothesis (7), yields

$$\begin{aligned} a \int_{\Omega_t} [(|\nabla' u_\ell|^p + k|\nabla'' u_\ell|^{p-k}|\nabla' u_\ell|^k) + k|\nabla''(u_\infty - u_\ell)|^p] dx \\ \leq \int_{\Omega_s \setminus \Omega_t} G(\nabla v_1(x)) dx, \end{aligned} \quad (27)$$

where $a > 0$ is a small generic constant that only depends on the other constants involved.

We now focus on estimating the right-hand side of (27). We have

$$\begin{cases} \nabla' v_1 = (1 - \alpha\rho_{s,t})\nabla' u_\ell + \alpha\nabla' \rho_{s,t}(u_\infty - u_\ell), \\ \nabla'' v_1 = (1 - \alpha\rho_{s,t})\nabla'' u_\ell + \alpha\rho_{s,t}\nabla'' u_\infty. \end{cases} \quad (28)$$

Based on (28) and the definition of $\rho_{s,t}$, we have the following estimates for any exponent q :

$$\begin{cases} |\nabla' v_1|^q \leq 2^{q-1}|\nabla' u_\ell|^q + 2^{q-1}\frac{K^q}{(s-t)^q}|u_\infty - u_\ell|^q, \\ |\nabla'' v_1|^q \leq 2^{q-1}|\nabla'' u_\ell|^q + 2^{q-1}|\nabla''(u_\infty - u_\ell)|^q. \end{cases} \quad (29)$$

We will use exponents $q = p$ and $q = k$ for the first line and $q = p - k$ for the second line. Due to the growth hypothesis (7), we have

$$\begin{aligned} G(\nabla v_1) &\leq A \left(|\nabla' u_\ell|^p + \frac{1}{(s-t)^p}|u_\infty - u_\ell|^p \right. \\ &\quad \left. + k(|\nabla'' u_\ell|^{p-k} + |\nabla''(u_\infty - u_\ell)|^{p-k}) \left(|\nabla' u_\ell|^k + \frac{1}{(s-t)^k}|u_\infty - u_\ell|^k \right) \right), \end{aligned} \quad (30)$$

where A is a large generic constant that only depends on the other constants involved. For $k \geq 1$, three of the four product terms that appear need to be estimated. For this purpose, we will use Young's inequality in the following form

$$a^k b^{p-k} \leq \frac{k}{p} a^p + \frac{p-k}{p} b^p$$

for $a, b \geq 0$ (recall that $p > k$). We thus obtain

$$G(\nabla v_1) \leq A \left(|\nabla' u_\ell|^p + \frac{1}{(s-t)^p} |u_\infty - u_\ell|^p + k \left(|\nabla'' u_\ell|^{p-k} |\nabla' u_\ell|^k + |u_\infty - u_\ell|^p + \frac{1}{(s-t)^{kp/(p-k)}} |\nabla'' u_\ell|^p + |\nabla''(u_\infty - u_\ell)|^p \right) \right), \quad (31)$$

where A is another generic constant. We integrate this inequality over $\Omega_s \setminus \Omega_t$ and use Poincaré's inequality in the vertical variables to obtain

$$\int_{\Omega_s \setminus \Omega_t} G(\nabla v_1) dx \leq A \int_{\Omega_s \setminus \Omega_t} (|\nabla' u_\ell|^p + k |\nabla'' u_\ell|^{p-k} |\nabla' u_\ell|^k + k |\nabla''(u_\infty - u_\ell)|^p) dx + \frac{A}{(s-t)^p} \|\nabla''(u_\infty - u_\ell)\|_{L^p(\Omega_s \setminus \Omega_t)}^p + \frac{Ak}{(s-t)^{kp/(p-k)}} \|\nabla'' u_\ell\|_{L^p(\Omega_s \setminus \Omega_t)}^p, \quad (32)$$

with A yet another generic constant.

We now consider two different cases. First, for $0 \leq k \leq p/2$, let us set

$$h(t) = \int_{\Omega_t} (|\nabla' u_\ell|^p + k |\nabla'' u_\ell|^{p-k} |\nabla' u_\ell|^k + k |\nabla''(u_\infty - u_\ell)|^p) dx.$$

Inequalities (27) and (32) may be rewritten as

$$h(t) \leq \theta h(s) + \frac{1}{(s-t)^p} \|\nabla''(u_\infty - u_\ell)\|_{L^p(\Omega_s \setminus \Omega_t)}^p + \frac{k}{(s-t)^{kp/(p-k)}} \|\nabla'' u_\ell\|_{L^p(\Omega_s \setminus \Omega_t)}^p, \quad (33)$$

with $\theta = \frac{A}{A+a} \in]0, 1[$. Let $t \leq t_1 < s_1 \leq s$. We invoke Lemma 3.3, with $\nu_1 = p$, $C = \|\nabla''(u_\infty - u_\ell)\|_{L^p(\Omega_s \setminus \Omega_t)}^p$, $\nu_2 = kp/(p-k)$, $D = k \|\nabla'' u_\ell\|_{L^p(\Omega_s \setminus \Omega_t)}^p$, to conclude that

$$h(t_1) \leq c(C(s_1 - t_1)^{-\nu_1} + D(s_1 - t_1)^{-\nu_2}). \quad (34)$$

The result follows in this case by letting $t_1 \rightarrow t$ and $s_1 \rightarrow s$ since the constant c only depends on ν_1 , ν_2 and θ , and h is continuous (recall that $\delta = 1$).

Now the second case is when $p/2 < k < p$. Estimate (32) still holds true, but we now use Young's inequality once more in the form

$$a^w b^{p-w} \leq \frac{w}{p} a^p + \frac{p-w}{p} b^p$$

with $w = \frac{p(2k-p)}{k}$ to deduce that

$$\frac{1}{(s-t)^p} = 1^w \left(\frac{1}{(s-t)^{p/(p-w)}} \right)^{p-w} \leq \frac{p-k}{k} + \frac{2k-p}{k} \frac{1}{(s-t)^{kp/(p-k)}},$$

so that we can actually write

$$h(t) \leq \theta h(s) + \frac{k}{(s-t)^{kp/(p-k)}} (\|\nabla''(u_\infty - u_\ell)\|_{L^p(\Omega_s \setminus \Omega_t)}^p + \|\nabla'' u_\ell\|_{L^p(\Omega_s \setminus \Omega_t)}^p), \quad (35)$$

with the same function h , but with another value for θ , which we do not write here. We conclude as before with Lemma 3.3 and the first constant $C = 0$ for instance. \square

The following is an immediate consequence of the previous estimate.

Corollary 3.5 *We have, for all $\ell \geq \ell_0$,*

$$\begin{aligned} & \|\nabla' u_\ell\|_{L^p(\Omega_{\ell_0}; \mathbb{R}^r)}^p + k \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_{\ell_0}; \mathbb{R}^{n-r})}^p \\ & \leq \frac{\delta c_3}{(\ell - \ell_0)^p} \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})}^p \\ & + \frac{c_3 k}{(\ell - \ell_0)^{kp/(p-k)}} \left\{ \|\nabla'' u_\ell\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})}^p + (1 - \delta) \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})}^p \right\}, \end{aligned} \quad (36)$$

where $\delta = 1$ if $0 \leq k \leq p/2$, $\delta = 0$ otherwise.

Proof. Indeed, we take $s = \ell$, $t = \ell_0$ and notice that $\Omega_\ell \setminus \Omega_{\ell_0} \subset \Omega_\ell$. \square

Let us remark that if $k = 0$ and there is actually no strict convexity assumption made on F'' , i.e., F'' may well be not strictly convex, the previous result boils down to

$$\|\nabla' u_\ell\|_{L^p(\Omega_{\ell_0}; \mathbb{R}^r)}^p \leq \frac{c_3}{(\ell - \ell_0)^p} \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_\ell; \mathbb{R}^{n-r})}^p.$$

However, when $k > 0$, we make crucial use of the uniform strict convexity to derive the estimate.

Let us close this section with an estimate similar to that obtained in Lemma 3.2. Recall that u_ℓ is a minimizer on Ω_ℓ , whereas the following estimate is on Ω_{ℓ_0} .

Lemma 3.6 *There exist constants $\bar{\ell}$ and c_4 , independent of ℓ , such that for all $\bar{\ell} \leq \ell_0 \leq \ell$,*

$$\int_{\Omega_{\ell_0}} |\nabla u_\ell|^p dx \leq c_4 \bar{\ell}_0^r. \quad (37)$$

Proof. Let $1 \leq t \leq \ell - 1$ and set $\rho_t = \rho_{t+1,t}$. We take $v_{t,\ell} = (1 - \rho_t)u_\ell$ as a test-function in problem (2). This test-function is equal to u_ℓ “far away” and is 0 in Ω_t . We obtain

$$\begin{aligned} \int_{\Omega_\ell} F(\nabla u_\ell) dx & \leq \int_{\Omega_\ell} (F(\nabla v_{t,\ell}) - f''(v_{t,\ell} - u_\ell)) dx \\ & = \int_{\Omega_t} F(0) dx + \int_{\Omega_t} f'' u_\ell dx + \int_{\Omega_{t+1} \setminus \Omega_t} [F(\nabla v_{t,\ell}) + f'' \rho_t u_\ell] dx \\ & \quad + \int_{\Omega_\ell \setminus \Omega_{t+1}} F(\nabla u_\ell) dx. \end{aligned}$$

Therefore, we see that

$$\int_{\Omega_{t+1}} F(\nabla u_\ell) dx \leq At^r + \int_{\Omega_{t+1}} \nu_t f'' u_\ell dx + \int_{\Omega_{t+1} \setminus \Omega_t} F(\nabla v_{t,\ell}) dx,$$

with $A = F(0)\mathcal{L}^r(\omega')\mathcal{L}^{n-r}(\omega'')$ and $\nu_t = \mathbf{1}_{\Omega_t} + \rho_t \mathbf{1}_{\Omega_{t+1} \setminus \Omega_t}$.

By the coerciveness and growth hypotheses (6), we infer that

$$\lambda \int_{\Omega_{t+1}} |\nabla u_\ell|^p dx \leq Bt^r + \int_{\Omega_{t+1}} |f'' u_\ell| dx + \Lambda \int_{\Omega_{t+1} \setminus \Omega_t} |\nabla v_{t,\ell}|^p dx,$$

for some constant B , since $0 \leq \nu_t \leq 1$ and the Lebesgue measure of $\Omega_{t+1} \setminus \Omega_t$ is of the order of t^{r-1} .

In $\Omega_{t+1} \setminus \Omega_t$, we have

$$|\nabla v_{t,\ell}|^p = |(1 - \rho_t)\nabla u_\ell - u_\ell \nabla \rho_t|^p \leq 2^{p-1}(|\nabla u_\ell|^p + K^p |u_\ell|^p).$$

Clearly, estimate (12) is also valid on $\Omega_{t+1} \setminus \Omega_t$, thus,

$$\int_{\Omega_{t+1} \setminus \Omega_t} |u_\ell|^p dx \leq c_1^p \int_{\Omega_{t+1} \setminus \Omega_t} |\nabla'' u_\ell|^p dx \leq c_1^p \int_{\Omega_{t+1} \setminus \Omega_t} |\nabla u_\ell|^p dx,$$

so that

$$\int_{\Omega_{t+1} \setminus \Omega_t} |\nabla v_{t,\ell}|^p dx \leq 2^{p-1}(1 + c_1^p K^p) \int_{\Omega_{t+1} \setminus \Omega_t} |\nabla u_\ell|^p dx.$$

Furthermore,

$$\begin{aligned} \int_{\Omega_{t+1}} |f'' u_\ell| dx &\leq \frac{\varepsilon}{p} \int_{\Omega_{t+1}} |u_\ell|^p dx + \frac{(t+1)^r}{\varepsilon^{p'/p} p'} \mathcal{L}^r(\omega') \|f''\|_{L^{p'}(\omega'')}^{p'} \\ &\leq \frac{\varepsilon c_1^p}{p} \int_{\Omega_{t+1}} |\nabla u_\ell|^p dx + \frac{C}{\varepsilon^{p'/p}} t^r, \end{aligned}$$

with $\varepsilon > 0$ to be chosen afterwards.

Let us set

$$h(t) = \int_{\Omega_t} |\nabla u_\ell|^p dx.$$

Putting all the above estimates together, it follows that

$$\lambda' h(t+1) \leq E(h(t+1) - h(t)) + Dt^r, \quad (38)$$

with $\lambda' = \lambda - \frac{\varepsilon c_1^p}{p}$, $D = B + \frac{C}{\varepsilon^{p'/p}}$ and $E = 2^{p-1} \Lambda(1 + c_1^p K^p)$. We now pick ε in such a way that $\lambda' > 0$. Inequality (38) may be rewritten as

$$h(t) \leq \theta h(t+1) + Ht^r, \quad (39)$$

where $\theta = 1 - \frac{\lambda'}{E} \in]0, 1[$ and $H = \frac{D}{E}$ depend neither on t nor on ℓ . Iterating inequality (39), we see that for $n = \lfloor \ell - t \rfloor$, we have

$$h(t) \leq \theta^n h(t+n) + H \sum_{m=0}^{n-1} (t+m)^r \theta^m. \quad (40)$$

Let us now set $t = \ell_0$. We have $h(\ell_0 + \lfloor \ell - \ell_0 \rfloor) \leq h(\ell) \leq c_2 \ell^r$ by Lemma 3.2. Hence

$$\theta^{\lfloor \ell - \ell_0 \rfloor} h(t + \lfloor \ell - \ell_0 \rfloor) \leq c_2 \theta^{\lfloor \ell - \ell_0 \rfloor} \ell^r \leq c_2 \theta^{\ell - \ell_0 - 1} \ell^r.$$

Now, for $\ell_0 \geq -\frac{r}{\ln \theta}$, the function in the right-hand side is decreasing, hence maximum for $\ell = \ell_0$. Therefore,

$$\theta^{\lfloor \ell - \ell_0 \rfloor} h(t + \lfloor \ell - \ell_0 \rfloor) \leq \frac{c_2}{\theta} \ell_0^r,$$

for $\ell \geq \ell_0 \geq -\frac{r}{\ln \theta}$. Moreover, for $\ell_0 \geq 1$,

$$\begin{aligned} \sum_{m=0}^{\lfloor \ell - \ell_0 \rfloor - 1} (\ell_0 + m)^r \theta^m &= \ell_0^r \sum_{m=0}^{\lfloor \ell - \ell_0 \rfloor - 1} \left(1 + \frac{m}{\ell_0}\right)^r \theta^m \\ &\leq \ell_0^r \sum_{m=0}^{\lfloor \ell - \ell_0 \rfloor - 1} (1 + m)^r \theta^m \leq \frac{\sum_{m=1}^{+\infty} m^r \theta^m}{\theta} \ell_0^r, \end{aligned}$$

which completes the proof with $\bar{\ell} = \max(1, -\frac{r}{\ln \theta})$. \square

We now turn to the convergence results. As a consequence of Lemma 3.6, we have, without any restriction on r with respect to p and k ,

Theorem 3.7 *There exists a subsequence $\ell \rightarrow +\infty$ and a function $u^* \in W_{\text{loc}}^{1,p}(\Omega_\infty)$ such that, for all ℓ_0 ,*

$$u_{\ell|\Omega_{\ell_0}} \rightharpoonup u_{|\Omega_{\ell_0}}^* \text{ weakly in } W^{1,p}(\Omega_{\ell_0}). \quad (41)$$

Moreover, $u^* = 0$ on $\partial\Omega_\infty$.

Note that the weak convergence above implies that $u_\ell \rightharpoonup u^*$ weakly in $W_{\text{loc}}^{1,p}(\Omega_\infty)$. We will sometimes omit the restriction notation in the sequel when unnecessary.

Proof. By estimates (12) and (37), for all $n \in \mathbb{N}^*$, u_ℓ is bounded in $W^{1,p}(\Omega_n)$. Using the diagonal procedure, we thus construct a sequence ℓ_n such that for all m , $u_{\ell_n|\Omega_m} \rightharpoonup u_m^*$ weakly in $W^{1,p}(\Omega_m)$, with $u_m = 0$ on $\omega'_m \times \partial\omega''$. Now, since $\Omega_m \subset \Omega_{m'}$ as soon as $m \leq m'$, it follows that $u_m^* = u_{m'|\Omega_m}^*$, so that we have constructed a single limit function u^* in the desired class. Furthermore, for all ℓ_0 , if we choose an integer $m \geq \ell_0$, we see that convergence (41) holds true. \square

In the sequel, we will always consider a weakly convergent subsequence u_ℓ in the sense of Theorem 3.7.

4 Identification of the limit when $\ell \rightarrow +\infty$

In this section, we do not make any further use of assumption (9) of uniform strict convexity of F'' , other than the fact that we used it to establish Theorem 3.4.² The results will only hold for values of r small enough depending on p . We let $\Omega_\infty = \mathbb{R}^r \times \omega''$.

Let us first show that the asymptotic behavior of u_ℓ is independent of the elongated dimension if r is small enough.

Theorem 4.1 *Assume that $r < p$ if $k = 0$, or that $r < kp/(p-k)$ if $0 < k < p$. Then we have $\nabla' u^* = 0$ and u^* may be identified with a function in the x'' variable only, still denoted u^* , which belongs to $W_0^{1,p}(\omega'')$.*

Proof. By estimates (13) and (36) and the triangle inequality, it follows that

$$\|\nabla' u_\ell\|_{L^p(\Omega_{\ell_0}; \mathbb{R}^r)}^p \leq C \left(\frac{\delta}{(\ell - \ell_0)^p} + \frac{k}{(\ell - \ell_0)^{kp/(p-k)}} \right) \ell^r \rightarrow 0 \quad (42)$$

²Keep in mind that this hypothesis is void for $k = 0$ anyway.

when $\ell \rightarrow +\infty$ with ℓ_0 fixed. Indeed, when $0 < k \leq p/2$, we actually have $\frac{kp}{p-k} \leq p$ and since $\ell \rightarrow +\infty$, the first term in the right hand side of estimate (42) is bounded from above by the second term.

Now $\nabla' u_\ell \rightharpoonup \nabla' u^*$ weakly in $L^p_{\text{loc}}(\Omega_\infty)$, hence we see that $\nabla' u^* = 0$, which concludes the proof of the Theorem. \square

In order to get a feeling of what Theorem 4.1 says, let us look at a few examples. For the Laplacian, we have $p = 2$ and we can take $k = 0$, which restricts this result to $r = 1$ (see Section 5 for a more general result with additional hypotheses, that applies in this case). For the p -Laplacian, $p > 2$, we can take $k = 2$ and the result is restricted to $r < 2p/(p-2)$. Note that $r = 1$ and $r = 2$ are allowed for any value of p . This is not optimal in this particular case, since it is known that $\ell \rightarrow +\infty$ convergence holds without restriction on the dimension with respect to p , see [8].

Let us now identify the limit function. We first need another estimate.

Lemma 4.2 *There exists a constant c_5 such that for all $t \leq s$,*

$$\limsup_{\ell \rightarrow +\infty} \int_{\Omega_s \setminus \Omega_t} |\nabla u_\ell|^p dx \leq c_5 (s^r - t^r). \quad (43)$$

Proof. We may assume that $t > 0$, since the case $t = 0$ is already covered by Lemma 3.6. We use here De Giorgi's classical slicing trick. Let n be an integer large enough so that $0 \leq t - \frac{1}{n} < s + \frac{1}{n} \leq \ell$. For each integer m , $1 \leq m \leq n$, we consider the cut-off function

$$\chi_{m,n}(x') = \rho_{s+\frac{m}{n^2}, s+\frac{m-1}{n^2}}(x') \left(1 - \rho_{t-\frac{m-1}{n^2}, t-\frac{m}{n^2}}(x') \right).$$

This cut-off function takes its values in $[0, 1]$, it is 0 whenever $g(x') \geq s + \frac{m}{n^2}$ or $g(x') \leq t - \frac{m}{n^2}$, it is 1 for $t - \frac{m-1}{n^2} \leq g(x') \leq s + \frac{m-1}{n^2}$, and $|\nabla \chi_{m,n}| \leq Kn^2$. Let us call $S_{m,n}$ the slice where $0 < \chi_{m,n}(x') < 1$. We observe that

$$\bigcup_{m=1}^n \overline{S_{m,n}} = \overline{\Omega_{s+\frac{1}{n}} \setminus \Omega_s} \cup \overline{\Omega_t \setminus \Omega_{t-\frac{1}{n}}} \subset \Omega_{s+1},$$

and that $S_{m,n} \cap S_{m',n} = \emptyset$ when $m \neq m'$.

Let us consider the test-function $v_{\ell,m,n} = (1 - \chi_{m,n})u_\ell + \chi_{m,n}u^*$. The minimization problem yields the estimate

$$\begin{aligned} \int_{\Omega_\ell} F(\nabla u_\ell) dx &\leq \int_{\Omega_\ell} F(\nabla v_{\ell,m,n}) dx - \int_{\Omega_\ell} f'' \chi_{m,n} (u^* - u_\ell) dx \\ &= \int_{\Omega_\ell} F(\nabla v_{\ell,m,n}) dx - \int_{\Omega_{s+1}} f'' \chi_{m,n} (u^* - u_\ell) dx. \end{aligned}$$

Taking into account the specific form of the cut-off function, this implies that

$$\begin{aligned}
\int_{\Omega_s \setminus \Omega_t} F(\nabla u_\ell) dx &\leq \int_{\Omega_{s+\frac{m}{n^2}} \setminus \Omega_{t-\frac{m}{n^2}}} F(\nabla u_\ell) dx \\
&\leq \int_{\Omega_{s+\frac{m}{n^2}} \setminus \Omega_{t-\frac{m}{n^2}}} F(\nabla v_{\ell,m,n}) dx - \int_{\Omega_{s+1}} f'' \chi_{m,n}(u^* - u_\ell) dx \\
&\leq \int_{S_{m,n}} F(\nabla v_{\ell,m,n}) dx + \int_{\Omega_{s+\frac{m-1}{n^2}} \setminus \Omega_{t-\frac{m-1}{n^2}}} F(\nabla u^*) dx \\
&\quad - \int_{\Omega_{s+1}} f'' \chi_{m,n}(u^* - u_\ell) dx. \tag{44}
\end{aligned}$$

Let us estimate each term in the right-hand side separately. First of all, we have

$$\left| \int_{\Omega_{s+1}} f'' \chi_{m,n}(u^* - u_\ell) dx \right| \leq A^{1/p'} (s+1)^{r/p'} \|f''\|_{L^{p'}(\omega'')} \|u^* - u_\ell\|_{L^p(\Omega_{s+1})}, \tag{45}$$

with $A = \mathcal{L}^r(\omega')$. Secondly, we see that

$$\left| \int_{\Omega_{s+\frac{m-1}{n^2}} \setminus \Omega_{t-\frac{m-1}{n^2}}} F(\nabla u^*) dx \right| \leq A \left(\left(s + \frac{1}{n} \right)^r - \left(t - \frac{1}{n} \right)^r \right) \|F''(\nabla'' u^*)\|_{L^1(\omega'')}. \tag{46}$$

We now come to the slicing argument *stricto sensu*. By the growth estimate (6), we have

$$\begin{aligned}
\int_{S_{m,n}} F(\nabla v_{\ell,m,n}) dx &\leq 2^{p-1} \Lambda \left(\int_{S_{m,n}} (|\nabla u_\ell|^p + |\nabla u^*|^p + 1) dx \right. \\
&\quad \left. + K^p n^{2p} \int_{S_{m,n}} |u^* - u_\ell|^p dx \right). \tag{47}
\end{aligned}$$

The only term that causes a difficulty is the last term coming from $\nabla \chi_{m,n}$. We now plug estimates (45), (46) and (47) into the right-hand side of estimate (44), sum for $m = 1$ to n and divide the result by n . Observing that the sum of integrals over the slices $S_{m,n}$ gives rise to integrals over the union of all slices, which is included in Ω_{s+1} , this yields

$$\int_{\Omega_s \setminus \Omega_t} F(\nabla u_\ell) dx \leq A^{1/p'} (s+1)^{r/p'} \|f''\|_{L^{p'}(\omega'')} \|u^* - u_\ell\|_{L^p(\Omega_{s+1})} \tag{48}$$

$$+ A \left(\left(s + \frac{1}{n} \right)^r - \left(t - \frac{1}{n} \right)^r \right) \|F''(\nabla'' u^*)\|_{L^1(\omega'')} \tag{49}$$

$$+ \frac{2^p \Lambda c_4}{n} (s+1)^r + 2^{p-1} \Lambda K^p n^{2p-1} \|u^* - u_\ell\|_{L^p(\Omega_{s+1})}^p. \tag{50}$$

We first let $\ell \rightarrow +\infty$. Due to the Rellich-Kondrašov theorem, $\|u^* - u_\ell\|_{L^p(\Omega_{s+1})} \rightarrow 0$ and it follows from the coerciveness estimate that

$$\begin{aligned}
\limsup_{\ell \rightarrow +\infty} \int_{\Omega_s \setminus \Omega_t} |\nabla u_\ell|^p dx &\leq \frac{A}{\lambda} \left(\left(s + \frac{1}{n} \right)^r - \left(t - \frac{1}{n} \right)^r \right) \|F''(\nabla'' u^*)\|_{L^1(\omega'')} \\
&\quad + \frac{2^p \Lambda c_4}{n \lambda} (s+1)^r.
\end{aligned}$$

We finally let $n \rightarrow +\infty$ to obtain the result with $c_5 = \frac{A}{\lambda} \|F''(\nabla'' u^*)\|_{L^1(\omega'')} \cdot \square$

We now are in a position to prove the main result of this section.

Theorem 4.3 *The function u^* is a minimizer of problem \mathcal{P}_∞ .*

Proof. Let $z \in W_0^{1,p}(\omega'')$ be arbitrary. We use the test function $v_\ell = (1 - \rho_t)u_\ell + \rho_t z$, with $\rho_t = \rho_{t+1,t}$, so that $v_\ell = u_\ell$ on $\Omega_\ell \setminus \Omega_{t+1}$ and $v_\ell = z$ on Ω_t . We thus have

$$\int_{\Omega_{t+1}} [F(\nabla u_\ell) - f''u_\ell] dx \leq \int_{\Omega_{t+1} \setminus \Omega_t} [F(\nabla v_\ell) - f''v_\ell] dx + \int_{\Omega_t} [F(\nabla z) - f''z] dx. \quad (51)$$

It follows from Lemma 4.2 that

$$\limsup_{\ell \rightarrow +\infty} \left| \int_{\Omega_{t+1} \setminus \Omega_t} [F(\nabla v_\ell) - f''v_\ell] dx \right| \leq C(t+1)^{r-1}$$

for some constant C independent of ℓ and t . The left-hand side of estimate (51) is weakly lower-semicontinuous, hence, letting $\ell \rightarrow +\infty$, we obtain

$$\begin{aligned} (t+1)^r \mathcal{L}^r(\omega') \int_{\omega''} [F(\nabla u^*) - f''u^*] dx' &\leq C(t+1)^{r-1} \\ &+ t^r \mathcal{L}^r(\omega') \int_{\omega''} [F(\nabla z) - f''z] dx' \end{aligned}$$

and the result follows from letting $t \rightarrow +\infty$, since $F(\nabla u^*) = F''(\nabla''u^*)$ and $F(\nabla z) = F''(\nabla''z)$. \square

We now apply a classical trick to obtain strong convergence when F'' is strictly convex. Of course, when $k > 0$, this is already the case by assumption (9). Strict convexity is only a new assumption if $k = 0$. In this case, the solution u_∞ of the limit problem is unique and this uniqueness implies the weak convergence of the whole family u_ℓ .

Theorem 4.4 *Assume that F'' is strictly convex. Then $u^* = u_\infty$ and $u_\ell \rightarrow u_\infty$ strongly in $W^{1,p}(\Omega_{\ell_0})$ for all ℓ_0 .*

We recall the following two lemmas that can be found *e.g.* in [3].

Lemma 4.5 *Let $F: \mathbb{R}^M \rightarrow \mathbb{R}$ be strictly convex. Let $\mu \in]0, 1[$ and $a_j, a \in \mathbb{R}^M$ such that*

$$\mu F(a_j) + (1 - \mu)F(a) - F(\mu a_j + (1 - \mu)a) \rightarrow 0 \text{ as } j \rightarrow +\infty.$$

Then $a_j \rightarrow a$.

The second lemma is a slight variation on Fatou's lemma.

Lemma 4.6 *Let $F_j, F, H_j, H \in L^1(\Omega)$ with $F_j \geq H_j \geq 0$ for all j , $F_j \rightarrow F$ and $H_j \rightarrow H$ a.e., and $\int_\Omega F_j dx \rightarrow \int_\Omega F dx$. Then*

$$\int_\Omega H_j dx \rightarrow \int_\Omega H dx.$$

Proof of Theorem 4.4. We already know that $\nabla' u_\ell \rightarrow 0 = \nabla' u^*$ strongly in $L^p(\Omega_t)$ by estimate (42). We thus just have to prove the strong convergence of $\nabla'' u_\ell$.

We use a similar slicing as before, with the test-functions $\rho_{t+\frac{m}{n^2}, t+\frac{m-1}{n^2}} u_\ell$ for n large enough, $1 \leq m \leq n$. Skipping the details, this slicing implies that

$$\limsup_{\ell \rightarrow +\infty} \int_{\Omega_t} F(\nabla u_\ell) dx \leq \int_{\Omega_t} F(\nabla u^*) dx.$$

On the other hand, for almost all x' , the function $u_{x', \ell}: x'' \mapsto u_\ell(x', x'')$ is an admissible test-function for the limit problem, so that

$$\int_{\omega''} [F''(\nabla'' u^*) - f'' u^*] dx'' \leq \int_{\omega''} [F''(\nabla'' u_{x', \ell}) - f'' u_{x', \ell}] dx''.$$

We integrate this inequality with respect to $x' \in t\omega'$ and obtain

$$\int_{\Omega_t} [F''(\nabla'' u^*) - f'' u^*] dx \leq \int_{\Omega_t} [F''(\nabla'' u_\ell) - f'' u_\ell] dx.$$

We now let $\ell \rightarrow +\infty$, which yields

$$\int_{\Omega_t} F''(\nabla'' u^*) dx \leq \liminf_{\ell \rightarrow +\infty} \int_{\Omega_t} F''(\nabla'' u_\ell) dx.$$

By hypothesis (7), $G \geq 0$, which implies that $F''(\xi'') \leq F(\xi', \xi'')$ for any ξ' . It follows that

$$\int_{\Omega_t} F''(\nabla'' u_\ell) dx \rightarrow \int_{\Omega_t} F''(\nabla'' u^*) dx \quad (52)$$

when $\ell \rightarrow +\infty$, since $F''(\nabla'' u^*) = F(\nabla u^*)$.

Let us pick $\mu \in]0, 1[$ and set

$$g_\ell = \mu F''(\nabla'' u_\ell) + (1 - \mu) F''(\nabla'' u^*) - F''(\mu \nabla'' u_\ell + (1 - \mu) \nabla'' u^*).$$

By weak lower semicontinuity, it is clear that

$$\liminf_{\ell \rightarrow +\infty} \int_{\Omega_t} F''(\mu \nabla'' u_\ell + (1 - \mu) \nabla'' u^*) dx \geq \int_{\Omega_t} F''(\nabla'' u^*) dx.$$

Therefore

$$0 \leq \limsup_{\ell \rightarrow +\infty} \int_{\Omega_t} g_\ell dx \leq \int_{\Omega_t} F''(\nabla'' u^*) dx - \int_{\Omega_t} F''(\nabla'' u^*) dx = 0,$$

so that $g_\ell \rightarrow 0$ a.e. (up to a subsequence). We then apply Lemma 4.5 to deduce that $\nabla'' u_\ell \rightarrow \nabla'' u^*$ a.e. up to that same subsequence.

We now let

$$H_\ell = |\nabla'' u_\ell - \nabla'' u^*|^p \leq 2^{p-1} (F''(\nabla'' u_\ell) + |\nabla'' u^*|^p) = F_\ell,$$

and invoke Lemma 4.6 and (52) to obtain the result for $\ell_0 = t$. To conclude for all ℓ_0 , we use the diagonal process. \square

5 Convergence rates

In the previous section, we obtained convergence results without taking advantage of the term involving k in the left-hand side of estimate (36). This makes them valid in particular for $k = 0$ without strict or uniform strict convexity. It should however be clear that for $k > 0$, the term in question can be used to obtain a much shorter convergence proof with convergence rate, which we do not detail here. More precisely,

Theorem 5.1 *Under the previous hypotheses with $0 < k < p$ and $r < kp/(p - k)$, we have*

$$\|u_\ell - u_\infty\|_{W^{1,p}(\Omega_{\ell_0})}^p \leq C\ell^r - \frac{kp}{p-k}.$$

The proof is a direct consequence of Corollary 3.5 and Lemma 3.2.

In any case, the estimates do not seem to allow a convergence proof without any restriction on r with respect to p in all generality, whereas it is known in some cases, for instance in the case of the Laplacian, that convergence holds true for all values of r .

In order to partially overcome these shortcomings, we assume now that $k = 0$ and that F'' is uniformly strictly convex in the sense that

$$F''(\theta\xi'' + \mu\zeta'') \leq \theta F''(\xi'') + \mu F''(\zeta'') - \beta\theta\mu(\theta^{p-1} + \mu^{p-1})|\xi'' - \zeta''|^p, \quad (53)$$

for some $\beta > 0$. Note that this is equivalent to allowing $k = p$ in hypotheses (7) and (9). In some sense, $\frac{kp}{p-k}$ is then infinite and it is to be expected that there should be no restriction on the allowed dimensions r , plus faster than polynomial convergence. This is what we now proceed to show.

Under assumption (53), it is fairly clear that we still have an estimate similar to that of Theorem 3.4, namely,

$$\|\nabla' u_\ell\|_{L^p(\Omega_t)}^p + \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_t)}^p \leq \frac{C}{(s-t)^p} \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_s \setminus \Omega_t)}^p. \quad (54)$$

Let us thus prove that not only does convergence hold without restrictions on the elongated dimension r , but that it also occurs at an exponential rate. The extra control makes things actually much easier.

Theorem 5.2 *Under hypotheses (6)-(7) with $k = 0$ and (53), then for all $r \leq n$ and all ℓ_0 , there exist constants C and $\alpha > 0$ independent of ℓ such that we have*

$$\|\nabla(u_\ell - u_\infty)\|_{L^p(\Omega_{\ell_0})} \leq Ce^{-\alpha\ell}.$$

Proof. We take $s = t + 1$ in estimate (54), which yields

$$\begin{aligned} \|\nabla' u_\ell\|_{L^p(\Omega_t)}^p + \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_t)}^p &\leq C\|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_{t+1} \setminus \Omega_t)}^p \\ &\leq C\|\nabla' u_\ell\|_{L^p(\Omega_{t+1} \setminus \Omega_t)}^p \\ &\quad + C\|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_{t+1} \setminus \Omega_t)}^p. \end{aligned}$$

Setting

$$g(t) = \|\nabla' u_\ell\|_{L^p(\Omega_t)}^p + \|\nabla''(u_\ell - u_\infty)\|_{L^p(\Omega_t)}^p,$$

we have just shown that

$$g(t) \leq C(g(t+1) - g(t)),$$

or in other words

$$g(t) \leq \theta g(t+1) \tag{55}$$

with $\theta = \frac{C}{1+C} \in]0, 1[$.

We iterate inequality (55) using the sequence $t_n = n + \ell_0$, $n = 0, \dots, \lfloor \ell - \ell_0 \rfloor$. Obviously

$$g(\ell_0) = g(t_0) \leq \theta^n g(t_n)$$

for all such n , and in particular for the last one,

$$g(\ell_0) \leq \theta^{\lfloor \ell - \ell_0 \rfloor} g(t_{\lfloor \ell - \ell_0 \rfloor}) \leq \theta^{\ell - \ell_0 - 1} g(\ell) \leq C \theta^{-\ell_0 - 1} e^{\ell \ln \theta} \ell^r,$$

with $\ln \theta < 0$. Now, for all r , we can pick α such that $\ln \theta < p\alpha < 0$ and $e^{\ell \ln \theta} \ell^r \leq e^{-p\alpha \ell}$ for ℓ large enough, which completes the proof since $\nabla' u_\infty = 0$. \square

Theorem 5.2 applies to energies of the form $F(\xi) = F'(\xi') + F''(\xi'')$, for instance. We recover in particular the known result for the case of the 2-Laplacian.

6 Extension to the vectorial case

We have written everything so far in the context of a scalar problem, *i.e.*, the functions u_ℓ are scalar-valued. All previous developments only made use of the minimization problem, under various convexity assumptions. Now clearly, absolutely nothing is changed if we consider instead vector-valued problems in the calculus of variations, with functions u_ℓ taking their values in some \mathbb{R}^N , if the energies are supposed to satisfy the same growth, coercivity and convexity assumptions as before, and the same convergence results hold true.

Unfortunately, in the vectorial case of the calculus of variations, the relevant condition that guarantees lower-semicontinuity of the energy functional is not convexity, but much weaker conditions such as quasiconvexity, or in the case of energies that can take the value $+\infty$, as is the case in nonlinear elasticity, polyconvexity, see [9]. Indeed, convexity is not suitable in nonlinear elasticity for well-known modeling reasons. This explains why we have striven to use as little convexity as possible (in some sense) at any given point in the sequence of arguments. This comment should however be mitigated by the fact that some instances of our uses of convexity will also work with rank-1-convexity, which is a reasonable assumption in the vectorial case. There are also notions of strict uniform quasiconvexity that may apply, see [10].

The fact that the Euler-Lagrange equation is not available in nonlinear elasticity is also an incentive to try and only use the minimization problem. Now, it is at this point unclear to us how to attack the elongation problem in such nonconvex vectorial cases, since we still heavily rely on (strict uniform) convexity at crucial points of the proofs. Moreover, the Dirichlet boundary condition considered here is not necessarily the most interesting one in the context of nonlinear elasticity, in particular if we have the Saint Venant principle in mind.

Even the potential limit problem is not so clear. In another dimension reduction context, when considering a body whose thickness goes to zero, and

with different boundary conditions, it can be seen that quasiconvexity is not conserved through an “algebraic” formula of the kind found here, and that a relaxation step is necessary, see for instance [13]. Physically, this due to the possibility of crumpling such a thin body. A similar phenomenon may quite possibly happen here, but maybe not in the same fashion.

To the best of our knowledge, the nonconvex vectorial case remains open.

References

- [1] H. Attouch and D. Aze Approximation and regularization of arbitrary functions in Hilbert spaces by the Lasry-Lions method, *Annales de l’Institut Henri Poincaré (C) Non Linear Analysis*, 10 (1993), 289–312.
- [2] J.M. Ball, Convexity conditions and existence theorems in nonlinear elasticity, *Arch. Rational Mech. Anal.*, 63 (1977), 337–403.
- [3] J.M. Ball and J.E. Marsden, Quasiconvexity at the boundary, positivity of the second variation and elastic stability, *Arch. Rat. Mech. Anal.*, 86, 251–277, 1984.
- [4] M. Chipot, *ℓ goes to plus infinity*, Birkhäuser 2002.
- [5] M. Chipot, A. Mojsic and P. Roy, On some variational problems set on domains tending to infinity. *Discrete Contin. Dyn. Syst.* 36 (2016), no. 7, 3603–3621.
- [6] M. Chipot and A. Rougirel, Sur le comportement asymptotique de la solution de problèmes elliptiques dans des domaines cylindriques tendant vers l’infini, *C. R. Acad. Sci. Paris Sér. I Math.* 331 (2000), no. 6, 435–440.
- [7] M. Chipot and A. Rougirel, On the asymptotic behaviour of the solution of parabolic problems in cylindrical domains of large size in some directions. *Discrete Contin. Dyn. Syst. Ser. B* 1 (2001), no. 3, 319–338.
- [8] M. Chipot and Y. Xie, Some issues on the p -Laplace equation in cylindrical domains. (English summary) *Tr. Mat. Inst. Steklova* 261 (2008), Differ. Uravn. i Din. Sist., 293–300 ; translation in *Proc. Steklov Inst. Math.* 261 (2008), no. 1, 287–294.
- [9] B. Dacorogna, *Direct Methods in the Calculus of Variations*, second edition. Springer Verlag, 2000.
- [10] L.C. Evans, Quasiconvexity and partial regularity in the calculus of variations, *Arch. Rational Mech. Anal.*, 95 (1986), 227–252.
- [11] M. Giaquinta, *Introduction to regularity theory for nonlinear elliptic systems*. Lectures in Mathematics ETH Zürich. Birkhäuser Verlag, Basel, 1993.
- [12] A. Juditsky and Y. Nesterov, Deterministic and stochastic primal-dual sub-gradient algorithms for uniformly convex minimization, *Stoch. Syst.*, Volume 4 (2014), Number 1, 44–80.

- [13] H. Le Dret and A. Raoult, The nonlinear membrane model as variational limit of nonlinear three-dimensional elasticity. *J. Maths. Pures Appl.* 74 (1995), 549–578.
- [14] A. Mielke, Normal hyperbolicity of center manifolds and Saint-Venant’s principle, *Arch. Rational Mech. Anal.*, 110 (1990), 353–372. doi:10.1007/BF00393272
- [15] R.A. Toupin, Saint-Venant’s principle, *Arch. Rational Mech. Anal.*, 18 (1965), 83–96.