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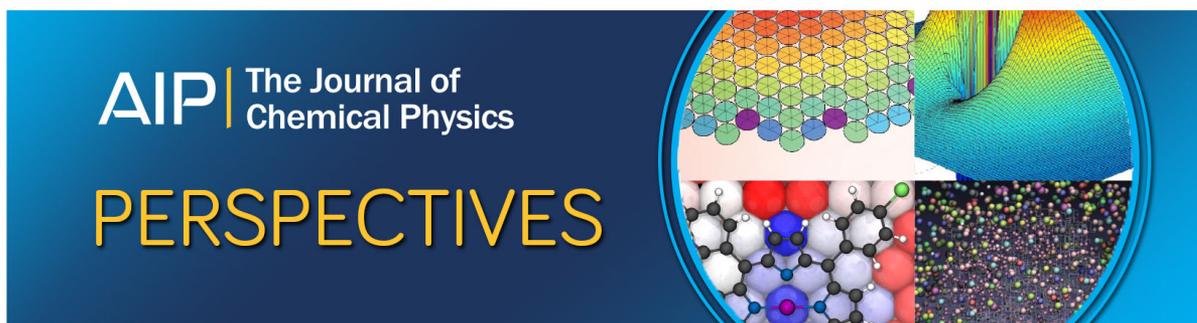
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Double cross polarization for the indirect detection of nitrogen-14 nuclei in magic angle spinning NMR spectroscopy

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Nitrogen-14 NMR spectra at fast magic-angle spinning rates can be acquired indirectly by means of two-dimensional techniques based on double cross polarization transfer $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$. Experimental evidence is given for polycrystalline samples of glycine, L-histidine, and the dipeptide Ala-Gly. Either one-bond or long-range correlations can be favored by choosing the length of the cross polarization contact pulses. Longer contact pulses allow the detection of unprotonated nitrogen sites. In contrast to earlier methods that exploited second-order quadrupolar/dipolar cross-terms, cross polarization operates in the manner of the method of Hartmann and Hahn, even for ^{14}N quadrupolar couplings up to 4 MHz. Simulations explain why amorphous samples tend to give rise to featureless spectra because the ^{14}N quadrupolar interactions may vary dramatically with the lattice environment. The experiments are straightforward to set up and are shown to be effective for different nitrogen environments and robust with respect to the *rf*-field strengths and to the ^{14}N carrier frequency during cross polarization. The efficiency of indirect detection of ^{14}N nuclei by double cross polarization is shown to be similar to that of isotopically enriched ^{13}C nuclei. *Published by AIP Publishing.* <https://doi.org/10.1063/1.5000689>

INTRODUCTION

Nitrogen is a fundamental and ubiquitous element that plays a key role in a variety of fields such as biology, organic chemistry, and materials science. Appearing both in amino acids and nucleotides—the building blocks of proteins and nucleic acids—nitrogen lies at the heart of the “structural patterns” upon which life is based on Earth. The possibility to acquire nitrogen NMR spectra is therefore essential. Unfortunately, its spin-1/2 isotope ^{15}N has a very low natural abundance of ca. 0.37%. Most NMR studies relying on ^{15}N , such as the popular investigations of the structure and dynamics of proteins, must resort to costly isotopic enrichment. In contrast, the highly abundant ^{14}N isotope (ca. 99.63%) has a quadrupolar nucleus with $I = 1$. This feature results in (i) solution-state signals that are usually too broad to be observed, except for highly symmetrical ammonium ions and (ii) solid-state signals that suffer from inhomogeneous broadening up to several MHz if the environments have a low symmetry. In contrast to half-integer quadrupolar spins, nitrogen-14 does not have any central transition $m_I = +1/2 \leftrightarrow m_I = -1/2$, which is usually much easier to observe since it is only affected by small second-order interactions that typically lie in the kHz range.¹ As a result, the NMR observation of ^{14}N is very challenging, even in solid samples that are spinning about the magic angle.

A commonly employed technique for the direct observation of less abundant nuclei such as ^{13}C and ^{15}N is cross polarization (CP).^{2,3} Typically CP exploits the transfer of magnetization from abundant high- γ nuclei with $I = 1/2$ to less abundant low- γ nuclei of interest with $S = 1/2$. The CP process yields a gain in sensitivity that approaches the ratio γ_I/γ_S . This transfer is mediated by dipolar couplings between the two nuclei. In order to achieve this transfer, the Hartmann-Hahn matching condition has to be fulfilled,² with Rabi nutation frequencies $\omega_I = \omega_S$ or $\omega_I = \omega_S \pm n\omega_R$ for static or spinning samples, respectively, and $n = 1$ or 2 . For quadrupolar nuclei with $S > 1/2$, CP is usually less effective. The sensitivity enhancement is more difficult to achieve and the matching conditions are more complex. The introduction of multiple-quantum (MQ)⁴ and satellite-transition (ST)⁵ methods has led to increased interest in solid-state NMR of half-integer spins. As a consequence, interest in CP involving quadrupolar nuclei has been renewed.^{6–10} Nonetheless, mostly half-integer quadrupolar spins with $S = 3/2, 5/2, \text{ or } 7/2$ —all of which have a favorable central transition—have been investigated in this context, with only very few studies concerning nuclei with $S = 1$.^{11–13}

Current methods for solid-state ^{14}N NMR spectroscopy rely on (i) piecewise acquisition¹⁴ or frequency-swept excitation^{15–17} in static samples, (ii) magic angle spinning (MAS) spectra utilizing small flip angle pulses so as to achieve broad excitation profiles,^{18,19} (iii) overtone irradiation,^{20–28} or (iv) heteronuclear multiple-quantum correlation (HMQC) experiments for the indirect detection of ^{14}N spectra via “spy” nuclei such as ^1H or ^{13}C .^{29–43} The HMQC approach offers advantages

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in terms of sensitivity as magnetization on the target spin ^{14}N is created initially by transfer from ^1H and ^{13}C nuclei. Detailed reviews are available of both direct⁴⁴ and indirect⁴⁵ methods for the acquisition of nitrogen-14 NMR spectra.

In this letter, we investigate the use of double cross polarization (DCP) $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ for the indirect detection of nitrogen-14 spectra in powdered samples under MAS conditions. In contrast with early studies that utilized cross polarization for the indirect detection of double-quantum (DQ) nitrogen-14 coherences in static single crystals,^{46,47} we shall focus on single-quantum (SQ) nitrogen-14 coherences in polycrystalline powders that are spinning rapidly about the magic angle.

RESULTS AND DISCUSSION

Figure 1 depicts the pulse sequence proposed in this study. In contrast to HMQC-type experiments, where either SQ or DQ ^{14}N coherences are created via ^{14}N *rf* pulses acting on anti-phase heteronuclear single-quantum spin states such as $H_y N_z$ and $H_x N_z^2$ that result from an initial free evolution of H_x coherence under scalar interactions or second-order cross terms, a simple cross-polarization step in the manner of Hartmann and Hahn will be used in this study to convert in-phase H_y coherence directly into in-phase transverse N_x magnetization. In terms of irreducible tensor operators T_{lm} of rank l and coherence order m , the HMQC approach that we and others have advocated for a spin pair with $I=1/2$ and $S=1$ may be depicted by $T_{1\pm 1}^I \rightarrow T_{1\pm 1}^I T_{10}^S + T_{1\pm 1}^I T_{20}^S$ followed by *rf* irradiation applied to the ^{14}N nuclei to create anti-phase SQ terms $T_{1\pm 1}^I T_{1\pm 1}^S + T_{1\pm 1}^I T_{2\pm 1}^S$ as well as anti-phase DQ terms $T_{1\pm 1}^I T_{2\pm 2}^S$. Such heteronuclear multiple quantum coherences tend to have short lifetimes since the $T_{1\pm 1}^I$ tensor components decay rapidly due to strong proton-proton interactions. By contrast, the much simpler CP process can be described by $T_{1\pm 1}^I \rightarrow T_{1\pm 1}^S$, leading to the creation of pure in-phase ^{14}N single-quantum coherence that may have a longer lifetime and hence give rise to narrower signals. This $T_{1\pm 1}^S$ state is then allowed to evolve under chemical shifts and second-order quadrupolar interactions during an evolution interval t_1 that must be stepped in rotor-synchronized time increments $\Delta t_1 = n\tau_R = n/v_R$ ($n=1, 2, 3, \dots$) so as to refocus the effects of first-order quadrupolar interactions that cancel at the top of rotational echoes, provided the magic angle is set accurately. An identical CP step is then used to transfer the ^{14}N single-quantum $T_{1\pm 1}^S$ coherence back to the proton coherence $T_{1\pm 1}^I$, the signals of which are finally acquired. Importantly,

since we are dealing with *single-spin* states $T_{1\pm 1}^S$ rather than *product* states such as $T_{1\pm 1}^I T_{1\pm 1}^S$ or $T_{1\pm 1}^I T_{2\pm 1}^S$, our double CP experiments do not need any π pulse applied to the protons in the middle of the t_1 evolution interval in order to refocus modulations due to the offsets of the protons. Such refocusing pulses obviously fail to cancel the effects of proton-proton dipole-dipole interactions, which therefore contribute to the decay of $T_{1\pm 1}^I T_{1\pm 1}^S$ and $T_{1\pm 1}^I T_{2\pm 1}^S$ terms and hence to the linewidths in the ω_1 dimension of HMQC-type experiments. These homogeneous losses—leading to line-broadenings—may obscure the inhomogeneous features due to second-order quadrupolar couplings and hinder the extraction of the quadrupolar coupling constant C_Q and asymmetry parameter η_Q .

Figure 2 shows experimental 2D spectra acquired with our novel DCP approach for a polycrystalline powder of L-histidine. The spectrum in Fig. 2(a) was acquired with a contact time $\tau_{CP} = 240 \mu\text{s} = 15\tau_R$ (e.g., $n=15$), a duration that favors one-bond transfers in this sample. Strong one-bond correlations arising from the NH_3^+ and $\text{N}^{\epsilon 2}\text{H}$ sites are clearly identified. Weaker correlations between the NH_3^+ nitrogen-14 and more remote protons may also be distinguished. Figure 2(b) shows an analogous spectrum acquired with a longer contact time $\tau_{CP} = 6.4 \text{ ms} = 400\tau_R$ (e.g., $n=400$), a duration that allows polarization to be transferred between nitrogen-14 sites and both close and remote protons, so that all ^{14}N sites appear to be correlated with all protons. Most importantly, the unprotonated $\text{N}^{\delta 1}$ site is clearly identified in the latter spectrum. To the best of our knowledge, this nitrogen-14 site has never been observed before by indirect detection. The position of the carrier frequency $\nu_{rf}(^{14}\text{N})$ in the spectrum of Fig. 2(b) leads to two horizontal ridges (corresponding to the NH_3^+ and $\text{N}^{\epsilon 2}\text{H}$ protons) that appear near the edge of the spectral width in the vertical ω_1 domain. This finding suggests that the experiment is remarkably robust with respect to the position of the carrier frequency $\nu_{rf}(^{14}\text{N})$. The ^{14}N line shapes obtained in the ω_1 projections of these two ridges reveal how second-order features and intensities may vary critically depending on the length of the CP contact time and the carrier position. For a given set of experimental parameters, the efficiency of the CP process may therefore not be uniform for all crystallite orientations, thus resulting in a partial excitation of the line shapes. The quadrupolar coupling constant C_Q may therefore be underestimated, and the asymmetry parameter η_Q derived from fitting the line shape may not be quite accurate.

Figure 3(a) shows the DCP spectrum of a polycrystalline powder of the dipeptide Ala-Gly acquired with a contact time

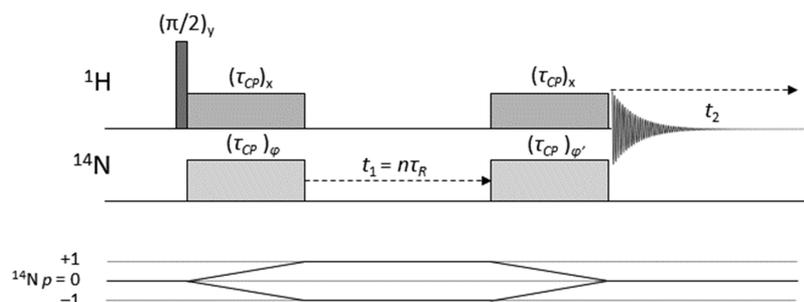


FIG. 1. Pulse sequence for double cross polarization (DCP) for the indirect detection of ^{14}N signals with corresponding coherence transfer pathway. Phase cycling is given in the [supplementary material](#).

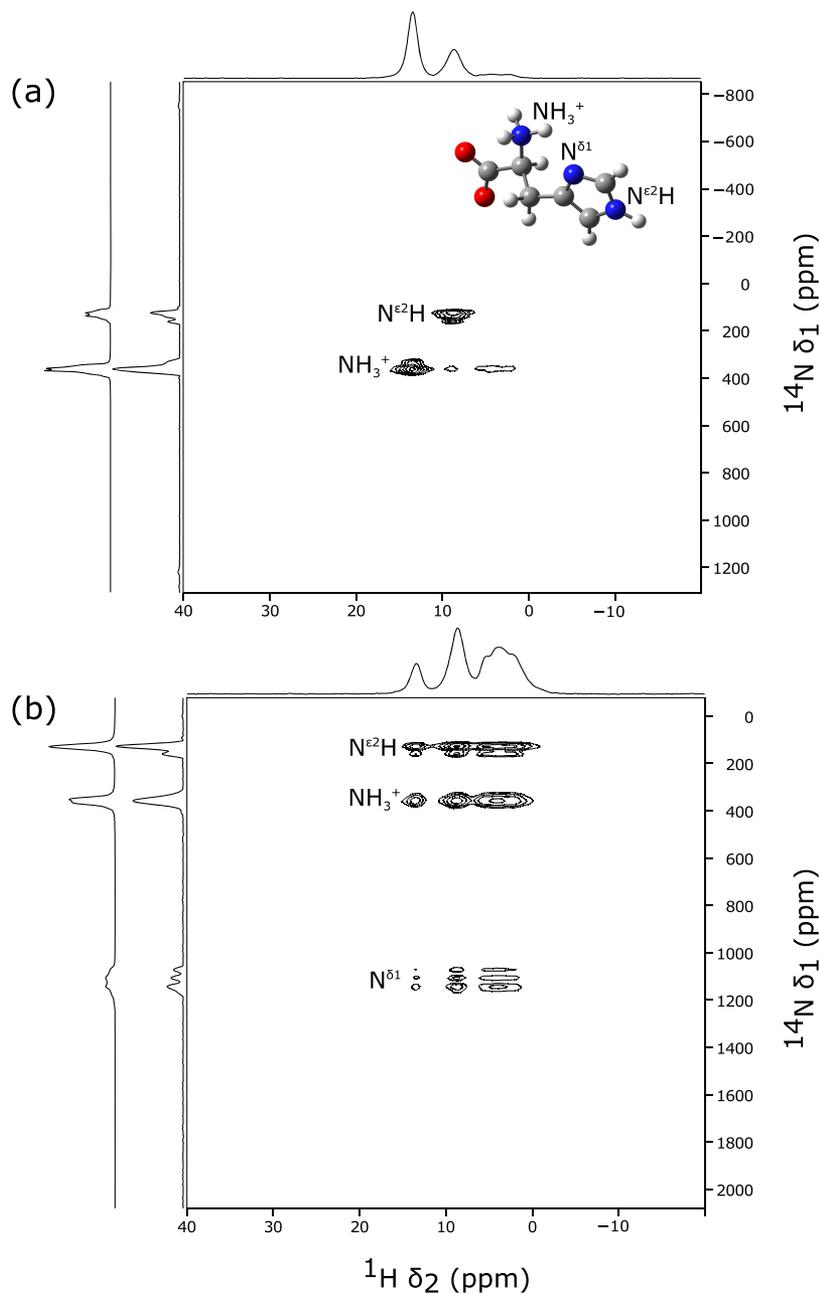


FIG. 2. (a) Experimental 2D ^1H - ^{14}N single-quantum correlation spectrum of L-histidine at $B_0 = 9.4$ T (400.1 MHz for ^1H , 28.9 MHz for ^{14}N) recorded with a spinning frequency of $\nu_R = 62.5$ kHz (rotor period $\tau_R = 16$ μs) and with a contact time of $\tau_{CP} = 15\tau_R = 240$ μs . The most intense one-bond correlations correspond to the NH_3^+ and $\text{N}^{\epsilon 2}\text{H}$ sites. Weaker longer-range correlations corresponding to more distant protons can also be identified. (b) Spectrum recorded with a longer contact time of $\tau_{CP} = 400\tau_R = 6.4$ ms. The unprotonated nitrogen $\text{N}^{\delta 1}$ of the imidazole moiety can now be identified. Numerical simulations of the projections onto the ω_1 axes are shown on the left side for both cases. A recovery delay of 1 s was used and 64 transients were averaged in both cases. The number of t_1 increments was 200 for (a) and 160 for (b).

$\tau_{CP} = 160$ $\mu\text{s} = 10\tau_R$, which is most effective for one-bond transfers. Signals of both NH_3^+ and amide NH fragments, which are expected to have very different C_Q constants of ca. 1 and 3 MHz, respectively, can be clearly identified. This demonstrates that our method is efficient over a wide range of C_Q values and is therefore suitable for detecting a wide variety of nitrogen-14 sites. Numerical simulations of the ω_1 projections are shown on the left side along with the experimental ones, with relevant parameters given in the [supplementary material](#).

In order to evaluate the sensitivity of our $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ DCP method, a $^1\text{H} \rightarrow ^{13}\text{C} \rightarrow ^1\text{H}$ DCP spectrum was acquired for comparison on a sample of ^{13}C - and ^{15}N -enriched L-histidine. The grey contours in Fig. 3(b) show signals of all four $^1\text{H}_n$ ^{13}C groups. A contact time $\tau_{CP} = 96$ $\mu\text{s} = 6\tau_R$ was chosen to maximize one-bond transfers. A $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ DCP spectrum of natural-abundance L-histidine was acquired

at the same spinning rate $\nu_R = 62.5$ kHz, with the same number of 32 scans and 64 increments in t_1 (so that $t_1^{\max} = 64\tau_R \approx 1$ ms), the same recovery delay of 1 s and identical signal processing. The ^{14}N spectrum is shown with blue contours overlaid on top of the grey ^{13}C spectrum. The corresponding projections for the two 2D spectra are shown along the sides using the same color-coding. No ^1H decoupling was applied in the evolution interval t_1 in either spectra. The digital resolution in ω_1 is determined by $(t_1^{\max})^{-1} = (64\tau_R)^{-1} \approx 1$ kHz, which does not allow one to resolve any inhomogeneous features in the ^{14}N line shapes. The more pronounced wiggles observed in the indirect ω_1 dimension of the ^{13}C spectrum are due to digital truncation, indicating longer lifetimes of ^{13}C coherences compared with those of ^{14}N , i.e., $T_2^{\text{SQC}}(^{13}\text{C}) > T_2^{\text{SQC}}(^{14}\text{N})$. We observe a gain in intensity of ca. 3.5 when going from ^{14}N to ^{13}C nuclei, close to the ratio $\gamma(^{13}\text{C})/\gamma(^{14}\text{N}) = 3.48$. This shows that our technique allows us to obtain ^{14}N spectra with

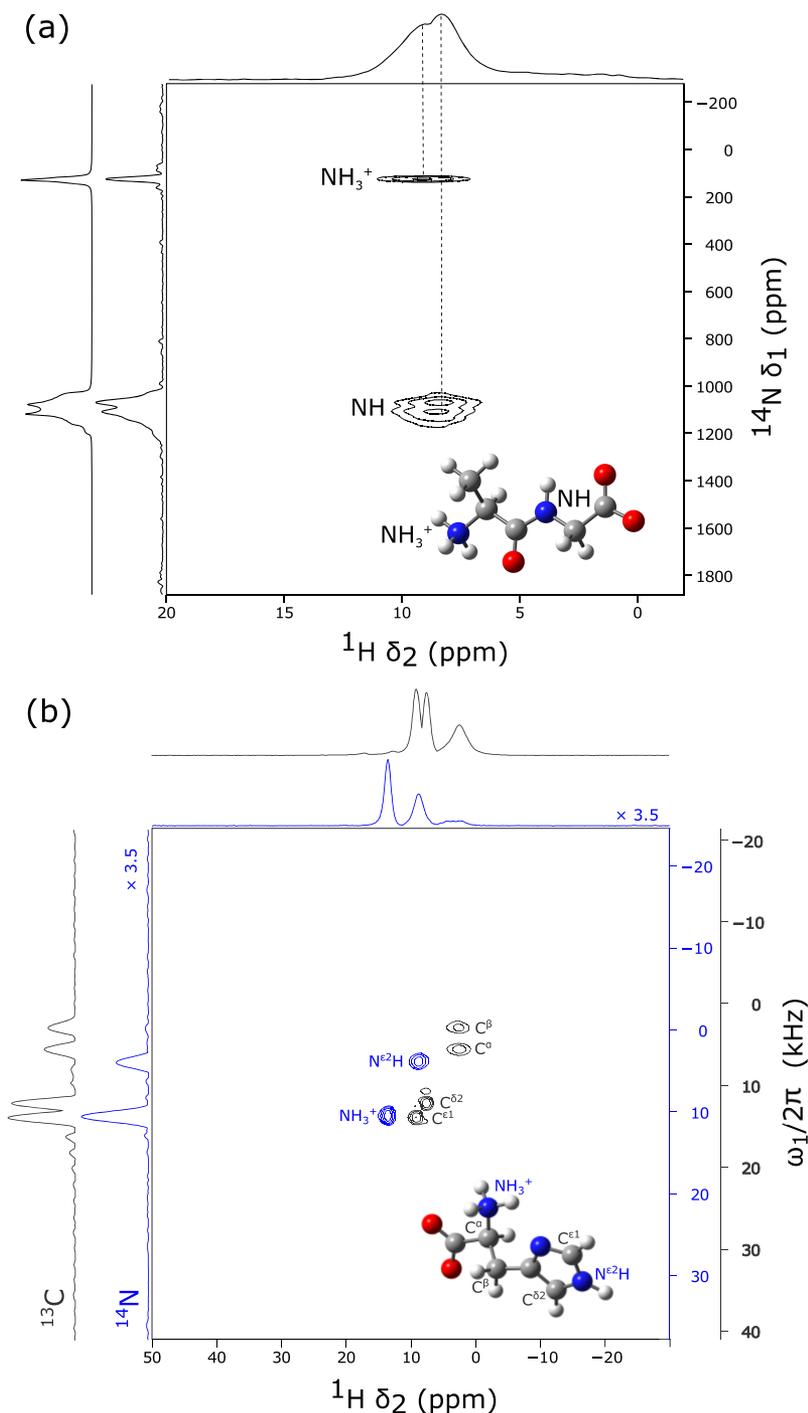


FIG. 3. (a) Experimental ^1H - ^{14}N single-quantum correlation spectrum of the dipeptide Ala-Gly at $B_0 = 9.4$ T recorded with a spinning frequency of $\nu_R = 62.5$ kHz ($\tau_R = 16$ μs) and with a contact time of $\tau_{CP} = 10\tau_R = 160$ μs and 64 transients for each of 160 t_1 increments. Both one-bond correlations of the narrow NH_3^+ and the broad NH sites can be identified. Numerical fits of the ω_1 projections are shown on the left side. (b) Overlay of a ^1H - ^{14}N correlation spectrum of natural abundance L-histidine (blue: $\tau_{CP} = 15\tau_R = 240$ μs) and a ^1H - ^{13}C correlation spectrum of ^{13}C -enriched L-histidine (gray: $\tau_{CP} = 6\tau_R = 96$ μs). The ^1H - ^{14}N and ^1H - ^{13}C spectra in (b) were recorded with exactly the same parameters for acquisition and processing: 32 transients were averaged and 64 t_1 increments were recorded in both cases, in a total experimental time of ca. 34 min. The protonated nitrogen-14 and carbon-13 sites lead to signals with similar intensities.

an efficiency that is comparable to that of ^{13}C in ^{13}C -enriched samples. Of course, larger rotor sizes of, say 4 mm, spun at lower spinning rates of, say, 10 kHz, would surely grant higher sensitivity for 1D ^{13}C CP spectra.

In order to investigate the efficiency of the $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ transfer by double cross polarization, experiments were conducted on polycrystalline powder samples of glycine and L-histidine. One can monitor the CP process as a function of the duration τ_{CP} of the cross-polarization contact pulses and the rf -field strengths, in a manner that is very similar to conventional ^1H - ^{13}C one-way cross polarization. Figure 4(a) shows proton signals observed with the DCP approach in glycine, with $\tau_{CP} = 320$ $\mu\text{s} = n\tau_R$ with $n = 20$, at a spinning

rate of $\nu_R = 62.5$ kHz ($\tau_R = 16$ μs) in steps of $\Delta\tau_{CP} = \tau_R$. The black dots show the build-up of the NH_3^+ proton signals, whereas empty circles show the CH_2 proton signals. The former profile is characterized by a shorter build-up time due to the short one-bond distances involved in the $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ transfers (ca. 1.1 \AA). By contrast, the latter profile shows a slower build-up since the protons involved in the CP process are found at longer distances from the ^{14}N nucleus, i.e., ca. 2.1 \AA if we assume that the transfers occur within the same molecule. Interestingly, an inversion of the sign of the signals is observed for both pathways around $\tau_{CP} = 5\tau_R$. Figure 4(b) shows build up curves similar to those in Fig. 4(a) over a longer time scale $\tau_{CP} = 250\tau_R = 4$ ms in

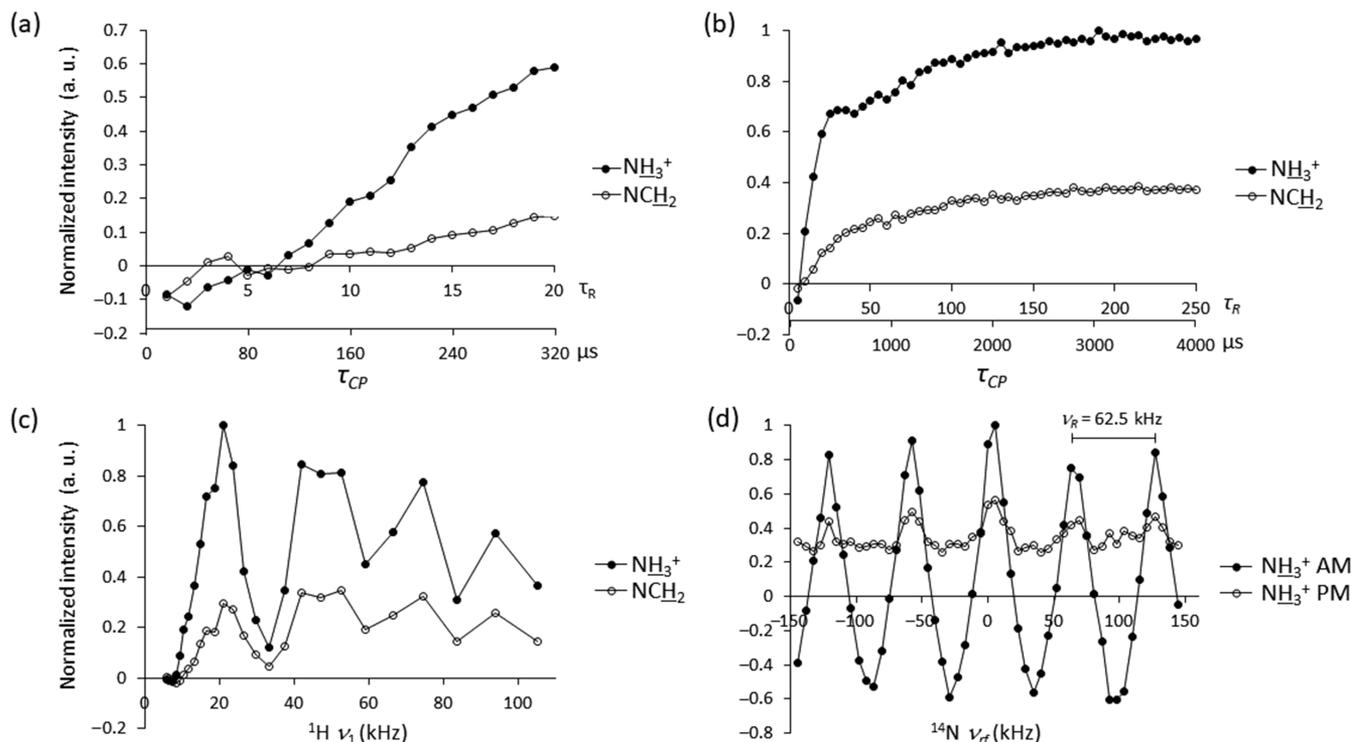


FIG. 4. (a) Experimental signal amplitudes of distinguishable protons of glycine ($\text{NH}_3^+\text{CH}_2\text{CO}_2^-$) observed after a double $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ transfer as a function of the length of the contact pulses $\tau_{\text{CP}} = n\tau_R$, equal to an integer number n of rotor periods $\tau_R = 16 \mu\text{s}$, for $1 \leq n \leq 20$ incremented in steps of 1. Signals detected on the NH_3^+ protons due to one-bond transfers are indicated by black dots, whereas signals detected on the $\text{N}^\alpha\text{CH}_2$ protons that arise from long-range transfers are indicated by empty circles. (b) Similar signals for $5 \leq n \leq 250$, incremented in steps of 5. (c) Similar signals recorded as a function of the rf field strength $\nu_1(^1\text{H})$ applied to the proton channel during a contact time $\tau_{\text{CP}} = 225\tau_R$. (d) Similar signals recorded as a function of the nitrogen-14 carrier frequency $\nu_{\text{rf}}(^{14}\text{N})$ during cross polarization. Both amplitude-modulated (AM) and phase-modulated (PM) experiments are shown.

steps of $\Delta\tau_{\text{CP}} = 5\tau_R$. Filled and empty circles have the same meaning as in Fig. 1(a). The build-up curve of the NH_3^+ signal reaches a plateau near $25\tau_R < \tau_{\text{CP}} < 30\tau_R$. Subsequently, a second surge of signal intensity attributed to more remote protons starts to contribute. This second component is similar to the build-up curve of the CH_2 protons represented by the empty circles.

The conditions for matching the rf -amplitudes $\nu_1(^1\text{H})$ and $\nu_1(^{14}\text{N})$ for the best $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ transfer were also investigated. Figure 4(c) shows signal intensities in one-dimensional DCP experiments carried out on glycine by varying the proton rf -field strength $\nu_1(^1\text{H})$ while keeping $\nu_1(^{14}\text{N}) = 83 \text{ kHz}$. The length of the two contact intervals was $\tau_{\text{CP}} = 225\tau_R = 3.6 \text{ ms}$. Filled and empty circles refer to signals detected on the NH_3^+ and CH_2 protons, respectively, after transfer via ^{14}N , selected by phase cycling. The rf -matching profiles are very similar, with optimal transfer conditions appearing at intervals of ca. 23 kHz.

The robustness of DCP experiments with respect to the position of the ^{14}N carrier frequency $\nu_{\text{rf}}(^{14}\text{N})$, which was varied within the manifold of spinning sidebands that is determined by first-order quadrupolar couplings, was also investigated on the same glycine sample. Figure 4(d) shows signal intensities for DCP experiments performed by varying the carrier $\nu_{\text{rf}}(^{14}\text{N})$ over a range of 289 kHz (ca. 10 000 ppm for a ^{14}N Larmor frequency of 28.9 MHz at $B_0 = 9.4 \text{ T}$). This range of 289 kHz amounts to ca. $4.6\nu_R$, i.e., 4.6 times the spectral width in the indirect dimension, which is identical to the

spinning frequency $\nu_R = 62.5 \text{ kHz}$ since the experiments are rotor-synchronized. Only the NH_3^+ signals are shown for simplicity, the profiles of the CH_2 protons being almost identical. Filled circles show modulations of the signal intensity with a periodicity that matches the spinning frequency $\nu_R = 62.5 \text{ kHz}$, passing from positive to negative values. As remarked by O'Dell *et al.*,⁴⁸ this latter feature is related to the amplitude modulation (AM) of the signal if the evolution interval chosen in the one-dimensional DCP experiments is fixed (i.e., $t_1 = \tau_R$ in the present case), as determined by the selection of both $p(^{14}\text{N}) = \pm 1$ coherence orders. In such 1D experiments, the signal may be written as $S_{\text{AM}}(t_1, t_2) = (1/2) \cos(\Omega_S t_1) \exp(i\Omega_I t_2)$, where the fixed evolution interval is $t_1 = \tau_R$. Since the offset in the ω_1 dimension is $\Omega_S = \omega_{0,S} + 2\pi\nu_{\text{rf}}(^{14}\text{N})$, stepping the carrier frequency $\nu_{\text{rf}}(^{14}\text{N})$ leads to the observed amplitude modulation. If one utilizes a phase-modulated (PM) version of the experiment by selecting only $p(^{14}\text{N}) = -1$ coherences in t_1 , one obtains a nearly flat response, as shown by the empty circles in Fig. 4(d). In this case, one has $S_{\text{PM}}(t_1, t_2) = (1/4) \exp(i\Omega_S t_1) \exp(i\Omega_I t_2)$, so that stepping the carrier frequency $\nu_{\text{rf}}(^{14}\text{N})$ introduces a phase modulation. This can be compensated by a simple first-order phase correction after Fourier transformation. The two-fold reduction in amplitude observed when switching from amplitude- to phase-modulated experiments is because in the latter case half of the signal is discarded by the phase cycle. A PM version of the experiment may be preferred for the optimization of experimental parameters in 1D experiments but is not suitable *as is* for 2D

spectroscopy since it results in phase-twisted 2D line shapes. The nearly flat responses prove the robustness of the DCP transfer for all nitrogen-14 sites that occur in the detected window, provided the carrier frequency ν_{rf} (^{14}N) is positioned within the inhomogeneous line-width that is due to the first-order quadrupolar interactions in the static sample or within the manifold of sidebands in a spinning sample. This requirement can be easily fulfilled since the first-order quadrupolar broadening typically spans several MHz. Analogous responses acquired for L-histidine are substantially identical. Numerical data of these plots are given in the [supplementary material](#).

We compared the performance of our DCP approach with the conventional HMQC method. Both experiments were optimized on a sample of glycine and subsequently applied to samples of L-histidine and the dipeptide Ala-Gly without re-optimizing the parameters (see the [supplementary material](#)). These results indicate that our DCP approach is less demanding in terms of the adjustment of the parameters since it performed well on all samples considered without the need for sample-specific optimizations. Furthermore, when a sufficient number of t_1 increments are acquired, the resulting 2nd-order features of the ^{14}N line shapes are better resolved in the DCP experiment than those produced by HMQC, with FWHH of ca. 1.3 and 1.9 kHz, respectively, obtained on the glycine sample (see the [supplementary material](#)). As discussed in the Introduction, we ascribe this latter finding to homogeneous losses due to transverse relaxation $T_2(^1\text{H})$ and to proton-proton dipole-dipole interactions that affect product states such as $T_{1\pm 1}^I T_{1\pm 1}^S$ and $T_{1\pm 1}^I T_{2\pm 1}^S$. Obviously, a refocusing pulse applied in the center of the t_1 evolution interval fails to cancel the effects of dipole-dipole proton-proton interactions, which may be regarded as the Achilles' heel of all HMQC methods. These interactions evidently cannot contribute to the decay of single-spin states such as $T_{1\pm 1}^S$.

We have attempted to tackle non-crystalline systems such as a lyophilized sample of lysozyme with disappointing results. This suggests that the distribution of molecular conformations and environments in amorphous solids—or a high number of crystallographically distinct sites in ordered materials—may result in excessive broadening of ^{14}N line shapes. In this respect, it is important to consider (i) the isotropic shifts due to second-order quadrupolar couplings that contribute to the spread of resonance frequencies if there is a range of C_Q values and (ii) the fact that one must necessarily observe transitions $m_I = \pm 1 \leftrightarrow m_I = 0$ that are broadened by the first-order quadrupolar interaction, so that the width in the rotor-synchronized indirect dimension is limited by the spinning rate and may result in extensive folding and overlap. A computational investigation of these limitations is presented in the [supplementary material](#), where we show how inhomogeneities in local environments result in a spread of resonances, which is relatively small for the isotropic chemical shift but much more severe for the quadrupolar interactions. These simulations show why amorphous samples may tend to give rise to featureless spectra because the ^{14}N quadrupolar interactions may vary dramatically with the local environment.

All build-up plots in Fig. 4 show signal amplitudes observed in one-dimensional experiments resulting from the averaging of 32 transients. Considering the small amounts of

ca. 3 mg that can be packed in 1.3 mm rotors, the relatively small number of scans required for DCP experiments indicates that indirect detection of ^{14}N via cross polarization at high spinning rates may be accomplished with an efficiency comparable to that of common direct one-dimensional ^{13}C detection with proton decoupling in ^{13}C -enriched samples. Furthermore, the only parameters that need to be adjusted experimentally are exactly the same as one has to optimize for common ^{13}C CP MAS spectra, i.e., the contact time τ_{CP} and rf -field strength ν_1 of one of the two channels. These features indicate that our approach grants both maximal sensitivity and ease of use. Furthermore, the pulse sequence does not have to be symmetric about the center of the t_1 delay. Thus, a longer contact may be chosen for the first $^1\text{H} \rightarrow ^{14}\text{N}$ transfer to exploit the full proton bath as a source of polarization, and a shorter contact may be used after the delay to favor one-bond transfers. By inspection of the build-up profiles of Figs. 4(a) and 4(b) and Figs. 1(a) and 1(b) of the [supplementary material](#), one may conclude that different contact lengths for the first and second CP steps may result in improvements or deteriorations in terms of signal intensity, which are site- as well as sample-specific. Rotor-synchronization of the length contact pulses is not required and their deliberate desynchronization has little effect on the signal intensity.

CONCLUSIONS

A simple double cross polarization (DCP) experiment provides an efficient approach for the indirect detection of nitrogen-14 NMR spectra under fast magic-angle spinning conditions on polycrystalline samples such as glycine, L-histidine, and the dipeptide Ala-Gly. The resulting two-dimensional correlation spectra can be tuned to unveil either short-range one-bond or long-range correlations. In the latter case, unprotonated ^{14}N sites can be detected. The method is shown experimentally to be effective over a range of quadrupolar coupling constants of at least 1-4 MHz. Amorphous samples tend to give rise to very broad ^{14}N spectra because the quadrupolar interactions are very sensitive to the lattice environment. The method is straightforward to implement since it only requires a few simple optimizations, much like standard CP methods that are used routinely for ^{13}C and other nuclei with $S = 1/2$. The efficiency of the polarization transfer has been investigated as a function of the length of the contact pulses, the rf -field strength, and rf -carrier frequency. The efficiency of the transfer of polarization and the sensitivity in our ^1H - ^{14}N experiments is comparable to similar ^1H - ^{13}C experiments of samples that are isotopically enriched in carbon-13.

EXPERIMENTAL DETAILS

All DCP MAS spectra were recorded at the Ecole Normale Supérieure de Paris, on a wide-bore Bruker 400 spectrometer (9.4 T) with an Avance-III console, using 1.3 mm ZrO_2 rotors spinning at $\nu_R = 62.5$ kHz in a triple-resonance MAS probe in a two-channel configuration. Typical amounts of samples packed in the rotors were ca. 3 mg (for L-histidine). The rf -field amplitude for proton excitation pulses was $\nu_1 = 178$ kHz, corresponding to 90° flip-angle pulse length of $\tau_p = 1.4$ μs .

The rf -field strengths for $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ cross polarization steps were $\nu_1(^1\text{H}) = 47$ and $\nu_1(^{14}\text{N}) = 83$ kHz, whereas those for $^1\text{H} \rightarrow ^{13}\text{C} \rightarrow ^1\text{H}$ were $\nu_1(^1\text{H}) = 66$ and $\nu_1(^{13}\text{C}) = 50$ kHz. A recycling delay of 1 s was used in all cases. Phase cycling for the DCP experiment: $\phi_1 = 4\{y\}$, $4\{-y\}$, $4\{y\}$, $4\{-y\}$, $\phi_2 = 8\{x\}$, $8\{-x\}$, $\phi_3 = 4\{x, -x\}$, $4\{-x, x\}$, $\phi_4 = 2\{x\}$, $2\{-x\}$, $\phi_5 = x$, $\phi_{\text{rec}} = x, -x, -x, x, -x, x, x, -x, -x, x, x, -x, x, -x, -x, x$ (see the [supplementary material](#)). The amplitude of the pulses was kept constant during all CP steps. No ^1H decoupling was applied during the t_1 interval (see the [supplementary material](#)). The chemical shifts were referenced with respect to adamantane and NH_4Cl whereas the magic angle was adjusted using KBr.

SUPPLEMENTARY MATERIAL

See [supplementary material](#) for details and discussion about profiles analogous to those of Fig. 4 recorded on a sample of L-histidine, comparison between linewidths produced by the DCP and HMQC approaches, the effect of distribution of environments in terms of chemical shift and quadrupolar interaction, the effect of proton decoupling during t_1 evolution, spurious signal eliminated by the adopted phase cycling, and considerations on the experimental times.

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