

The unit of time: Present and future directions Sebastien Bize

▶ To cite this version:

Sebastien Bize. The unit of time: Present and future directions. Comptes Rendus. Physique, 2019, 20 (1-2), pp.153-168. 10.1016/j.crhy.2019.02.002 . hal-02164954

HAL Id: hal-02164954 https://hal.sorbonne-universite.fr/hal-02164954v1

Submitted on 25 Jun 2019

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés. Contents lists available at ScienceDirect

Comptes Rendus Physique

www.sciencedirect.com

The new International System of Units / Le nouveau Système international d'unités

The unit of time: Present and future directions

L'unité de temps : directions actuelles et futures

Sébastien Bize

SYRTE, Observatoire de Paris, Université PSL, CNRS, Sorbonne Université, LNE, 61, avenue de l'Observatoire, 75014 Paris, France

A R T I C L E I N F O

Article history: Available online 26 February 2019

Keywords: Atomic fountain Timescale Optical frequency standard Fundamental physics test Chronometric geodesy Redefinition of the SI second

Mots-clés : Fontaine atomique Échelle de temps Étalon de fréquence optique Test de physique fondamentale Géodésie chronométrique Redéfinition de la seconde du SI

ABSTRACT

Some 50 years ago, physicists, and after them the entire world, started to found their time reference on atomic properties instead of motions of the Earth that have been in use since the origin. Far from being an arrival point, this decision marked the beginning of an adventure characterized by an improvement by 6 orders of magnitude in the uncertainty of realization of atomic frequency and time references. Ever-progressing atomic frequency standards and time references derived from them are key resources for science and for society. We will describe how the unit of time is realized with a fractional accuracy approaching 10^{-16} and how it is delivered to users via the elaboration of the international atomic time. We will describe the tremendous progress of optical frequency standards with fractional uncertainties of 10^{-18} . We will describe existing and emerging applications of atomic frequency standards. We will describe existing and emerging applications of atomic frequency standards in science.

© 2019 Académie des sciences. Published by Elsevier Masson SAS. This is an open access article under the CC BY-NC-ND licenses (http://creativecommons.org/licenses/by-nc-nd/4.0/).

RÉSUMÉ

Il y a une cinquantaine d'années, les physiciens, et après eux le monde entier, ont commencé à fonder leur référence temporelle sur les propriétés atomiques au lieu des mouvements de la Terre, qui étaient utilisés depuis l'origine. Loin d'être un point d'arrivée, cette décision a marqué le début d'une aventure caractérisée par une amélioration par six ordres de grandeur de l'incertitude de la réalisation des références atomiques de fréquence et de temps. Les étalons de fréquence atomique en constante progression et les références de temps qui en découlent sont des ressources essentielles pour la science et pour la société. Nous décrirons comment l'unité de temps est réalisée avec une précision relative approchant 10^{-16} et comment elle est mise à la disposition des utilisateurs via l'élaboration du temps atomique international. Nous montrerons les progrès considérables de la métrologie des fréquences optiques au cours des vingt dernières années, qui ont conduit à une nouvelle génération d'étalons de fréquence optique avec des incertitudes relatives de 10^{-18} . Nous décrirons les travaux en vue d'une éventuelle redéfinition de la

E-mail address: sebastien.bize@obspm.fr.

https://doi.org/10.1016/j.crhy.2019.02.002

1631-0705/© 2019 Académie des sciences. Published by Elsevier Masson SAS. This is an open access article under the CC BY-NC-ND license (http://creativecommons.org/licenses/by-nc-nd/4.0/).







seconde du SI basée sur ces étalons. Nous décrirons les applications scientifiques existantes et émergentes des étalons atomiques de fréquence.

© 2019 Académie des sciences. Published by Elsevier Masson SAS. This is an open access article under the CC BY-NC-ND licenses (http://creativecommons.org/licenses/by-nc-nd/4.0/).

1. The unit of time

In 1967, the 13th General Conference on Weights and Measures (CGPM) defined the SI second as "the duration of 9 192 631 770 periods of the radiation corresponding to the transition between two hyperfine levels of the ground state of the cesium 133 atom" [1]. Atoms have quantized energy levels. With any pair of quantized levels of energy E_g and E_e , a frequency ν is associated via Planck–Einstein's relation: $h\nu = E_e - E_g$. One fundamental idea behind the definition is that these atomic frequencies are perfectly stable and universal. Observations support this idea. Atoms can be regarded as "perfect" frequency standards given by nature. The immutability of atomic frequencies is integrated into fundamental theories underpinning physics: general relativity and the standard model of particle physics.

Accessing the atomic frequency and transferring its qualities to a macroscopic usable signal requires a device called atomic frequency standard. The typical architecture of such device is shown and explained in Fig. 1. The output frequency v(t) does not coincide with the atomic frequency v_{at} . It is perturbed by noises and biases of both technical and fundamental nature. These noises and biases determine two mains characteristics of a given atomic frequency standard. Noises limit the uncertainty with which a frequency measurement can be made with this standard in a given duration. It is characterized by the fractional frequency instability $\sigma_y(\tau)$, which is a function of the measurement duration τ . Biases offset the mean output frequency with respect to the unperturbed atomic frequency v_{at} . If known and stable, a bias can be taken into account. What really matters is the uncertainty to which biases are known. This uncertainty, often reported in fractional terms, defines the level to which the standard actually gives access to the unperturbed atomic frequency. It summarizes the capability of the standard to realize the SI second (for ¹³³Cs primary standards), to be used for fundamental physics and for other applications.

In practice, only well-chosen atomic transitions are suitable to realize standards with lowest stability and uncertainty. Levels must have long lifetimes to enable a high atomic quality factor Q_{at} (see Fig. 1). Transition must have low sensitivity to external fields (electric, magnetic, thermal radiation, etc.). The atomic structure must be compatible with methods needed to manipulate and detect atoms. It is also essential to consider practical criteria such as reliability, operability, possibility to generate and use the interrogation field at the transition frequency. Given the state of knowledge and technology at the time of the 13th CGPM, a frequency standard based on the ¹³³Cs ground-state hyperfine transition was one of the best possibilities, almost two decades after the first observation of the transition [2]. It had shown sufficient maturity and had been accurately measured with respect to the ephemeris time [3,4]. Since then, fundamental and technological breakthroughs in many areas led to major changes in ways to realize and use highly accurate atomic frequency standards. Their uncertainty improved by 6 orders of magnitude. References [5] and [6] provide an overview of these developments.

In November 2018, the 26th CGPM adopted a major redefinition of the international system of units. The redefinition concerns the units of mass, of temperature, of electrical current and of amount of substance. The new system is defined by adopting conventional values for Planck's constant h, for Boltzmann's constant k_B , for the elementary charge e and for Avogadro's constant N_A , similarly to what was already done since 1983 [7] with the speed of light c to define the unit of length. In essence, the definition of the second remains unchanged. The form of the definition, however, is modified to resemble those of other units, i.e. "the International System of Units, the SI, is the system of units in which: the unperturbed ground state hyperfine transition frequency of the caesium 133 atom Δv_{CS} is 9 192 631 770 Hz [...]" [8].

2. Research on highly accurate atomic frequency standards

Section 1 above gave fundamental concepts behind atomic time and atomic frequency standards. These concepts are amazingly simple and have been used for more than five decades. Still, they continue being at the basis of an active field of research in the following ways.

Search for advancing atomic frequency standards with extremely low uncertainties goes hand in hand with exploration of atomic systems interacting with electromagnetic fields. Progress in atomic frequency standards reveals new phenomena and provides means of investigating them. And vice versa, progress in other areas (e.g., laser-cooling of atoms, physics of collisions and interactions, quantum entanglement...) enables the improvement of atomic frequency standards.

These devices with extremely low uncertainties are tools of choice to probe the structure of space-time and to test fundamental laws of nature. They provide experimental inputs to the quest for a unified theory of gravitation and quantum mechanics [9–11] and to the search for dark matter [12,13] and to other applications.

On the applied side, atomic frequency standards are key to major applications of utmost importance in modern science and society. The practical realization of the unit of time of the SI system is one of them, in particular via the elaboration of the international atomic time and of the universal coordinate time (TAI/UTC). Another example is global navigation satellite

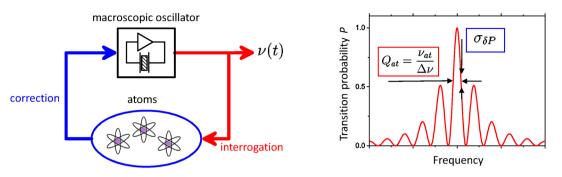


Fig. 1. On the left: architecture of highly accurate atomic frequency standards. An oscillator generates a macroscopic, practically usable electromagnetic signal. A fraction of this signal is split to probe the atomic transition chosen as a reference. The response of the atomic sample is detected to determine the probability of exciting the transition. This information is used to stabilize the frequency of the oscillator to the spectroscopic signal. An other right: representative spectroscopic signal. Key features of this signal are the atomic quality factor $Q_{at} = v_{at}/\Delta v$ and the noise in the measurement of the transition probability $\sigma_{\delta P}$. The frequency v_{at} : $v(t) = v_{at} \times (1 + \varepsilon + y(t))$, where ε is a fractional frequency offset and y(t) represents fractional frequency fluctuations.

systems (GNSS) such as GPS, GALILEO, etc. In turn, these realizations are key for telecommunications, transports, finance, digital economy, etc.

Search for ultra-low uncertainties in atomic frequency standards is a steady driver for innovation in multiple technological areas such as lasers, low-noise electronics, ultra-stable oscillators. Also, it creates knowledge necessary to define trade-offs between performance and other requirements for industrial frequency standards or space clocks. This knowledge is often applicable not only to novel frequency standards, but also to other types of atom-based instruments like, for example, accelerometers and gyrometers based on matter–wave interferometry or magnetometers.

Progress in reducing the uncertainty of atomic frequency standards leads to novel applications. For instance, because of its much lower uncertainty, the new generation of optical frequency standards enables chronometric geodesy, i.e. the determination of differences in terrestrial gravitational potential via the measurement of Einstein's gravitational red shift.

3. Primary frequency standards based on atomic fountains

The first generation of 1^{33} Cs primary frequency standards, which led to the adoption of the atomic time, was based on the thermal atomic beam technology and the separated oscillatory fields method [14,15]. The development of laser cooling of atoms in the 1980s [16–18] enables a second generation called atomic fountains [19–21]. Atomic samples laser-cooled to temperatures near 1 μ K enable atomic quality factors higher than 10¹⁰, a factor 100 times higher than in atomic beam standards. Fig. 2 displays a schematic of an atomic fountain. When an ultra-low noise microwave source is used to interrogate the atomic transition, the quantum projection noise limit is reached [22]. This limit is the fundamental limit set by the quantum measurement process for unentangled particles. This yields short-term fractional frequency instabilities as low as 1.6×10^{-14} at 1 s [23]. Nowadays, the accuracy of the best atomic fountains ranges between 1 and 3 parts in 10^{16} . This is the result of refining models of the interrogation and detection processes, and of stringently testing these models with experiments. Such experiments typically require many frequency measurements with statistical uncertainties near 10^{-16} . which is reached, in the very best case, after several full days of measurement. In other words, fountains have reached the situation where the stability imposes a practical limit to studying systematic shifts and improving the accuracy. Example of effects whose modeling improved significantly includes the effects of phase gradients in the interrogation microwave cavity [24,25], the effects of cold collisions [26–28], the effects of collisions with background gases [29,30], and the effects of the microwave field on external atomic motion ("microwave lensing") [31–33]. The implementation of fountains with cryogenic interrogation environment enabled a new direct measurement of the effects of thermal radiation shift ("blackbody radiation shift") [34–37]. Recently, a complete accuracy budget was published for continuous fountain at the level of 1.99×10^{-15} [38]. More details on atomic fountains can be found, for instance, in [39], [23] and [40].

The ⁸⁷Rb ground-state hyperfine transition is also used to realize atomic fountain frequency standards. LNE–SYRTE developed a dual fountain using Rb and Cs, which truly realizes two state-of-the-art microwave standards in a single setup [41,42]. The Rb part of this fountain has an uncertainty of 3.2×10^{-16} similar to the one of the best cesium fountains. This dual fountain enables highly accurate comparisons of the ⁸⁷Rb and ¹³³Cs hyperfine frequencies. These comparisons find many applications (see sections 6 and 7 below). In particular, they led to the adoption of this transition as a secondary representation of the SI second [42,43].

Agreement between fountains is well tested by means of specific remote comparisons by satellite methods [44] and recently by optical links between LNE–SYRTE and PTB [45]. Also, the elaboration of TAI provides a vehicle to compare fountain frequency standards (see section 4 below).

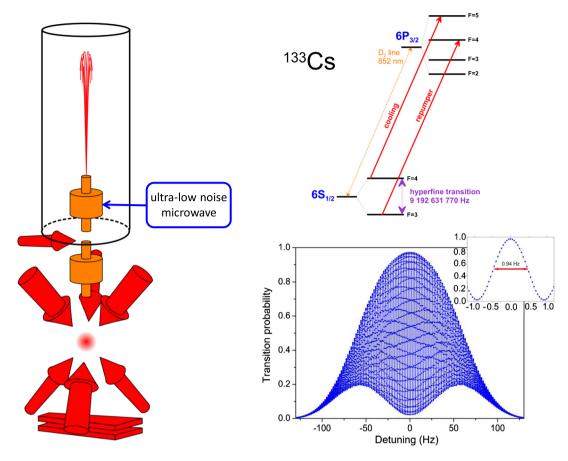


Fig. 2. On the left: Schematic of an atomic fountain frequency standard. A cloud of cold atoms is captured at the crossing of 3 pairs of counter-propagating laser beams. It is launched upwards at a typical speed of 4 m s⁻¹ and with temperature of ~ 1 μ K. During the ballistic flight, atoms are state-selected with a first microwave interaction and a push laser beam. Then, they pass upward and downward through a microwave cavity where they interact with the signal from the interrogation oscillator. They continue falling through two laser beam sheets in the detection region. On the right: Top: Energy levels of the cesium hyperfine transition in an atomic fountain showing Ramsey fringes with a width < 1 Hz and an atomic quality factor Q_{at} of 10¹⁰. Each point is a single measurement of the transition probability at a rate of ~ 1 per second with a typical noise $\sigma_{\delta P} \sim 2 \times 10^{-4}$.

4. Elaboration of TAI: accuracy of atomic fountains delivered to users

One essential outcome of time and frequency metrology is the construction of the international atomic time (TAI) and of the universal coordinate time (UTC). It is worth noting here that the 26th CGPM adopted a resolution on the definition of time scales [8], which corrects for the lack so far of a self-contained definition of TAI. We remind that a meaningful definition and realization of a time scale valid globally in the vicinity of the Earth requires the framework of general relativity, in particular to properly account for Einstein's gravitational redshift, which is about 10^{-16} per meter of elevation at the surface of the Earth. A description of the elaboration of TAI by the BIPM can be found, for example, in [46] and references therein. By means of satellite-based comparisons, data from about 450 continuously operated commercial clocks are used to compute the free atomic time scale (EAL). The large number of devices and locations ensure the permanence of EAL. Data from a much smaller number of primary and secondary frequency standards (about 20) in an even more limited number of institutes are used to calibrate the frequency of EAL against the ¹³³Cs hyperfine transition, according to the definition of the second, and to steer EAL to realize TAI. UTC is derived from TAI by inserting leap seconds that maintain agreement with the universal time (UT1) and the observed rotation of the Earth. The result of this process is published by the BIPM in its monthly *Circular T* [47]. This information, in turn, is used by national metrology institutes and other participants to steer their local physical representation of UTC, from which disseminations in society start by various means.

Circular T and calculations done at the BIPM also yield an estimation of the performance of TAI and provide a vehicle to compare frequency standards worldwide. Differences between frequency standards observed by this means are consistent with uncertainties [48,49]. The accuracy to which the scale interval is determined with respect to ¹³³Cs hyperfine transition now reaches 2×10^{-16} [46]. In other words, the performance of frequency standards is transferred to the time scale and thereby to users. The 2×10^{-16} fractional frequency uncertainty translates into an error of less than 10 ns after one year. This improvement by a factor of 10 since 2000 is the result of the commitment of a few metrology institutes to provide

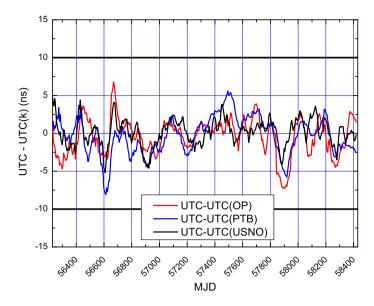


Fig. 3. Local realizations of UTC using atomic fountain data. The plot shows time differences UTC-UTC(k) as a function of the Modified Julian Date (MJD) for UTC(USNO), UTC(PTB) and UTC(OP). Data are shown between 29 October 2012 and 8 November 2018.

regular calibrations of TAI. Nowadays, about four or five calibrations by fountains are typically available each month, as can be seen in section 3 of *Circular T* [47]. Over the last 15 years, LNE–SYRTE made 40% of all worldwide calibrations with fountains. The adoption of secondary representations of the SI second (see [43] and section 7 below) led to the possibility to calibrate TAI with other atomic transitions than the ¹³³Cs hyperfine transition. This was done for the first time with the ⁸⁷Rb hyperfine transition by LNE–SYRTE, which provided close to 100 calibrations by this mean [42]. LNE–SYRTE also pioneered providing calibrations based on an optical transition. This was done with the ⁸⁷Sr¹S₀–³P₀ transition at 698 nm. Regular calibrations of TAI by optical frequency standards are an important prerequisite for a possible redefinition of the second based on optical transition(s) [50]. The transfer of long-term stability and accuracy of primary frequency standards to TAI enables highly accurate SI-traceable frequency measurements without a local primary frequency [51,52] and fundamental physics tests [53,54].

Progress in the reliability of atomic fountains made it possible for a few institutes to steer their local realizations of UTC with fountain data [55–57]. Local realization of UTC by institute k is denoted UTC(k). Fig. 3 shows time differences between such timescales and UTC. They are realized using hydrogen masers whose frequency is calibrated and steered with atomic fountains, and for the long term, with the time difference UTC(k)–UTC provided every month in *Circular T*. They typically deviate from UTC and from each others by no more than a few nanoseconds. It is worth noting another significant evolution in timekeeping, namely the rapid realization of UTC by the BIPM [58]. UTCr implements faster exchange of data than *Circular T* does between participating laboratories and the BIPM. It improves the level to which a given laboratory can verify the synchronization of its UTC(k) to UTC. This is of particular interest and significance for laboratories where resources allocated to the local timescale are limited (i.e. to a few commercial Cs beam standards).

5. A new generation of optical atomic frequency standards

Optical frequency metrology is advancing at high pace, in particular since the introduction of optical frequency combs [59,60]. Optical frequency standards refer to atomic standards relying on transitions whose frequency corresponds to the optical domain of the electromagnetic spectrum. Their frequency is by 10^4 to 10^5 higher than the 133 Cs hyperfine frequency and their potential atomic quality factor can exceed 10^{16} instead of 10^{10} for atomic fountain standards. To date, both neutral atoms and ions are studied with the aim to obtain the lowest possible uncertainties. Fig. 4 explains the principle of operation of these two types of standards. Optical frequency standards achieve fractional frequency uncertainties close to 10^{-18} [61–68]. For example, reference [61] reports an uncertainty of 1.4×10^{-18} for an 171 Yb optical lattice standard. References [64] and [62] report uncertainties of 2.1×10^{-18} and 4.8×10^{-18} , respectively, for 87 Sr optical lattice standards. Reference [63] reports an uncertainty of 3.2×10^{-18} for a 171 Yb⁺ ion standard. Many physical effects had to be understood and controlled at an even lower level individually. Some notable ones are frequency shifts due to black-body radiation [64–71] and to electric fields [72], shifts induced by the lattice light [73,74], shifts induced by interactions [75], light-shifts induced by probe light [76]. This list of examples must not be considered as complete either in terms of effects or in terms of references. Comparisons between optical frequency standards using the same transition were performed with improving uncertainties [65,77–80], down to below 10^{-18} for the most recent work [61]. A recent and more complete account on optical frequency standards can be found in [81]. Fig. 5 shows a recent measurement of the stability of a 87 Sr optical lattice frequency standard from LNE–SYRTE.

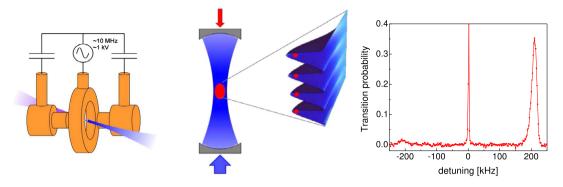


Fig. 4. On the left: Schematic of single ion optical frequency standard. The ion is confined in a Paul trap made of electrodes with oscillating electric potentials. The ion is laser-cooled and probed with ultra-stable laser light. Center: Schematic of optical lattice neutral atom frequency standard. An ensemble of atoms is confined in a corrugated potential formed by an intense standing laser field (blue). Atoms are probed with ultra-stable light aligned with the trap axis (red). On the right: Spectroscopy of the ⁸⁷Sr reference transition at 429 THz in an optical lattice. The spectrum is typical of Lamb–Dicke regime used in optical frequency standards. Because of the confinement, the external motion of atoms has quantized vibrational levels. The spectrum exhibits resolved sidebands with frequency offsets corresponding to the trap vibrational frequency. The central "carrier" resonance corresponds to excitations that do not change the external atomic state. It is essentially unaffected by the first order Doppler effect and by the recoil effect [82]. This resonance can be as narrow as few 100 mHz corresponding to atomic quality factors Q_{at} of several 10¹⁵, at the origin of the superior performance of optical frequency standards. The counterpart of using confined particles is the need to care about effects of the trapping fields. Fundamentally, the ion is held thanks to its electric charge without strong perturbation of its internal structure. It is nevertheless necessary to care about effects of micro-motions induced by the oscillating electric field [83,84]. Trapping of neutral atoms relies on polarizing them with the intense lattice light and thereby on perturbing their internal structure. In optical frequency standards, the lattice trap can be non-perturbing if the trap wavelength has a specific wavelength called *magic wavelength* where the polarizability is the same for the two levels of the reference transition [85–87].

Evaluation of uncertainties and comparisons at the 10^{-18} level are made possible by the excellent short-term stabilities reached by optical frequency standards [88–92]. Stabilities as low as 1.6×10^{-16} at 1 s can be observed for a single optical lattice frequency standard and even lower between two regions of the same atomic cloud [93]. Ion-based optical frequency standards show significantly worse stabilities because they are so far using a single ion instead of several hundreds of atoms or more in neutral atom standards. Synchronized and correlated interrogation of two ion-based standards enables comparison with stability in the mid- 10^{-16} at 1 s [94]. Such stabilities became possible only with the progress of ultra-stable lasers, which remain the subject of active developments. Classical ultra-stable lasers based on Fabry–Pérot prove being limited by thermal Brownian noise in dielectric mirror coatings [95]. Several approaches were and are still investigated to mitigate this limit. Extending cavity length can already give significant improvement [89,96,97]. Crystalline silicon cavity at cryogenic temperatures showed exquisite laser instability and laser linewidth (4×10^{-17} and < 10 mHz, respectively) [98, 99]. Another promising method could be to use crystalline coatings that exhibit lower thermal noise [100]. Other approaches shift away from the Fabry–Pérot cavity. Prospects exist to use spectral hole burning in rare-earth doped ions in crystalline matrices at cryogenic temperatures [101–103]. Lasers using ultra-narrow atomic transitions are another proposed alternative [104–106].

Interest in quantum technologies increased considerably in the last years because they promise major breakthroughs in computation, communication, and sensing [107,108]. Research in optical frequency standards already gave striking examples of quantum enhanced metrology. Operation of the single Al^+ ion frequency standard relies on a quantum gate between Al⁺ and a companion ion used for state readout [77,109,110]. To date, there is still a large potential for optical frequency standards to further exploit quantum metrology. Tailored quantum superposition of internal states can reduce sensitivity to external field perturbations [111]. Entangled states of several ions can improve the stability below the quantum projection noise limit [112]. In parallel, progress were made in designing traps that can support multiple-ion chains while maintaining low uncertainty due to motional effects [113]. A challenge for future ion-based standards will be to merge all methods in a single device. Neutral atom standards already use samples that comprise hundreds or thousands of atoms. Quantum non-destructive measurements performed on such samples can generate entangled states [114] that are metrologically useful like, for example, spin-squeezed states [115]. Proof-of-principle experiments using hyperfine transitions are reported in [116,117] and in [118] where a specular 10-fold reduction (20 dB squeezing) below the quantum projection limit is obtained. Non-destructive detection by optical phase shift measurement in ⁸⁷Sr optical lattice frequency standard is developed at LNE-SYRTE [119,120]. A promising sensitivity level of a few atoms only is achieved and the path toward a quantum nondestructive regime is clarified. Optical lattice frequency standards constitute an excellent platform to harvest the benefit of both classical and quantum non-destructive detection. Classical non-destructive detection could help reduce dead times in the probing sequence and thereby limit the negative influence of laser frequency noise on stability. Atomic phase lock method [121,122] could be used to extend the interrogation duration beyond the probe laser coherence time. Quantum non-destructive detection could be used to beat the quantum projection noise limit. These schemes promise stabilities at 10^{-17} at 1 s or below. It will remain to investigate to which extent non-destructive detection introduces additional sources of uncertainty. Many other schemes to generate entangled states do exist, which are potentially interesting to improve frequency standards, like, for instance, [123–125]. The present account must not be considered complete.

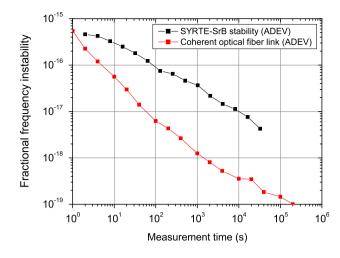


Fig. 5. Stability of optical frequency standards and coherent optical fiber links. In black: fraction frequency instability of a ⁸⁷Sr optical lattice frequency standard from LNE–SYRTE. The curve shows the stability of a single standard (SYRTE–SrB) inferred from the analysis of a multipartite comparison with SYRTE–Sr2, NPL–Sr, and PTB–Yb+. In red: stability of a Paris–Strasbourg–Paris link developed within equipex REFIMEVE+ project, coordinated by LPL.

Optical frequency combs are of utmost importance for optical frequency standards. They enable comparisons between optical frequency references at different wavelengths. They enable conversion of optical frequency references to the microwave domain and thereby, connection with all existing time and frequency methods and infrastructures, in particular comparisons between optical and microwave standards. Optical frequency combs developed and found applications in many other different fields, which are reviewed for example in [126,127]. Here, we highlight developments directly connecting to highly accurate frequency standards. One key aspect was the development of reliable femtosecond lasers based on erbium-doped fiber technology. Combs based on this technology readily connect to the wavelength of 1.5 μ m used in optical fiber links (see below) and in best ultra-stable lasers to date. Such combs can transfer stability between ultra-stable optical references with degradation no higher than 4×10^{-18} at 1 s, a level that surpasses by far the stability of the best ultra-stable lasers [128]. Optical-to-microwave conversion was also developed to minimize noises from all processes involved in the comb architecture [129,130]. Reliable optical frequency comb systems can deal with multiple wavelengths simultaneously and support complex optical and microwave frequency standard comparisons and microwave generation applied to atomic fountains [45,131–138].

Optical fiber links made tremendous progress over the last decade. Coherent optical fiber links transmit an ultra-stable laser light at 1.5 μ m and are therefore directly adapted to the comparison of distant optical frequency standards. From the initial proof-of-concept experiments [139–141], this method was extended to continental distances [142,143]. It enabled comparison of optical frequency standards over continental distances with unprecedented stability and accuracy [137]. Fig. 5 shows the stability of a Paris–Strasbourg–Paris link obtained during two weeks of operation. This link enables comparison to 10^{-18} in less than 2000 s. In France, the first industrial-grade link of this type was recently implemented and tested [144]. It proved the readiness of this technology for commercialization and applications. Coherent optical fiber links reach fractional frequency instability below 10^{-18} for a measurement duration of 10^4 s for lengths of several hundreds to 1500 km. The accuracy of the frequency transfer is verified to better than 10^{-19} . Over continental distances, they surpass satellite-based methods by 3 to 4 orders of magnitude. Besides frequency transfer using a coherent optical carrier, time transfer using optical fibers is actively developed with, here also, the promise to surpass existing methods by orders of magnitude (see, for instance [145–148]). Relativistic effects in optical fiber links are studied from a theoretical standpoint in [149]. In addition to comparisons of optical frequency standards and dissemination of time and frequency references, optical fiber links are used or considered for applications in chronometric geodesy (see section 8 below), in the measurement of the rotation of the Earth [150,151], in earthquake detection [152], and in fundamental physics (see, e.g., [153]).

6. Test of fundamental laws of physics

Physics is underpinned by two fundamental theories: general relativity and the standard model of particle physics. This theoretical framework is extremely successful in describing a huge number of observations. Nonetheless, it is not free of significant problems. One difficulty is that the standard model is a quantum field theory, while general relativity is not. This is giving an heterogeneous picture of the three fundamental interactions. To date, a unified theory of gravitational, electroweak, and strong interactions is still missing. A second difficulty arises in the need to introduce dark matter and dark energy to reproduce rotation curves of galaxies and the accelerating expansion of the Universe in the so-called Λ CDM cosmological model of the Universe. In this model, dark matter and dark energy introduced in an *ad hoc* manner represent 95% of mass and energy in the Universe.

In this context, many experiments and many observations are used to test fundamental laws of physics. One can refer, for instance, to [9–11,154,155] for an overview. Here, we will highlight tests using highly accurate atomic frequency standards. Unlike tests based on observations over geological and cosmological timescales, tests using atomic frequency standards are present-day laboratory experiments that do not require a particular cosmological model for their interpretation. Also, they can be repeated for verification purposes. For these reasons, they are good tools for revealing physics beyond general relativity and the standard model and for providing experimental constraints to guide the development of unified theories (see, for instance, [156–158]).

Within its uncertainty, an atomic frequency standard gives access to the unperturbed atomic frequency of the chosen transition. Comparisons of atomic frequency standards based on different atomic transitions determine atomic frequency ratios. These frequency ratios are dimensionless quantities given by nature, independently of any system of units. A variation of atomic frequency ratios would be a violation of the framework defined by general relativity and the standard model of particle physics and of its founding principles. This can be tested with repeated measurements of atomic frequency ratios. Time series of such measurements are used to search for signals indicative of variations induced by putative phenomena. Linear variations with time could be a present-day effect of evolution of the Universe over the cosmological timescale. Variations synchronous with the Earth's motion around the Sun could be an effect of an extraneous coupling with the gravitational field of the Sun or to additional field(s) finding its source in the Sun. Sinusoidal variations could be due to an extraneous coupling with certain candidate dark matter field condensed on our galaxy. For example, linear variations of the fine-structure constant α , of the electron-to-proton mass ratio $\mu = m_e/m_p$ and of the quark mass m_q/Λ_{QCD} are tested with uncertainties of 1.0×10^{-7} , 8.8×10^{-7} and 2.3×10^{-6} (see, for instance, [11,159] and references therein). Searches for dark matter with frequency standards are reported in [12,13,160]. All these tests improve regularly with the steady progress of atomic frequency standards.

Time and frequency metrology provide many other tests of local position invariance and local Lorentz invariance. New possibilities are frequently emerging, like, for instance, the use of fiber links [153] for special relativity tests, of GALILEO satellites for gravitational redshift test [161]. The ACES space mission [162–165] will provide improved tests and new opportunities. It is worth noting here that the third component of Einstein's equivalence principle, the universality of free-fall, was recently tested with improved uncertainty with the MICROSCOPE space mission (see [166] and references therein).

7. Toward a redefinition of the SI second

Optical frequency standards now surpass ¹³³Cs fountain primary standards by more than 2 orders of magnitude, both in stability and in uncertainty. This naturally creates a strong incentive to redefine the SI second based on optical transition(s). The advent of this possibility was anticipated shortly after the demonstration of the first optical frequency comb. In 2001, the Consultative Committee for Time and Frequency (CCTF) recommended that a list of secondary representations of the second (SRS) be established [167], as one of the key processes to engage into a possible redefinition of the second. In 2004, the International Committee for Weights and Measures (CIPM) adopted the first SRS, the ⁸⁷Rb hyperfine frequency based on developments and measurements made at LNE–SYRTE [168,169]. To date, the list of SRS comprises nine transitions, eight of which are optical transitions [50]. SRS are atomic transitions used to realize frequency standards with excellent uncertainties and which are measured in the SI system with accuracies close to the limit of ¹³³Cs fountain standards. The recommended fractional frequencies of the CIPM [43,50]. The recommended values of SRS come from the work of the joint CCL–CCTF Frequency Standard working group of the CIPM. Values are determined based on the least-square adjustments to high accuracy measurements found in peer-reviewed publications [50,170,171]. This work enables checking the consistency of published measurements and the status of atomic frequency standards with the lowest uncertainties.

In view of a possible redefinition of the second, a key aspect of the process is to link the optical frequency domain to the present definition. In the current status of the list of recommended standard frequencies, this is achieved by some 55 highly accurate absolute frequency measurements of optical transitions. The 698 nm ${}^{87}Sr^{1}S_{0}-{}^{3}P_{0}$ transition used in optical lattice standards gathers by far the largest number of measurements and the ones with the highest accuracy from LNE–SYRTE [78,136] and PTB [172]. These measurements are limited by the accuracy of the best realizations of the second based on the ${}^{133}Cs$ hyperfine transition. Measurements used to establish the list of recommended values also include a still limited number of optical-to-optical frequency ratios (eight measurements of five ratios). Such measurements are not limited by ${}^{133}Cs$ fountain. The best measurements so far have a fractional frequency uncertainty near 5 parts in 10^{17} [173, 174]. Among optical-to-optical frequency ratios, only the ${}^{199}Hg/{}^{87}Sr$ is measured independently in at least two laboratories (RIKEN [175] and LNE–SYRTE [138]) with uncertainties below the current realization of the SI second. Optical-to-optical ratios measured independently provide means to check the status of optical frequency metrology beyond the limit imposed by ${}^{133}Cs$ fountains. Ongoing work shall increase the number of such tests and push their uncertainties near one part in 10^{18} .

A redefinition of the SI second must not have negative consequences on the elaboration of TAI. This implies that the sustainable capability of calibrating TAI scale interval against optical transition(s) must be proven before the redefinition. Most of the current TAI architecture and infrastructure operates in the microwave domain and it will remain so for many more years. Combs enable dividing the frequency of an optical reference to the microwave and RF domains with a limited

degradation of its accuracy. An optical frequency standard can thereby readily be inserted in lieu of ¹³³Cs fountains in the TAI calibration processes of metrology institutes, provided that its reference transition is adopted as a secondary representation of the second [42]. The minimum requirement is to being able to calibrate the mean frequency of a local oscillator linked to TAI (typically, a hydrogen maser) over five days of a conventional grid. This was pioneered by LNE–SYRTE using the ⁸⁷Sr¹S₀–³P₀ transition. A first series of TAI calibrations with this SRS is reported in the BIPM *Circular T350* [47]. These first calibrations were not used to steering the frequency of TAI. At the beginning of 2019, *Circular T372* was published with the first TAI calibrations by optical standards used for steering, still with ⁸⁷Sr. Reliable operation and use of ⁸⁷Sr optical lattice standards to this end is described in [136,176]. Along the same line, application of a ⁸⁷Sr optical lattice standards over intercontinental baseline via satellite-based methods and for almost one day is reported in [177]. In both cases, the architecture of the timescale and links is in the microwave domain. It remains a topic of investigation to define novel architectures and infrastructures that can take full benefit of the 2 orders of magnitude better characteristics of optical frequency standards [178]. Such architectures and infrastructures may include optical fibers links, optical local oscillators, optical clockwork like the one demonstrated in [179], and novel methods.

Altogether, CCTF, in its strategy document, keeps a list of requirements for a redefinition of the second based on optical transition(s) to become possible. The aim is to make a choice that will last long, to ensure continuity, to guaranty gapless dissemination in particular via the elaboration of TAI and to validate the uncertainty of optical frequency standards [50]. CCTF and the joint CCL–CCTF Frequency Standard working group monitor progress of the field in this direction.

8. Chronometric geodesy

Remote comparisons between two identical frequency standards show a frequency ratio $v_2/v_1 = [1 - (U_2 - U_1)/c^2]$, where U_1 and U_2 are gravitational potentials at the location of the two standards. This gravitational redshift already mentioned in section 4 amounts to $\sim 10^{-18}$ per centimeter of elevation in the vicinity of the Earth's surface. Frequency standards can be viewed as sensors to determine gravitational potential differences for the purpose of geodesy and Earth's science. This method, called chronometric geodesy, was proposed several decades ago (see, for example, [180] and references therein). However, it is only now that progress in optical frequency standards and in optical frequency metrology makes this method potentially relevant.

A few years ago, chronometric geodesy gained significant interest from both time and frequency metrology and geodesy communities. Several aspects need being covered. One is to develop methodologies to use data from frequency standards in conjunction with other data already used in gravity field modeling. In this context, optical frequency standards can be considered as sensors of the potential with a point-like spatial response function, since the spread of the atomic sample is less than 1 mm. Instead, space-based gravity measurements (GRACE, GOCE missions) gives first or higher derivatives of the potential with a spatial response function with a characteristic size of 100 km. Gravimeters and gradiometers on ground give the first and second derivatives of the potential, respectively, with a point-like response (0.1 m–1 m), and so on. Geodetic methods to link the potential at the location of a frequency standard to larger and global scale are described in [137,181,182]. These methods were applied to several places in Europe and provided improvements over previous similar determinations [183–185]. Specific gravimetry and leveling measurements in the vicinity (up to few 10 km) of the location of the frequency standards are required to connect to the regional and global scale to the 10^{-18} or equivalently to the 1 cm level. Time variations must be taken into account, since tides can produce local gravity potential changes of up to 10^{-16} . This is described in [186].

Another aspect of the investigations of chronometric geodesy is to understand and define cases where data from optical frequency standards can bring the most to applications in geodesy and geophysics. Reference [187] reports one of the first quantitative studies along this line. Representative cases of hilly areas were studied to determine how much adding data from optical frequency standards can improve the modeling of gravity with high spatial resolution. A fully synthetic simulation framework was developed, which will be further refined and applied to other cases of potential interest, like for instance, coastal areas. Interest can be twofold: improvement of reference systems and studies of geophysical phenomena. Some possibilities are discussed in [188–191]. Within the last few years, working groups within the International Association of Geodesy (IAG) were initiated to consider potentialities of chronometric geodesy.

A third aspect is the further development of instrumental capabilities. The long-term goal would be to have ruggedized field-compatible optical frequency standards with uncertainties $\leq 10^{-18}$ and means of comparing them from any place of interest on Earth. First proof-of-concept experiments used coherent optical fiber links between laboratories with already existing state-of-the-art optical frequency metrology programs [62,137]. Over the years, transportable optical frequency standards were developed and currently achieve uncertainties in the 10^{-17} range [192–194], also with the aim to increase technology readiness for space. Recently, one of these devices was used for another proof-of-concept experiment [195], again in conjunction with coherent optical fiber link. Solutions for remote comparisons that reach 10^{-18} , firstly without optical fiber link and secondly over intercontinental distances, are still missing. For the former, free-space optical links studied in [196–199] could be extended to distances of a few 100 km with the help of an airborne platform.

The further advancement of all the above studies shall lead to first pilot chronometric geodesy programs designed for geodesy, geophysics, and Earth's science applications.

9. Conclusions

We described the status of the realization of the SI unit of time according to its current definition. We reported on the tremendous progress of optical frequency metrology over the last 20 years. We mentioned how the CIPM monitors the situation of the field in view of a possible redefinition. To date, work remains to be done to meet all milestones of the CCTF strategy document and to prepare for long-term commitments of the post redefinition era. When this will be achieved, progress of frequency and time metrology and its applications will not need being limited by ¹³³Cs standards any more. Secondary representations of the second and highly accurate redundant measurements of a sufficient number of optical-to-optical frequency ratios will permit a smooth passage through the redefinition. On the basis of these considerations and taking into account the time for due process at the CIPM and at the CGPM, a redefinition seems possible before 2030. The new definition must last long. Paradoxically, the persistent vitality of research on optical frequency standards could be a reason to delay the redefinition, if new breakthroughs suggest even better choices than those being currently consolidated. One possible choice for the redefinition would be to mimic the present definition and select one single optical transition. In this case, secondary representations will be useful, even needed in practice. The possibility, the shape, and the relevance of a system using several transitions on an even basis remain to be explored.

The choice of one particular transition of one particular atom (or a set of them) remains a kind of artifact. Atoms are already complex assemblages of elementary particles and transition frequencies are by far not calculable at the level of uncertainty to which they can be realized by atomic frequency standards. The article of Christian Bordé in the present issue of the *Comptes rendus Physique* presented the deeper foundations of the international system of units, its links with the geometry of space-time and with the system of units introduced by Planck based on five fundamental constants. The system adopted by the 26th CGPM is based on fixing Δv_{CS} , *c*, *h*, *e*, *k*_B. Moving to a system of units founded solely on most fundamental aspects of physical laws would require to abandon atomic transition(s) for the definition of the unit of time. Referring to Planck's system, that would mean fixing the gravitational constant *G* to define Planck's time $\sqrt{\hbar G/c^5}$. Another somewhat intermediate possibility could be to fix the mass of an elementary particle, for instance, the electron m_e , which defines the unit of time as h/m_ec^2 . Given the utmost importance of the practical implementation of the SI system, it does not seem presently possible to opt for one of these fundamental definitions of time, because we do not know how to realize them with the same exquisite accuracy of atomic frequency standards. Research shall investigate how to measure *G* or m_e with radically improved accuracy in the SI system of units attached to atomic transitions.

Acknowledgements

The author is grateful to Christophe Salomon for his patience, his careful reading of the manuscript, and his suggestions to improve it.

References

- [1] Comptes rendus de la 13e CGPM (1967/68) (1969) 103, https://www.bipm.org/en/worldwide-metrology/cgpm/resolutions.html.
- [2] H. Lyons, Spectral lines as frequency standards, Ann. N.Y. Acad. Sci. 55 (5) (1952) 831-871, https://doi.org/10.1111/j.1749-6632.1952.tb26600.x.
- [3] L. Essen, J.V.L. Parry, An atomic standard of frequency and time interval: a caesium resonator, Nature 176 (4476) (1955) 280-282, https://doi.org/10. 1038/176280a0.
- [4] W. Markowitz, R.G. Hall, L. Essen, J.V.L. Parry, Frequency of cesium in terms of ephemeris time, Phys. Rev. Lett. 1 (1958) 105–107, https://doi.org/10. 1103/PhysRevLett.1105.
- [5] T. Quinn, Fifty years of atomic time-keeping: 1955 to 2005, Metrologia 42 (3) (2005), https://doi.org/10.1088/0026-1394/42/3/E01.
- [6] C. Salomon, The measurement of time/La mesure du temps: Foreword, C. R. Physique 16 (5) (2015) 459-460, https://doi.org/10.1016/j.crhy.2015.05. 006.
- [7] Comptes rendus de la 17^e CGPM (1983) 10.
- [8] Comptes rendus de la 26^e CGPM (2018).
- [9] J.-P. Uzan, Varying constants, gravitation and cosmology, Living Rev. Relativ. 14 (2011) 2, http://www.livingreviews.org/lrr-2011-2.
- [10] C.M. Will, The confrontation between general relativity and experiment, Living Rev. Relativ. 17 (1) (2014) 4, https://doi.org/10.12942/lrr-2014-4.
- [11] M.S. Safronova, D. Budker, D. DeMille, D.F.J. Kimball, A. Derevianko, C.W. Clark, Search for new physics with atoms and molecules, Rev. Mod. Phys. 90 (2) (2018) 025008, https://doi.org/10.1103/RevModPhys.90.025008.
- [12] K. Van Tilburg, N. Leefer, L. Bougas, D. Budker, Search for ultralight scalar dark matter with atomic spectroscopy, Phys. Rev. Lett. 115 (2015) 011802, https://doi.org/10.1103/PhysRevLett.115.011802.
- [13] A. Hees, J. Guéna, M. Abgrall, S. Bize, P. Wolf, Searching for an oscillating massive scalar field as a dark matter candidate using atomic hyperfine frequency comparisons, Phys. Rev. Lett. 117 (2016) 061301, https://doi.org/10.1103/PhysRevLett.117.061301.
- [14] N. Ramsey, Experiments with separated oscillatory fields and hydrogen masers, Rev. Mod. Phys. 62 (1990) 541, https://doi.org/10.1103/RevModPhys. 62.541.
- [15] J. Vanier, C. Audoin, The Quantum Physics of Atomic Frequency Standards, Adam Hilger, 1989.
- [16] C.N. Cohen-Tannoudji, Nobel lecture: Manipulating atoms with photons, Rev. Mod. Phys. 70 (1998) 707-719, https://doi.org/10.1103/RevModPhys.70. 707.
- [17] W.D. Phillips, Nobel lecture: Laser cooling and trapping of neutral atoms, Rev. Mod. Phys. 70 (3) (1998) 721-741, https://doi.org/10.1103/RevModPhys. 70.721.
- [18] S. Chu, Nobel lecture: The manipulation of neutral particles, Rev. Mod. Phys. 70 (1998) 685–706, https://doi.org/10.1103/RevModPhys.70.685.
- [19] M. Kasevich, E. Riis, S. Chu, R. de Voe, RF spectroscopy in an atomic fountain, Phys. Rev. Lett. 63 (1989) 612.
- [20] A. Clairon, C. Salomon, S. Guellati, W. Phillips, Ramsey resonance in a Zacharias fountain, Europhys. Lett. 16 (1991) 165.
- [21] A. Clairon, P. Laurent, G. Santarelli, S. Ghezali, S. Lea, M. Bahoura, A cesium fountain frequency standard: recent results, IEEE Trans. Instrum. Meas. 44 (2) (1995) 128.

- [22] G. Santarelli, P. Laurent, P. Lemonde, A. Clairon, A.G. Mann, S. Chang, A.N. Luiten, C. Salomon, Quantum projection noise in an atomic fountain: a high stability cesium frequency standard, Phys. Rev. Lett. 82 (23) (1999) 4619.
- [23] J. Guéna, M. Abgrall, D. Rovera, P. Laurent, B. Chupin, M. Lours, G. Santarelli, P. Rosenbusch, M. Tobar, R. Li, K. Gibble, A. Clairon, S. Bize, Progress in atomic fountains at LNE-SYRTE, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 59 (3) (2012) 391–410, https://doi.org/10.1109/TUFFC.2012.2208.
- [24] J. Guéna, R. Li, K. Gibble, S. Bize, A. Clairon, Evaluation of Doppler shifts to improve the accuracy of primary atomic fountain clocks, Phys. Rev. Lett. 106 (13) (2011) 130801, https://doi.org/10.1103/PhysRevLett.106.130801.
- [25] R. Li, K. Gibble, Evaluating and minimizing distributed cavity phase errors in atomic clocks, Metrologia 47 (5) (2010) 534, http://stacks.iop.org/0026-1394/47/i=5/a=004.
- [26] F. Pereira Dos Santos, H. Marion, S. Bize, Y. Sortais, A. Clairon, C. Salomon, Controlling the cold collision shift in high precision atomic interferometry, Phys. Rev. Lett. 89 (2002) 233004.
- [27] D.J. Papoular, S. Bize, A. Clairon, H. Marion, S.J.J.M.F. Kokkelmans, G.V. Shlyapnikov, Feshbach resonances in cesium at ultralow static magnetic fields, Phys. Rev. A 86 (2012) 040701, https://doi.org/10.1103/PhysRevA.86.040701.
- [28] K. Szymaniec, W. Chałupczak, E. Tiesinga, C.J. Williams, S. Weyers, R. Wynands, Cancellation of the collisional frequency shift in caesium fountain clocks, Phys. Rev. Lett. 98 (15) (2007) 153002, https://doi.org/10.1103/PhysRevLett.98.153002.
- [29] K. Gibble, Scattering of cold-atom coherences by hot atoms: frequency shifts from background-gas collisions, Phys. Rev. Lett. 110 (2013) 180802, https://doi.org/10.1103/PhysRevLett.110.180802.
- [30] K. Szymaniec, S. Lea, K. Liu, An evaluation of the frequency shift caused by collisions with background gas in the primary frequency standard NPL-CsF2, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 61 (1) (2014) 203–206, https://doi.org/10.1109/TUFFC.2014.6689789.
- [31] K. Gibble, Difference between a photon's momentum and an atom's recoil, Phys. Rev. Lett. 97 (7) (2006) 073002, https://doi.org/10.1103/PhysRevLett. 97.073002.
- [32] R. Li, K. Gibble, K. Szymaniec, Improved accuracy of the NPL-CsF2 primary frequency standard: evaluation of distributed cavity phase and microwave lensing frequency shifts, Metrologia 48 (5) (2011) 283, http://stacks.iop.org/0026-1394/48/i=5/a=007.
- [33] K. Gibble, Ramsey spectroscopy, matter-wave interferometry, and the microwave-lensing frequency shift, Phys. Rev. A 90 (2014) 015601, https:// doi.org/10.1103/PhysRevA.90.015601.
- [34] W. Itano, L. Lewis, D. Wineland, Shift of ²S_{1/2} hyperfine splittings due to blackbody radiation, Phys. Rev. A 25 (1982) 1233.
- [35] R. Jefferts, S.P. Heavner, T.E. Parker, T.H. Shirley, J.A. Donley, E.N. Ashby, F. Levi, D. Calonico, G.A. Costanzo, High-accuracy measurement of the blackbody radiation frequency shift of the ground-state hyperfine transition in ¹³³Cs, Phys. Rev. Lett. 112 (2014) 050801, https://doi.org/10.1103/ PhysRevLett.112.050801.
- [36] T.P. Heavner, E.A. Donley, F. Levi, G. Costanzo, T.E. Parker, J.H. Shirley, N. Ashby, S. Barlow, S.R. Jefferts, First accuracy evaluation of NIST-F2, Metrologia 51 (3) (2014) 174, http://stacks.iop.org/0026-1394/51/i=3/a=174.
- [37] F. Levi, D. Calonico, C.E. Calosso, A. Godone, S. Micalizio, G.A. Costanzo, Accuracy evaluation of ITCsF2: a nitrogen cooled caesium fountain, Metrologia 51 (3) (2014) 270, http://stacks.iop.org/0026-1394/51/i=3/a=270.
- [38] A. Jallageas, L. Devenoges, M. Petersen, J. Morel, L.G. Bernier, D. Schenker, P. Thomann, T. Südmeyer, First uncertainty evaluation of the FoCS-2 primary frequency standard, Metrologia 55 (3) (2018) 366, https://doi.org/10.1088/1681-7575/aab3fa.
- [39] R. Wynands, S. Weyers, Atomic fountain clocks, Metrologia 42 (3) (2005) S64, http://stacks.iop.org/0026-1394/42/i=3/a=S08.
- [40] F. Riehle, 8th symposium on frequency standards and metrology 2015, J. Phys. Conf. Ser. 723 (1) (2016) 011001, http://stacks.iop.org/1742-6596/723/ i=1/a=011001.
- [41] J. Guéna, P. Rosenbusch, P. Laurent, M. Abgrall, D. Rovera, G. Santarelli, M. Tobar, S. Bize, A. Clairon, Demonstration of a dual alkali Rb/Cs fountain clock, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 57 (3) (2010) 647, https://doi.org/10.1109/TUFFC.2010.1461.
- [42] J. Guéna, M. Abgrall, A. Clairon, S. Bize, Contributing to TAI with a secondary representation of the SI second, Metrologia 51 (1) (2014) 108, http:// stacks.iop.org/0026-1394/51/i=1/a=108.
- [43] See the list of Secondary Representations of the SI second and of other recommended values of standard frequencies on the BIPM website, http:// www.bipm.org/en/publications/mep.html.
- [44] A. Bauch, J. Achkar, S. Bize, D. Calonico, R. Dach, R. Hlavac, L. Lorini, T. Parker, G. Petit, D. Piester, K. Szymaniec, P. Uhrich, Comparison between frequency standards in Europe and the USA at the 10⁻¹⁵ uncertainty level, Metrologia 43 (1) (2006) 109–120, http://stacks.iop.org/0026-1394/43/109.
- [45] J. Guéna, S. Weyers, M. Abgrall, C. Grebing, V. Gerginov, P. Rosenbusch, S. Bize, B. Lipphardt, H. Denker, N. Quintin, S.M.F. Raupach, D. Nicolodi, F. Stefani, N. Chiodo, S. Koke, A. Kuhl, F. Wiotte, F. Meynadier, E. Camisard, C. Chardonnet, Y.L. Coq, M. Lours, G. Santarelli, A. Amy-Klein, R.L. Targat, O. Lopez, P.E. Pottie, G. Grosche, First international comparison of fountain primary frequency standards via a long distance optical fiber link, Metrologia 54 (3) (2017) 348, http://stacks.iop.org/0026-1394/54/i=3/a=348.
- [46] G. Petit, F. Arias, G. Panfilo, International atomic time: status and future challenges, C. R. Physique 16 (5) (2015) 480–488, https://doi.org/10.1016/j. crhy.2015.03.002.
- [47] See Circular T and fountain reports on the BIPM website, https://www.bipm.org/fr/bipm-services/timescales/time-ftp/Circular-T.html.
- [48] T.E. Parker, Invited review article: The uncertainty in the realization and dissemination of the SI second from a systems point of view, Rev. Sci. Instrum. 83 (2) (2012) 021102, https://doi.org/10.1063/1.3682002.
- [49] G. Petit, G. Panfilo, Comparison of frequency standards used for TAI, IEEE Trans. Instrum. Meas. 62 (99) (2013) 1550, https://doi.org/10.1109/TIM.2012. 2228749.
- [50] F. Riehle, P. Gill, F. Arias, L. Robertsson, The CIPM list of recommended frequency standard values: guidelines and procedures, Metrologia 55 (2) (2018) 188, http://stacks.iop.org/0026-1394/55/i=2/a=188.
- [51] H. Hachisu, G. Petit, F. Nakagawa, Y. Hanado, T. Ido, SI-traceable measurement of an optical frequency at the low 10⁻¹⁶ level without a local primary standard, Opt. Express 25 (8) (2017) 8511–8523, https://doi.org/10.1364/OE.25.008511.
- [52] W.F. McGrew, X. Zhang, H. Leopardi, R.J. Fasano, D. Nicolodi, K. Beloy, J. Yao, J.A. Sherman, S.A. Schäffer, J. Savory, R.C. Brown, S. Römisch, C.W. Oates, T.E. Parker, T.M. Fortier, A.D. Ludlow, Towards adoption of an optical second: verifying optical clocks at the SI limit, arXiv:1811.05885 [physics].
- [53] N. Ashby, T.P. Heavner, S.R. Jefferts, T.E. Parker, A.G. Radnaev, Y.O. Dudin, Testing Local Position Invariance with four cesium-fountain primary frequency standards and four NIST hydrogen masers, Phys. Rev. Lett. 98 (2007) 070802, https://doi.org/10.1103/PhysRevLett.98.070802.
- [54] N. Ashby, T.E. Parker, B.R. Patla, A null test of general relativity based on a long-term comparison of atomic transition frequencies, Nat. Phys. 14 (2018) 822–826, https://doi.org/10.1038/s41567-018-0156-2.
- [55] A. Bauch, S. Weyers, D. Piester, E. Staliuniene, W. Yang, Generation of UTC(PTB) as a fountain-clock based time scale, Metrologia 49 (3) (2012) 180, http://stacks.iop.org/0026-1394/49/i=3/a=180.
- [56] S. Peil, J.L. Hanssen, T.B. Swanson, J. Taylor, C.R. Ekstrom, Evaluation of long term performance of continuously running atomic fountains, Metrologia 51 (3) (2014) 263, http://stacks.iop.org/0026-1394/51/i=3/a=263.
- [57] G.D. Rovera, S. Bize, B. Chupin, J. Guéna, P. Laurent, P. Rosenbusch, P. Uhrich, M. Abgrall, UTC(OP) based on LNE-SYRTE atomic fountain primary frequency standards, Metrologia 53 (3) (2016) S81, http://stacks.iop.org/0026-1394/53/i=3/a=S81.
- [58] G. Petit, F. Arias, A. Harmegnies, G. Panfilo, L. Tisserand, UTCr: a rapid realization of UTC, Metrologia 51 (1) (2014) 33, http://stacks.iop.org/0026-1394/51/i=1/a=33.

- [59] T.W. Hansch, Nobel lecture: Passion for precision, Rev. Mod. Phys. 78 (4) (2006) 1297, https://doi.org/10.1103/RevModPhys.78.1297.
- [60] J.L. Hall, Nobel lecture: Defining and measuring optical frequencies, Rev. Mod. Phys. 78 (4) (2006) 1279, https://doi.org/10.1103/RevModPhys.78.1279.
 [61] W.F. McGrew, X. Zhang, R.J. Fasano, S.A. Schäffer, K. Beloy, D. Nicolodi, R.C. Brown, N. Hinkley, G. Milani, M. Schioppo, T.H. Yoon, A.D. Ludlow, Atomic clock performance enabling geodesy below the centimetre level, Nature 564 (2018) 87–90, https://doi.org/10.1038/s41586-018-0738-2.
- [62] T. Takano, M. Takamoto, I. Ushijima, N. Ohmae, T. Akatsuka, A. Yamaguchi, Y. Kuroishi, H. Munekane, B. Miyahara, H. Katori, Geopotential measurements with synchronously linked optical lattice clocks, Nat. Photonics 10 (2016) 662–666, https://doi.org/10.1038/nphoton.2016.159.
- [63] N. Huntemann, C. Sanner, B. Lipphardt, C. Tamm, E. Peik, Single-ion atomic clock with 3 × 10⁻¹⁸ systematic uncertainty, Phys. Rev. Lett. 116 (2016) 063001, https://doi.org/10.1103/PhysRevLett.116.063001.
- [64] T.L. Nicholson, S.L. Campbell, R.B. Hutson, G.E. Marti, B.J. Bloom, R.L. McNally, W. Zhang, M.D. Barrett, M.S. Safronova, G.F. Strouse, W.L. Tew, J. Ye, Systematic evaluation of an atomic clock at 2 × 10⁻¹⁸ total uncertainty, Nat. Commun. 6 (2015) 6896, https://doi.org/10.1038/ncomms7896.
- [65] I. Ushijima, M. Takamoto, M. Das, T. Ohkubo, H. Katori, Cryogenic optical lattice clocks, Nat. Photonics 9 (2015) 185, https://doi.org/10.1038/nphoton. 2015.5.
- [66] B.J. Bloom, T.L. Nicholson, J.R. Williams, S.L. Campbell, M. Bishof, X. Zhang, W. Zhang, S.L. Bromley, J. Ye, An optical lattice clock with accuracy and stability at the 10⁻¹⁸ level, Nature 506 (2014) 71, https://doi.org/10.1038/nature12941.
- [67] K. Beloy, N. Hinkley, N.B. Phillips, J.A. Sherman, M. Schioppo, J. Lehman, A. Feldman, L.M. Hanssen, C.W. Oates, A.D. Ludlow, Atomic clock with 1×10^{-18} room-temperature blackbody stark uncertainty, Phys. Rev. Lett. 113 (2014) 260801, https://doi.org/10.1103/PhysRevLett.113.260801.
- [68] P. Dubé, A.A. Madej, M. Tibbo, J.E. Bernard, High-accuracy measurement of the differential scalar polarizability of a ⁸⁸Sr⁺ clock using the time-dilation effect, Phys. Rev. Lett. 112 (2014) 173002, https://doi.org/10.1103/PhysRevLett.112.173002.
- [69] T. Middelmann, S. Falke, C. Lisdat, U. Sterr, High accuracy correction of blackbody radiation shift in an optical lattice clock, Phys. Rev. Lett. 109 (2012) 263004, https://doi.org/10.1103/PhysRevLett.109.263004.
- [70] J.A. Sherman, N.D. Lemke, N. Hinkley, M. Pizzocaro, R.W. Fox, A.D. Ludlow, C.W. Oates, High-accuracy measurement of atomic polarizability in an optical lattice clock, Phys. Rev. Lett. 108 (2012) 153002, https://doi.org/10.1103/PhysRevLett.108.153002.
- [71] M. Doležal, P. Balling, P.B.R. Nisbet-Jones, S.A. King, J.M. Jones, H.A. Klein, P. Gill, T. Lindvall, A.E. Wallin, M. Merimaa, C. Tamm, C. Sanner, N. Huntemann, N. Scharnhorst, I.D. Leroux, P.O. Schmidt, T. Burgermeister, T.E. Mehlstäubler, E. Peik, Analysis of thermal radiation in ion traps for optical frequency standards, Metrologia 52 (6) (2015) 842, http://stacks.iop.org/0026-1394/52/i=6/a=842.
- [72] J. Lodewyck, M. Zawada, L. Lorini, M. Gurov, P. Lemonde, Observation and cancellation of a perturbing dc stark shift in strontium optical lattice clocks, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 59 (3) (2012) 411–415, https://doi.org/10.1109/TUFFC.2012.2209.
- [73] P.G. Westergaard, J. Lodewyck, L. Lorini, A. Lecallier, E.A. Burt, M. Zawada, J. Millo, P. Lemonde, Lattice-induced frequency shifts in Sr optical lattice clocks at the 10⁻¹⁷ level, Phys. Rev. Lett. 106 (21) (2011) 210801, https://doi.org/10.1103/PhysRevLett.106.210801.
- [74] R.C. Brown, N.B. Phillips, K. Beloy, W.F. McGrew, M. Schioppo, R.J. Fasano, G. Milani, X. Zhang, N. Hinkley, H. Leopardi, T.H. Yoon, D. Nicolodi, T.M. Fortier, A.D. Ludlow, Hyperpolarizability and operational magic wavelength in an optical lattice clock, Phys. Rev. Lett. 119 (25) (2017) 253001, https://doi.org/10.1103/PhysRevLett.119.253001.
- [75] S.L. Campbell, R.B. Hutson, G.E. Marti, A. Goban, N. Darkwah Oppong, R.L. McNally, L. Sonderhouse, J.M. Robinson, W. Zhang, B.J. Bloom, J. Ye, A Fermi-degenerate three-dimensional optical lattice clock, Science 358 (6359) (2017) 90–94, https://doi.org/10.1126/science.aam5538, http://science.sciencemag.org/content/358/6359/90.full.pdf.
- [76] V.I. Yudin, A.V. Taichenachev, C.W. Oates, Z.W. Barber, N.D. Lemke, A.D. Ludlow, U. Sterr, C. Lisdat, F. Riehle, Hyper-Ramsey spectroscopy of optical clock transitions, Phys. Rev. A 82 (1) (2010) 011804, https://doi.org/10.1103/PhysRevA.82.011804.
- [77] C.W. Chou, D.B. Hume, J.C.J. Koelemeij, D.J. Wineland, T. Rosenband, Frequency comparison of two high-accuracy Al⁺ optical clocks, Phys. Rev. Lett. 104 (7) (2010) 070802, https://doi.org/10.1103/PhysRevLett.104.070802.
- [78] R. Le Targat, L. Lorini, Y. Le Coq, M. Zawada, J. Guéna, M. Abgrall, M. Gurov, P. Rosenbusch, D.G. Rovera, B. Nagórny, R. Gartman, P.G. Westergaard, M.E. Tobar, M. Lours, G. Santarelli, A. Clairon, S. Bize, P. Laurent, P. Lemonde, J. Lodewyck, Experimental realization of an optical second with strontium lattice clocks, Nat. Commun. 4 (2013) 2109, https://doi.org/10.1038/ncomms3109.
- [79] B.J. Bloom, T.L. Nicholson, J.R. Williams, S.L. Campbell, M. Bishof, X. Zhang, W. Zhang, S.L. Bromley, J. Ye, A new generation of atomic clocks: accuracy and stability at the 10⁻¹⁸ level, arXiv:1309.1137.
- [80] Y. Huang, H. Guan, P. Liu, W. Bian, L. Ma, K. Liang, T. Li, K. Gao, Frequency comparison of two ⁴⁰Ca⁺ optical clocks with an uncertainty at the 10⁻¹⁷ level, Phys. Rev. Lett. 116 (2016) 013001, https://doi.org/10.1103/PhysRevLett.116.013001.
- [81] A.D. Ludlow, M.M. Boyd, J. Ye, E. Peik, P.O. Schmidt, Optical atomic clocks, Rev. Mod. Phys. 87 (2015) 637–701, https://doi.org/10.1103/RevModPhys. 87.637.
- [82] J. Bergquist, W. Itano, D. Wineland, Recoilless optical absorption and Doppler sidebands of a single trapped ion, Phys. Rev. A 36 (1987) 428.
- [83] W. Paul, Electromagnetic traps for charged and neutral particles, Rev. Mod. Phys. 62 (1990) 531.
- [84] D. Leibfried, R. Blatt, C. Monroe, D. Wineland, Quantum dynamics of single trapped ions, Rev. Mod. Phys. 75 (1) (2003) 281–324, https://doi.org/10. 1103/RevModPhys.75.281.
- [85] H. Katori, Spectroscopy of strontium atoms in the Lamb-Dicke confinement, in: Proc. of the 6th Symposium on Frequency Standards and Metrology, World Scientific, Singapore, 2001, p. 323.
- [86] H. Katori, M. Takamoto, V.G. Pal'chikov, V.D. Ovsiannikov, Ultrastable optical clock with neutral atoms in an engineered light shift trap, Phys. Rev. Lett. 91 (17) (2003) 173005, https://doi.org/10.1103/PhysRevLett.91.173005.
- [87] J. Ye, H.J. Kimble, H. Katori, Quantum state engineering and precision metrology using state-insensitive light traps, Science 320 (5884) (2008) 1734–1738, https://doi.org/10.1126/science.1148259, http://www.sciencemag.org/content/320/5884/1734.full.pdf.
- [88] M. Schioppo, R.C. Brown, W.F. McGrew, N. Hinkley, R.J. Fasano, K. Beloy, T.H. Yoon, G. Milani, D. Nicolodi, J.A. Sherman, N.B. Phillips, C.W. Oates, A.D. Ludlow, Ultrastable optical clock with two cold-atom ensembles, Nat. Photonics 11 (1) (2017) 48–52, https://doi.org/10.1038/nphoton.2016.231.
- [89] N. Hinkley, J.A. Sherman, N.B. Phillips, M. Schioppo, N.D. Lemke, K. Beloy, M. Pizzocaro, C.W. Oates, A.D. Ludlow, An atomic clock with 10⁻¹⁸ instability, Science 341 (6151) (2013) 1215–1218, https://doi.org/10.1126/science.1240420, http://www.sciencemag.org/content/341/6151/1215.full.pdf.
- [90] M. Takamoto, T. Takano, H. Katori, Frequency comparison of optical lattice clocks beyond the Dick limit, Nat. Photonics 5 (2011) 288–292, https:// doi.org/10.1038/nphoton.2011.34.
- [91] T.L. Nicholson, M.J. Martin, J.R. Williams, B.J. Bloom, M. Bishof, M.D. Swallows, S.L. Campbell, J. Ye, Comparison of two independent Sr optical clocks with 1 × 10⁻¹⁷ stability at 10³ s, Phys. Rev. Lett. 109 (2012) 230801, https://doi.org/10.1103/PhysRevLett.109.230801.
- [92] A. Al-Masoudi, S. Dörscher, S. Häfner, U. Sterr, C. Lisdat, Noise and instability of an optical lattice clock, Phys. Rev. A 92 (2015) 063814, https:// doi.org/10.1103/PhysRevA.92.063814.
- [93] G.E. Marti, R.B. Hutson, A. Goban, S.L. Campbell, N. Poli, J. Ye, Imaging optical frequencies with 100 μHz precision and 1.1 μm resolution, Phys. Rev. Lett. 120 (10) (2018) 103201, https://doi.org/10.1103/PhysRevLett.120.103201.
- [94] C.W. Chou, D.B. Hume, M.J. Thorpe, D.J. Wineland, T. Rosenband, Quantum coherence between two atoms beyond $Q = 10^{15}$, Phys. Rev. Lett. 106 (16) (2011) 160801, https://doi.org/10.1103/PhysRevLett.106.160801.
- [95] K. Numata, A. Kemery, J. Camp, Thermal-noise limit in the frequency stabilization of lasers with rigid cavities, Phys. Rev. Lett. 93 (25) (2004) 250602, https://doi.org/10.1103/PhysRevLett.93.250602.

- [96] M. Swallows, M. Martin, M. Bishof, C. Benko, Y. Lin, S. Blatt, A. Rey, J. Ye, Operating a ⁸⁷Sr optical lattice clock with high precision and at high density, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 59 (3) (2012) 416–425, https://doi.org/10.1109/TUFFC.2012.2210.
- [97] S. Häfner, S. Falke, C. Grebing, S. Vogt, T. Legero, M. Merimaa, C. Lisdat, U. Sterr, 8 × 10⁻¹⁷ fractional laser frequency instability with a long room-temperature cavity, Opt. Lett. 40 (9) (2015) 2112–2115, https://doi.org/10.1364/OL40.002112.
- [98] T. Kessler, C. Hagemann, C. Grebing, T. Legero, U. Sterr, F. Riehle, M.J. Martin, L. Chen, J. Ye, A sub-40-mHz-linewidth laser based on a silicon singlecrystal optical cavity, Nat. Photonics 6 (10) (2012) 687-692, https://doi.org/10.1038/nphoton.2012.217.
- [99] D.G. Matei, T. Legero, S. Häfner, C. Grebing, R. Weyrich, W. Zhang, L. Sonderhouse, J.M. Robinson, J. Ye, F. Riehle, U. Sterr, 1.5 μm lasers with sub 10 mHz linewidth, Phys. Rev. Lett. 118 (2017) 263202, https://doi.org/10.1103/PhysRevLett.118.263202.
- [100] G.D. Cole, W. Zhang, M.J. Martin, J. Ye, M. Aspelmeyer, Tenfold reduction of brownian noise in high-reflectivity optical coatings, Nat. Photonics 7 (8) (2013) 644–650, https://doi.org/10.1038/nphoton.2013.174.
- [101] Q.-F. Chen, A. Troshyn, I. Ernsting, S. Kayser, S. Vasilyev, A. Nevsky, S. Schiller, Spectrally narrow, long-term stable optical frequency reference based on a Eu³⁺: Y₂SiO₅ crystal at cryogenic temperature, Phys. Rev. Lett. 107 (2011) 223202, https://doi.org/10.1103/PhysRevLett.107.223202.
- [102] M.J. Thorpe, L. Rippe, T.M. Fortier, M.S. Kirchner, T. Rosenband, Frequency stabilization to 6 × 10⁻¹⁶ via spectral-hole burning, Nat. Photonics 6 (2011) 688, https://doi.org/10.1038/nphoton.2011.215.
- [103] O. Gobron, K. Jung, N. Galland, K. Predehl, R.L. Targat, A. Ferrier, P. Goldner, S. Seidelin, Y.L. Coq, Dispersive heterodyne probing method for laser frequency stabilization based on spectral hole burning in rare-earth doped crystals, Opt. Express 25 (13) (2017) 15539–15548, https://doi.org/10. 1364/OE.25.015539.
- [104] D. Yu, J. Chen, Optical clock with millihertz linewidth based on a phase-matching effect, Phys. Rev. Lett. 98 (5) (2007) 050801, https://doi.org/10. 1103/PhysRevLett.98.050801.
- [105] D. Meiser, J. Ye, D.R. Carlson, M.J. Holland, Prospects for a millihertz-linewidth laser, Phys. Rev. Lett. 102 (16) (2009) 163601, https://doi.org/10.1103/ PhysRevLett.102.163601.
- [106] M.A. Norcia, J.R.K. Cline, J.A. Muniz, J.M. Robinson, R.B. Hutson, A. Goban, G.E. Marti, J. Ye, J.K. Thompson, Frequency measurements of superradiance from the strontium clock transition, Phys. Rev. X 8 (2) (2018) 021036, https://doi.org/10.1103/PhysRevX.8.021036.
- [107] S. Wehner, D. Elkouss, R. Hanson, Quantum internet: a vision for the road ahead, Science 362 (6412) (2018) eaam9288, https://doi.org/10.1126/ science.aam9288.
- [108] C.L. Degen, F. Reinhard, P. Cappellaro, Quantum sensing, Rev. Mod. Phys. 89 (2017) 035002, https://doi.org/10.1103/RevModPhys.89.035002.
- [109] P.O. Schmidt, T. Rosenband, C. Langer, W.M. Itano, J.C. Bergquist, D.J. Wineland, Spectroscopy using quantum logic, Science 309 (5735) (2005) 749–752, https://doi.org/10.1126/science.1114375.
- [110] D.J. Wineland, Nobel lecture: Superposition, entanglement, and raising Schrödinger's cat, Rev. Mod. Phys. 85 (2013) 1103–1114, https://doi.org/10. 1103/RevModPhys.85.1103.
- [111] C.F. Roos, M. Chwalla, K. Kim, M. Riebe, R. Blatt, "Designer atoms" for quantum metrology, Nature 443 (7109) (2006) 316–319, https://doi.org/10. 1038/nature05101.
- [112] D. Leibfried, B. DeMarco, V. Meyer, M. Rowe, A. Ben-Kish, J. Britton, W.M. Itano, B. Jelenković, C. Langer, T. Rosenband, D.J. Wineland, Trapped-ion quantum simulator: Experimental application to nonlinear interferometers, Phys. Rev. Lett. 89 (2002) 247901, https://doi.org/10.1103/PhysRevLett.89. 247901.
- [113] J. Keller, T. Burgermeister, D. Kalincev, A. Didier, A.P. Kulosa, T. Nordmann, J. Kiethe, T.E. Mehlstäubler, Controlling systematic frequency uncertainties at the 10⁻¹⁹ level in linear Coulomb crystals, Phys. Rev. A 99 (1) (2019) 013405, https://doi.org/10.1103/PhysRevA.99.013405.
- [114] I. Bouchoule, K. Mølmer, Preparation of spin-squeezed atomic states by optical-phase-shift measurement, Phys. Rev. A 66 (4) (2002) 043811, https:// doi.org/10.1103/PhysRevA.66.043811.
- [115] D.J. Wineland, J.J. Bollinger, W.M. Itano, F.L. Moore, D.J. Heinzen, Spin squeezing and reduced quantum noise in spectroscopy, Phys. Rev. A 46 (1992) R6797-R6800, https://doi.org/10.1103/PhysRevA.46.R6797.
- [116] P.J. Windpassinger, D. Oblak, P.G. Petrov, M. Kubasik, M. Saffman, C.L.G. Alzar, J. Appel, J.H. Müller, N. Kjærgaard, E.S. Polzik, Nondestructive probing of Rabi oscillations on the cesium clock transition near the standard quantum limit, Phys. Rev. Lett. 100 (10) (2008) 103601, https://doi.org/10.1103/ PhysRevLett.100.103601.
- [117] A. Louchet-Chauvet, J. Appel, J.J. Renema, D. Oblak, N. Kjaergaard, E.S. Polzik, Entanglement-assisted atomic clock beyond the projection noise limit, New J. Phys. 12 (6) (2010) 065032, http://stacks.iop.org/1367-2630/12/i=6/a=065032.
- [118] O. Hosten, N.J. Engelsen, R. Krishnakumar, M.A. Kasevich, Measurement noise 100 times lower than the quantum-projection limit using entangled atoms, Nature 529 (7587) (2016) 505–508, https://doi.org/10.1038/nature16176.
- [119] J. Lodewyck, P.G. Westergaard, P. Lemonde, Nondestructive measurement of the transition probability in a Sr optical lattice clock, Phys. Rev. A 79 (6) (2009) 061401, https://doi.org/10.1103/PhysRevA.79.061401.
- [120] G. Vallet, E. Bookjans, U. Eismann, S. Bilicki, R.L. Targat, J. Lodewyck, A noise-immune cavity-assisted non-destructive detection for an optical lattice clock in the quantum regime, New J. Phys. 19 (8) (2017) 083002, http://stacks.iop.org/1367-2630/19/i=8/a=083002.
- [121] N. Shiga, M. Takeuchi, Locking the local oscillator phase to the atomic phase via weak measurement, New J. Phys. 14 (2) (2012) 023034, http:// stacks.iop.org/1367-2630/14/i=2/a=023034.
- [122] R. Kohlhaas, A. Bertoldi, E. Cantin, A. Aspect, A. Landragin, P. Bouyer, Phase locking a clock oscillator to a coherent atomic ensemble, Phys. Rev. X 5 (2015) 021011, https://doi.org/10.1103/PhysRevX.5.021011.
- [123] I. Kruse, K. Lange, J. Peise, B. Lücke, L. Pezzè, J. Arlt, W. Ertmer, C. Lisdat, L. Santos, A. Smerzi, C. Klempt, Improvement of an atomic clock using squeezed vacuum, Phys. Rev. Lett. 117 (14) (2016) 143004, https://doi.org/10.1103/PhysRevLett.117.143004.
- [124] I.D. Leroux, M.H. Schleier-Smith, V. Vuletić, Orientation-dependent entanglement lifetime in a squeezed atomic clock, Phys. Rev. Lett. 104 (25) (2010) 250801, https://doi.org/10.1103/PhysRevLett.104.250801.
- [125] I.D. Leroux, M.H. Schleier-Smith, V. Vuletić, Implementation of cavity squeezing of a collective atomic spin, Phys. Rev. Lett. 104 (7) (2010) 073602, https://doi.org/10.1103/PhysRevLett.104.073602.
- [126] S.A. Diddams, The evolving optical frequency comb [invited], J. Opt. Soc. Am. B 27 (11) (2010) B51-B62, https://doi.org/10.1364/JOSAB.27.000B51.
- [127] N.R. Newbury, Searching for applications with a fine-tooth comb, Nat. Photonics 5 (2011) 186–188, https://doi.org/10.1038/nphoton.2011.38.
- [128] D. Nicolodi, B. Argence, W. Zhang, R. Le Targat, G. Santarelli, Y. Le Coq, Spectral purity transfer between optical wavelengths at the 10⁻¹⁸ level, Nat. Photonics 8 (2014) 219, https://doi.org/10.1038/nphoton.2013.361.
- [129] X. Xie, R. Bouchand, D. Nicolodi, M. Giunta, W. Hänsel, M. Lezius, A. Joshi, S. Datta, C. Alexandre, M. Lours, P.-A. Tremblin, G. Santarelli, R. Holzwarth, Y. Le Coq, Photonic microwave signals with zeptosecond-level absolute timing noise, Nat. Photonics 11 (1) (2017) 44–47, https://doi.org/10.1038/ nphoton.2016.215.
- [130] F.N. Baynes, F. Quinlan, T.M. Fortier, Q. Zhou, A. Beling, J.C. Campbell, S.A. Diddams, Attosecond timing in optical-to-electrical conversion, Optica 2 (2) (2015) 141–146, https://doi.org/10.1364/OPTICA.2.000141.
- [131] N. Ohmae, N. Kuse, M.E. Fermann, H. Katori, All-polarization-maintaining, single-porter: fiber comb for high-stability comparison of optical lattice clocks, Appl. Phys. Express 10 (6) (2017) 062503, http://stacks.iop.org/1882-0786/10/i=6/a=062503.
- [132] Y. Le Coq, et al., Peignes de fréquences femtosecondes pour la mesure des fréquences optiques, Rev. Fr. Métrol. 32 (2012) 35.

- [133] J. Millo, M. Abgrall, M. Lours, E.M.L. English, H. Jiang, J. Guéna, A. Clairon, M.E. Tobar, S. Bize, Y.L. Coq, G. Santarelli, Ultralow noise microwave generation with fiber-based optical frequency comb and application to atomic fountain clock, Appl. Phys. Lett. 94 (14) (2009) 141105, https://doi.org/ 10.1063/1.3112574.
- [134] S. Weyers, B. Lipphardt, H. Schnatz, Reaching the quantum limit in a fountain clock using a microwave oscillator phase locked to an ultrastable laser, Phys. Rev. A 79 (3) (2009) 031803, https://doi.org/10.1103/PhysRevA.79.031803.
- [135] B. Lipphardt, V. Gerginov, S. Weyers, Optical stabilization of a microwave oscillator for fountain clock interrogation, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 64 (4) (2017) 761-766, https://doi.org/10.1109/TUFFC.2017.2649044.
- [136] J. Lodewyck, S. Bilicki, E. Bookjans, J.-L. Robyr, C. Shi, G. Vallet, R.L. Targat, D. Nicolodi, Y.L. Coq, J. Guéna, M. Abgrall, P. Rosenbusch, S. Bize, Optical to microwave clock frequency ratios with a nearly continuous strontium optical lattice clock, Metrologia 53 (4) (2016) 1123, http://stacks.iop.org/0026-1394/53/i=4/a=1123.
- [137] C. Lisdat, G. Grosche, N. Quintin, C. Shi, S. Raupach, C. Grebing, D. Nicolodi, F. Stefani, A. Al-Masoudi, S. Dorscher, S. Hafner, J.-L. Robyr, N. Chiodo, S. Bilicki, E. Bookjans, A. Koczwara, S. Koke, A. Kuhl, F. Wiotte, F. Meynadier, E. Camisard, M. Abgrall, M. Lours, T. Legero, H. Schnatz, U. Sterr, H. Denker, C. Chardonnet, Y. Le Coq, G. Santarelli, A. Amy-Klein, R. Le Targat, J. Lodewyck, O. Lopez, P.-E. Pottie, A clock network for geodesy and fundamental science, Nat. Commun. 7 (2016) 12443, https://doi.org/10.1038/ncomms12443.
- [138] R. Tyumenev, M. Favier, S. Bilicki, E. Bookjans, R.L. Targat, J. Lodewyck, D. Nicolodi, Y.L. Coq, M. Abgrall, J. Guéna, L.D. Sarlo, S. Bize, Comparing a mercury optical lattice clock with microwave and optical frequency standards, New J. Phys. 18 (11) (2016) 113002, http://stacks.iop.org/1367-2630/ 18/i=11/a=113002.
- [139] G. Grosche, B. Lipphardt, H. Schnatz, G. Santarelli, P. Lemonde, S. Bize, M. Lours, F. Narbonneau, A. Clairon, O. Lopez, A. Amy-Klein, C. Chardonnet, Transmission of an optical carrier frequency over a telecommunication fiber link, in: Conference on Lasers and Electro-Optics/Quantum Electronics and Laser Science Conference and Photonic Applications Systems Technologies, Optical Society of America, 2007, p. CMKK1, http://www.opticsinfobase. org/abstract.cfm?URI=CLEO-2007-CMKK1.
- [140] H. Jiang, F. Kéfélian, S. Crane, O. Lopez, M. Lours, J. Millo, D. Holleville, P. Lemonde, C. Chardonnet, A. Amy-Klein, G. Santarelli, Long-distance frequency transfer over an urban fiber link using optical phase stabilization, J. Opt. Soc. Am. B 25 (12) (2008) 2029–2035, http://josab.osa.org/abstract.cfm?URI= josab-25-12-2029.
- [141] G. Grosche, O. Terra, K. Predehl, R. Holzwarth, B. Lipphardt, F. Vogt, U. Sterr, H. Schnatz, Optical frequency transfer via 146 km fiber link with 10⁻¹⁹ relative accuracy, Opt. Lett. 34 (15) (2009) 2270–2272, https://doi.org/10.1364/OL.34.002270.
- [142] K. Predehl, G. Grosche, S.M.F. Raupach, S. Droste, O. Terra, J. Alnis, T. Legero, T.W. Hänsch, T. Udem, R. Holzwarth, H. Schnatz, A 920-kilometer optical fiber link for frequency metrology at the 19th decimal place, Science 336 (6080) (2012) 441–444, https://doi.org/10.1126/science.1218442, http:// www.sciencemag.org/content/336/6080/441.full.pdf.
- [143] O. Lopez, A. Haboucha, B. Chanteau, C. Chardonnet, A. Amy-Klein, G. Santarelli, Ultra-stable long distance optical frequency distribution using the internet fiber network, Opt. Express 20 (21) (2012) 23518–23526, https://doi.org/10.1364/OE.20.023518.
- [144] F. Guillou-Camargo, V. Ménoret, E. Cantin, O. Lopez, N. Quintin, E. Camisard, V. Salmon, J.-M.L. Merdy, G. Santarelli, A. Amy-Klein, P.-E. Pottie, B. Desruelle, C. Chardonnet, First industrial-grade coherent fiber link for optical frequency standard dissemination, Appl. Opt. 57 (25) (2018) 7203–7210, https://doi.org/10.1364/A0.57.007203.
- [145] F. Frank, F. Stefani, P. Tuckey, P.E. Pottie, A sub-ps stability time transfer method based on optical modems, IEEE Trans. Ultrason. Ferroelectr. Freq. Control 65 (6) (2018) 1001–1006, https://doi.org/10.1109/TUFFC.2018.2833389.
- [146] O. Lopez, F. Kéfélian, H. Jiang, A. Haboucha, A. Bercy, F. Stefani, B. Chanteau, A. Kanj, D. Rovera, J. Achkar, C. Chardonnet, P.-E. Pottie, A. Amy-Klein, G. Santarelli, Frequency and time transfer for metrology and beyond using telecommunication network fibres, C. R. Physique 16 (5) (2015) 531–539, https://doi.org/10.1016/j.crhy.2015.04.005.
- [147] L. Sliwczynski, et al., Dissemination of time and RF frequency via a stabilized fibre optic link over a distance of 420 km, Metrologia 50 (2) (2013) 133, http://stacks.iop.org/0026-1394/50/i=2/a=133.
- [148] O. Lopez, A. Kanj, P.-E. Pottie, D. Rovera, J. Achkar, C. Chardonnet, A. Amy-Klein, G. Santarelli, Simultaneous remote transfer of accurate timing and optical frequency over a public fiber network, Appl. Phys. B, Lasers Opt. 110 (2013) 3, https://doi.org/10.1007/s00340-012-5241-0.
- [149] J. Geršl, P. Delva, P. Wolf, Relativistic corrections for time and frequency transfer in optical fibres, Metrologia 52 (4) (2015) 552, http://stacks.iop.org/ 0026-1394/52/i=4/a=552.
- [150] S. Schiller, Feasibility of giant fiber-optic gyroscopes, Phys. Rev. A 87 (2013) 033823, https://doi.org/10.1103/PhysRevA.87.033823.
- [151] C. Clivati, D. Calonico, G.A. Costanzo, A. Mura, M. Pizzocaro, F. Levi, Large-area fiber-optic gyroscope on a multiplexed fiber network, Opt. Lett. 38 (7) (2013) 1092–1094, https://doi.org/10.1364/OL.38.001092.
- [152] G. Marra, C. Clivati, R. Luckett, A. Tampellini, J. Kronjäger, L. Wright, A. Mura, F. Levi, S. Robinson, A. Xuereb, B. Baptie, D. Calonico, Ultrastable laser interferometry for earthquake detection with terrestrial and submarine cables, Science 361 (6401) (2018) 486–490, https://doi.org/10.1126/science. aat4458.
- [153] P. Delva, J. Lodewyck, S. Bilicki, E. Bookjans, G. Vallet, R. Le Targat, P.-E. Pottie, C. Guerlin, F. Meynadier, C. Le Poncin-Lafitte, O. Lopez, A. Amy-Klein, W.-K. Lee, N. Quintin, C. Lisdat, A. Al-Masoudi, S. Dörscher, C. Grebing, G. Grosche, A. Kuhl, S. Raupach, U. Sterr, I.R. Hill, R. Hobson, W. Bowden, J. Kronjäger, G. Marra, A. Rolland, F.N. Baynes, H.S. Margolis, P. Gill, Test of special relativity using a fiber network of optical clocks, Phys. Rev. Lett. 118 (2017) 221102, https://doi.org/10.1103/PhysRevLett.118.221102.
- [154] S. Reynaud, C. Salomon, P. Wolf, Testing general relativity with atomic clocks, Space Sci. Rev. 148 (1) (2009) 233–247, https://doi.org/10.1007/s11214-009-9539-0.
- [155] D. Mattingly, Modern tests of Lorentz invariance, Living Rev. Relativ. 8 (5) (2005), http://www.livingreviews.org/lrr-2005-5.
- [156] W. Marciano, Time variation of the fundamental "constants" and Kaluza-Klein theories, Phys. Rev. Lett. 52 (1984) 489.
- [157] T. Damour, A. Polyakov, The string dilaton and a least coupling principle, Nucl. Phys. B 423 (1994) 532.
- [158] T. Damour, Theoretical aspects of the equivalence principle, Class. Quantum Gravity 29 (18) (2012) 184001, http://stacks.iop.org/0264-9381/29/i=18/ a=184001.
- [159] J. Guéna, M. Abgrall, D. Rovera, P. Rosenbusch, M.E. Tobar, P. Laurent, A. Clairon, S. Bize, Improved tests of local position invariance using ⁸⁷Rb and ¹³³Cs fountains, Phys. Rev. Lett. 109 (2012) 080801, https://doi.org/10.1103/PhysRevLett.109.080801.
- [160] P. Wcislo, P. Morzynski, M. Bober, A. Cygan, D. Lisak, R. Ciurylo, M. Zawada, Experimental constraint on dark matter detection with optical atomic clocks, Nat. Astron. 1 (2016) 0009, https://doi.org/10.1038/s41550-016-0009.
- [161] P. Delva, N. Puchades, E. Schönemann, F. Dilssner, C. Courde, S. Bertone, F. Gonzalez, A. Hees, C. Le Poncin-Lafitte, F. Meynadier, R. Prieto-Cerdeira, B. Sohet, J. Ventura-Traveset, P. Wolf, Gravitational redshift test using eccentric Galileo satellites, Phys. Rev. Lett. 121 (23) (2018) 231101, https:// doi.org/10.1103/PhysRevLett.121.231101.
- [162] F. Meynadier, P. Delva, C. Le Poncin-Lafitte, C. Guerlin, P. Wolf, Atomic clock ensemble in space (ACES) data analysis, Class. Quantum Gravity 35 (3) (2018) 035018, https://doi.org/10.1088/1361-6382/aaa279.
- [163] P. Laurent, D. Massonnet, L. Cacciapuoti, C. Salomon, The ACES/PHARAO space mission, C. R. Physique 16 (5) (2015) 540–552, https://doi.org/10.1016/ j.crhy.2015.05.002.
- [164] L. Cacciapuoti, C. Salomon, Space clocks and fundamental tests: the ACES experiment, Eur. Phys. J. Spec. Top. 172 (2009) 57–68, https://doi.org/10. 1140/epjst/e2009-01041-7.

- [165] L. Cacciapuoti, N. Dimarcq, G. Santarelli, P. Laurent, P. Lemonde, A. Clairon, P. Berthoud, A. Jornod, F. Reina, S. Feltham, C. Salomon, Atomic clock ensemble in space: scientific objectives and mission status, Nucl. Phys. B, Proc. Suppl. 166 (2007) 303–306, https://doi.org/10.1016/j.nuclphysbps. 2006.12.033, Proceedings of the Third International Conference on Particle and Fundamental Physics in Space.
- [166] P. Touboul, G. Métris, M. Rodrigues, Y. André, Q. Baghi, J. Bergé, D. Boulanger, S. Bremer, P. Carle, R. Chhun, B. Christophe, V. Cipolla, T. Damour, P. Danto, H. Dittus, P. Fayet, B. Foulon, C. Gageant, P.-Y. Guidotti, D. Hagedorn, E. Hardy, P.-A. Huynh, H. Inchauspe, P. Kayser, S. Lala, C. Lämmerzahl, V. Lebat, P. Leseur, F. m. c. Liorzou, M. List, F. Löffler, I. Panet, B. Pouilloux, P. Prieur, A. Rebray, S. Reynaud, B. Rievers, A. Robert, H. Selig, L. Serron, T. Sumner, N. Tanguy, P. Visser, MICROSCOPE mission: first results of a space test of the equivalence principle, Phys. Rev. Lett. 119 (2017) 231101, https://doi.org/10.1103/PhysRevLett.119.231101.
- [167] Consultative Committee for Time and Frequency (CCTF), Report of the 15th meeting (June 2001) to the International Committee for Weights and Measures, Tech. rep., BIPM, 2001.
- [168] H. Marion, F. Pereira Dos Santos, M. Abgrall, S. Zhang, Y. Sortais, S. Bize, I. Maksimovic, D. Calonico, J. Grünert, C. Mandache, P. Lemonde, G. Santarelli, P. Laurent, A. Clairon, C. Salomon, Search for variations of fundamental constants using atomic fountain clocks, Phys. Rev. Lett. 90 (2003) 150801, https://doi.org/10.1103/PhysRevLett.90.150801.
- [169] S. Bize, P. Laurent, M. Abgrall, H. Marion, I. Maksimovic, L. Cacciapuoti, J. Grünert, C. Vian, F. Pereira dos Santos, P. Rosenbusch, P. Lemonde, G. Santarelli, P. Wolf, A. Clairon, A. Luiten, M. Tobar, C. Salomon, Advances in ¹³³Cs fountains, C. R. Physique 5 (2004) 829.
- [170] H.S. Margolis, P. Gill, Least-squares analysis of clock frequency comparison data to deduce optimized frequency and frequency ratio values, Metrologia 52 (5) (2015) 628, http://stacks.iop.org/0026-1394/52/i=5/a=628.
- [171] L. Robertsson, On the evaluation of ultra-high-precision frequency ratio measurements: examining closed loops in a graph theory framework, Metrologia 53 (6) (2016) 1272–1280, https://doi.org/10.1088/0026-1394/53/6/1272.
- [172] C. Grebing, A. Al-Masoudi, S. Dörscher, S. Häfner, V. Gerginov, S. Weyers, B. Lipphardt, F. Riehle, U. Sterr, C. Lisdat, Realization of a timescale with an accurate optical lattice clock, Optica 3 (6) (2016) 563–569, https://doi.org/10.1364/OPTICA.3.000563.
- [173] T. Rosenband, D.B. Hume, P.O. Schmidt, C.W. Chou, A. Brusch, L. Lorini, W.H. Oskay, R.E. Drullinger, T.M. Fortier, J.E. Stalnaker, S.A. Diddams, W.C. Swann, N.R. Newbury, W.M. Itano, D.J. Wineland, J.C. Bergquist, Frequency ratio of Al⁺ and Hg⁺ single-ion optical clocks; metrology at the 17th decimal place, Science 319 (2008) 1808.
- [174] N. Nemitz, T. Ohkubo, M. Takamoto, I. Ushijima, M. Das, N. Ohmae, H. Katori, Frequency ratio of Yb and Sr clocks with 5 × 10⁻¹⁷ uncertainty at 150 seconds averaging time, Nat. Photonics 10 (4) (2016) 258–261, https://doi.org/10.1038/nphoton.2016.20.
- [175] K. Yamanaka, N. Ohmae, I. Ushijima, M. Takamoto, H. Katori, Frequency ratio of ¹⁹⁹Hg and ⁸⁷Sr optical lattice clocks beyond the SI limit, Phys. Rev. Lett. 114 (2015) 230801, https://doi.org/10.1103/PhysRevLett.114.230801.
- [176] H. Hachisu, F. Nakagawa, Y. Hanado, T. Ido, Months-long real-time generation of a time scale based on an optical clock, Sci. Rep. 8 (1) (2018) 4243, https://doi.org/10.1038/s41598-018-22423-5.
- [177] H. Hachisu, M. Fujieda, S. Nagano, T. Gotoh, A. Nogami, T. Ido, S. Falke, N. Huntemann, C. Grebing, B. Lipphardt, C. Lisdat, D. Piester, Direct comparison of optical lattice clocks with an intercontinental baseline of 9000 km, Opt. Lett. 39 (14) (2014) 4072–4075, https://doi.org/10.1364/0L.39.004072.
- [178] F. Riehle, Optical clock networks, Nat. Photonics 11 (1) (2017) 25–31, https://doi.org/10.1038/nphoton.2016.235.
- [179] D. Herman, S. Droste, E. Baumann, J. Roslund, D. Churin, A. Cingoz, J.-D. Deschênes, I.H. Khader, W.C. Swann, C. Nelson, N.R. Newbury, I. Coddington, Femtosecond timekeeping: slip-free clockwork for optical timescales, Phys. Rev. Appl. 9 (4) (2018) 044002, https://doi.org/10.1103/PhysRevApplied.9. 044002.
- [180] P. Delva, H. Denker, G. Lion, Chronometric geodesy: methods and applications, in: Relativistic Geodesy. Foundations and Applications, Springer, Berlin, 2019, pp. 25–85, https://www.springer.com/us/book/9783030114992.
- [181] H. Denker, L. Timmen, C. Voigt, S. Weyers, E. Peik, H.S. Margolis, P. Delva, P. Wolf, G. Petit, Geodetic methods to determine the relativistic redshift at the level of 10⁻¹⁸ in the context of international timescales: a review and practical results, J. Geod. 92 (5) (2018) 487-516, https://doi.org/10.1007/ s00190-017-1075-1.
- [182] H. Denker, Regional gravity field modeling: theory and practical results, in: G. Xu (Ed.), Sciences of Geodesy, vol. II, Springer, Berlin, Heidelberg, 2013, pp. 185–291, Chapter 5.
- [183] M. Weiss, N. Pavlis, A re-evaluation of the relativistic redshift on frequency standards at NIST, Boulder, Colorado, USA, Metrologia 54 (2017) 535, https://doi.org/10.1088/1681-7575/aa765c.
- [184] D. Calonico, A. Cina, I.H. Bendea, F. Levi, L. Lorini, A. Godone, Gravitational redshift at INRIM, Metrologia 44 (5) (2007) L44, http://stacks.iop.org/0026-1394/44/i=5/a=N03.
- [185] N.K. Pavlis, M.A. Weiss, The relativistic redshift with 3×⁻¹⁷ uncertainty at NIST, Boulder, Colorado, USA, Metrologia 40 (2) (2003) 66, http://stacks. iop.org/0026-1394/40/i=2/a=311.
- [186] C. Voigt, H. Denker, L. Timmen, Time-variable gravity potential components for optical clock comparisons and the definition of international time scales, Metrologia 53 (6) (2016) 1365, https://doi.org/10.1088/0026-1394/53/6/1365.
- [187] G. Lion, I. Panet, P. Wolf, C. Guerlin, S. Bize, P. Delva, Determination of a high spatial resolution geopotential model using atomic clock comparisons, J. Geod. 91 (6) (2017) 597–611, https://doi.org/10.1007/s00190-016-0986-6.
- [188] T.E. Mehlstäubler, G. Grosche, C. Lisdat, P.O. Schmidt, H. Denker, Atomic clocks for geodesy, Rep. Prog. Phys. 81 (6) (2018) 064401, https://doi.org/10. 1088/1361-6633/aab409.
- [189] J. Müller, D. Dirkx, S.M. Kopeikin, G. Lion, I. Panet, G. Petit, P.N.A.M. Visser, High performance clocks and gravity field determination, Space Sci. Rev. 214 (1) (2018) 5, https://doi.org/10.1007/s11214-017-0431-z.
- [190] R. Bondarescu, A. Schärer, A. Lundgren, G. Hetényi, N. Houlié, P. Jetzer, M. Bondarescu, Ground-based optical atomic clocks as a tool to monitor vertical surface motion, Geophys. J. Int. 202 (3) (2015) 1770–1774, https://doi.org/10.1093/gji/ggv246.
- [191] R. Bondarescu, M. Bondarescu, G. Hetényi, L. Boschi, P. Jetzer, J. Balakrishna, Geophysical applicability of atomic clocks: direct continental geoid mapping, Geophys. J. Int. 191 (1) (2012) 78–82, https://doi.org/10.1111/j.1365-246X.2012.05636.x.
- [192] K. Bongs, Y. Singh, L. Smith, W. He, O. Kock, D. Świerad, J. Hughes, S. Schiller, S. Alighanbari, S. Origlia, S. Vogt, U. Sterr, C. Lisdat, R.L. Targat, J. Lodewyck, D. Holleville, B. Venon, S. Bize, G.P. Barwood, P. Gill, I.R. Hill, Y.B. Ovchinnikov, N. Poli, G.M. Tino, J. Stuhler, W. Kaenders, Development of a strontium optical lattice clock for the SOC mission on the ISS, C. R. Physique 16 (5) (2015) 553–564, https://doi.org/10.1016/j.crhy.2015.03.009.
- [193] S.B. Koller, J. Grotti, S. Vogt, A. Al-Masoudi, S. Dörscher, S. Häfner, U. Sterr, C. Lisdat, Transportable optical lattice clock with 7 × 10⁻¹⁷ uncertainty, Phys. Rev. Lett. 118 (2017) 073601, https://doi.org/10.1103/PhysRevLett.118.073601.
- [194] S. Origlia, M.S. Pramod, S. Schiller, Y. Singh, K. Bongs, R. Schwarz, A. Al-Masoudi, S. Dörscher, S. Häfner, U. Sterr, C. Lisdat, Towards an optical clock for space: compact, high-performance optical lattice clock based on bosonic atoms, Phys. Rev. A 98 (5) (2018) 053443, https:// doi.org/10.1103/PhysRevA.98.053443.
- [195] J. Grotti, S. Koller, S. Vogt, S. Häfner, U. Sterr, C. Lisdat, H. Denker, C. Voigt, L. Timmen, A. Rolland, F.N. Baynes, H.S. Margolis, M. Zampaolo, P. Thoumany, M. Pizzocaro, B. Rauf, F. Bregolin, A. Tampellini, P. Barbieri, M. Zucco, G.A. Costanzo, C. Clivati, F. Levi, D. Calonico, Geodesy and metrology with a transportable optical clock, Nat. Phys. 14 (2018) 437–441, https://doi.org/10.1038/s41567-017-0042-3.
- [196] K.C. Cossel, E.M. Waxman, F.R. Giorgetta, M. Cermak, I.R. Coddington, D. Hesselius, S. Ruben, W.C. Swann, G.-W. Truong, G.B. Rieker, N.R. Newbury, Open-path dual-comb spectroscopy to an airborne retroreflector, Optica 4 (7) (2017) 724–728, https://doi.org/10.1364/OPTICA.4.000724.

- [197] J.-D. Deschênes, L.C. Sinclair, F.R. Giorgetta, W.C. Swann, E. Baumann, H. Bergeron, M. Cermak, I. Coddington, N.R. Newbury, Synchronization of distant optical clocks at the femtosecond level, Phys. Rev. X 6 (2016) 021016, https://doi.org/10.1103/PhysRevX.6.021016.
- [198] F.R. Giorgetta, W.C. Swann, L.C. Sinclair, E. Baumann, I. Coddington, N.R. Newbury, Optical two-way time and frequency transfer over free space, Nat. Photonics 7 (6) (2013) 434–438, https://doi.org/10.1038/nphoton.2013.69.
- [199] K. Djerroud, O. Acef, A. Clairon, P. Lemonde, C.N. Man, E. Samain, P. Wolf, Coherent optical link through the turbulent atmosphere, Opt. Lett. 35 (9) (2010) 1479–1481, http://ol.osa.org/abstract.cfm?URI=ol-35-9-1479.