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Authors: Franck Bourdelle [a,b], Karim Benzerara [a], Olivier Beyssac [a], Julie Cosmidis [a], Daniel Neuville [c], Gordon E. Brown Jr [d,e], Erwan Paineau [f,g]

Mailing addresses of the relevant institutions:

[a] IMPMC, UPMC-CNRS, Case courrier 115, 4 Place Jussieu, 75252 Paris cedex 05, France

[b] GeoRessources, Université de Lorraine, UMR 7359 CNRS, 54506 Vandœuvre-lés-Nancy, France

[c] Géochimie et Cosmochimie, Institut de Physique du Globe de Paris, Université Paris Diderot, Sorbonne Paris Cité, UMR 7154 CNRS, 1 rue Jussieu 75005 Paris, France

[d] Surface & Aqueous Geochemistry Group, Department of Geological and Environmental Sciences, Stanford University, Stanford, CA, 94305-2115, USA

[e] SLAC Natl Accelerator Lab, Stanford Synchrotron Radiation Lightsource, Menlo Pk, CA 94025 USA

[f] Laboratoire de Physique des Solides, Université Paris-Sud, UMR 8502 CNRS, Bâtiment 510, 91405 Orsay cedex, France

[g] CEA Saclay, IRAMIS, Laboratoire Interdisciplinaire sur l'Organisation Nanométrique et Supramoléculaire, 91191 Gif-sur-Yvette cedex, France

Corresponding authors:

1. Franck Bourdelle

Present mailing address: GeoRessources, Université de Lorraine, UMR 7359 CNRS, 54506 Vandœuvre-lés-Nancy, France

E-mail address: franck.bourdelle@gmail.com

Phone number: + 33 (0)3 83 68 47 28

2. Karim Benzerara

Mailing address: IMPMC, UPMC-CNRS, Case courrier 115, 4 Place Jussieu, 75252 Paris cedex 05, France

E-mail address: Karim.Benzerara@impmc.upmc.fr

Phone number: + 33 (0)1 44 27 75 42

Abstract

Estimation of $Fe^{3+}/\Sigma Fe$ ratios in materials at the submicrometer scale has been a long-standing challenge in the Earth and environmental sciences because of the usefulness of this ratio in estimating redox conditions as well as for geothermometry. To date, few quantitative methods with submicrometric resolution have been developed for this purpose, and most of them have used electron energy-loss spectroscopy (EELS) carried out in the UHV environment of a transmission electron microscope (TEM). Scanning transmission x-ray microscopy (STXM) is a relatively new technique complementary to TEM and is increasingly being used in the Earth sciences. Here, we detail an analytical procedure to quantify the $Fe^{3+}/\Sigma Fe$ ratio in silicates using Fe $L_{2,3}$-edge x-ray absorption near edge structure (XANES) spectra obtained by STXM, and we discuss its advantages and limitations. Two different methods for retrieving $Fe^{3+}/\Sigma Fe$ ratios from XANES spectra are calibrated using reference samples with known Fe^{3+} content by independent approaches. The first method uses the intensity ratio of the two major peaks at the L_3-edge. This method allows mapping of $Fe^{3+}/\Sigma Fe$ ratios at a spatial scale better than 50 nm by the acquisition of 5 images only. The second method employs a 2-eV-wide integration window centred on the L_2 maximum for Fe^{3+}, which is compared to the total integral intensity of the Fe L_2-edge. These two approaches are applied to metapelites from the Glarus massif (Switzerland), containing micrometer-sized chlorite and illite grains and prepared as ultrathin foils by Focused Ion Beam milling. Nanometrer-scale mapping of iron redox in these samples is presented, and shows evidence of compositional zonation. The existence of such zonation has crucial implications for geothermometry and illustrates the importance of being able to measure $Fe^{3+}/\Sigma Fe$ ratios at the submicrometer scale in geological samples.

Keywords: ferric/ferrous iron, STXM, XANES spectroscopy, $L_{2,3}$-edge, redox mapping, silicate.

Introduction

Determination of the redox state of iron and its spatial variations in sediments and rocks is of critical importance in both geosciences and environmental sciences, because of the need to understand redox state during their deposition or formation as well as subsequent

changes in redox state due to weathering and other processes (e.g., de Andrade et al. 2006; Muñoz et al. 2006; Bernard et al. 2010; Benzerara et al. 2011; Bolfan-Casanova et al. 2012; Stagno et al. 2013). In addition, quantification of $Fe^{3+}/\Sigma Fe$ ratios can yield a better insight into the chemistry of complex geological materials (e.g., Muñoz et al. 2006), or a better estimation of *P-T* conditions by geothermobarometers, when variations of the Fe^{3+} content within the crystals are taken into account (e.g., Schmid et al. 2003; de Andrade et al. 2006; Bourdelle et al. 2013a). Therefore, assessment of the $Fe^{3+}/\Sigma Fe$ ratio in minerals is an important and long-standing issue. Different techniques have been used extensively in the past for this purpose, including electron microprobe analysis (EMPA, e.g., Fialin et al. 2004), Mössbauer spectroscopy (e.g., Beaufort et al. 1992) x-ray photoelectron spectroscopy (XPS, e.g., Raeburn et al. 1997a, b), or x-ray absorption near edge structure (XANES) spectroscopy at the K edge (e.g., Waychunas et al. 1983; Bajt et al. 1994; Wilke et al. 2001, 2009; Berry et al. 2003, 2010). However, none of these methods provides spatial resolution at the few nanometers scale, which is particularly useful for studying chemical zonations patterns observed in low-temperature systems. Several studies (e.g., van Aken and Liebscher, 2002) have shown that electron energy-loss spectroscopy (EELS) carried out in a transmission electron microscope (TEM) is a powerful method for determining the redox state of iron at a submicrometer resolution. However it sometimes induces severe beam damage effects, such as electron beam-induced oxidation of iron (Lauterbach et al. 2000; Garvie et al. 2004), the effect of which can be corrected by measuring the signal as a function of time. Alternatively, XANES spectroscopy at the Fe $L_{2,3}$ edges carried out with a scanning transmission x-ray microscope (STXM) has been increasingly used in the Earth and environmental sciences to infer qualitatively $Fe^{3+}/\Sigma Fe$ ratios in geological and environmental samples at a spatial resolution of ~50-nm (e.g., Wasinger et al. 2003; Carlut et al. 2010; Lam et al. 2010; de Groot et al. 2010; Miot et al. 2011; Boulard et al. 2012). This technique has several advantages such as offering a high energy resolution (better than 0.1 eV at existing synchrotron facilities) and the possibility of maintaining samples under anoxic conditions before and during the measurement (e.g., Miot et al. 2009). However, no calibration of the STXM-based Fe $L_{2,3}$-edge XANES approach has yet been carried out, whereas calibration of the EELS approach was quantified by van Aken and Liebscher (2002). Fe $L_{2,3}$-edges result from $2p \rightarrow 3d$ electronic transitions, as shown by Wasinger et al. (2003). These authors described in detail the physical basis of Fe L edges, and showed that information about iron valency can be retrieved from XANES spectra by a multiplet calculation approach (e.g., van der Laan and Kirkman, 1992; Cressey et al. 1993). This approach is difficult to apply when dealing with

mineral phases for which we do not know the structure. Alternatively, fitting of XANES spectra with a linear combination of normalized reference spectra has been performed by Miot et al. (2009), but requires appropriate Fe^{2+} and Fe^{3+} end-member reference compounds with Fe in the same local coordination environment as in the sample of interest. Van Aken and Liebscher (2002) have shown the possibility of a third approach that they calibrated for EELS and which uses an empirical correlation between $Fe^{3+}/\Sigma Fe$ ratios and a parameter (i.e., modified integral white-line intensity ratio) which is directly retrieved from EELS spectra at the Fe $L_{2,3}$ edges, and is independent of the coordination environment of Fe to a the first-order approximation.

Here we propose an empirical approach similar to that of van Aken and Liebscher (2002) to calibrate the correlation between $Fe^{3+}/\Sigma Fe$ ratio and some parameters extracted from the STXM-derived XANES Fe $L_{2,3}$-edge spectra of reference silicate glasses and phyllosilicates. Two empirical calibrations are proposed, both of which offer a compromise between speed and accuracy of the analytical measurement. An application of this approach to ultra-thin sections of natural chlorites and micas is presented to illustrate the methodology and to further assess the range of applicability of this calibration for STXM.

Materials and Methods

Reference samples

The samples used in this study were reference synthetic silicate glasses, natural phyllosilicates and fayalite, prepared as powders or ultra-thin sections cut by Focused Ion Beam (FIB) milling. The bulk chemical compositions of the five synthetic glasses were previously determined by Magnien et al. (2004). All samples are composed of similar proportions of Si, Mg, Ca, Na and Fe. The SiO_2 and FeO contents are ~52 wt% and ~13 wt%, respectively. Bulk $Fe^{3+}/\Sigma Fe$ ratios were determined by wet chemistry, Mössbauer spectroscopy and EMPA and range from 0.09 to 0.94 (Table 1; Magnien et al. 2004). For STXM-XANES analyses, we ground these samples in deaerated and deionized water, inside an anoxic glovebox ($p(O_2) < 50$ ppm) to avoid oxidation during sample preparation.

The phyllosilicate samples have 2:1 and 2:1:1 structures and their bulk compositions were investigated previously by Joswig et al. (1986), Keeling et al. (2000), Shingaro et al. (2005), Rigault et al. (2010), and in the present study by EMP analyses. Total Fe contents vary significantly between samples and bulk $Fe^{3+}/\Sigma Fe$ ratios ranging between 0.03 and 1.0

were measured by Mössbauer spectroscopy, EXAFS (extended x-ray absorption fine structure) and/or EELS (Table 1). In addition, a fayalite sample was used as a pure Fe^{2+} reference. For STXM-XANES analyses, some samples (smectite Nau-2, chlorites GAB 42, VNI 92, VNI 114, fayalite) were prepared by grinding in deaerated and deionized water in an anoxic glovebox ($p(O_2) < 50$ ppm). Other samples (clintonite, chlorite 'prochlorite', chlorite Ch1, Ti-mica) where prepared by FIB milling.

Samples transparent to soft x-rays are needed to measure XANES spectra in the transmission mode of STXM, therefore requiring the preparation of thin samples. FIB foils were cut with a FEI Model 200 TEM FIB system at University Aix-Marseille using the protocol detailed by Heaney et al. (2001). A 30 kV Ga^+ beam operating at ~20 nA excavated the sample to a depth of 5 μm. The sample foil was then further thinned to ~80-100 nm at lower beam voltage (5 kV) and current (~100 pA), in order to remove the layer damaged by high-energy ions (Bourdelle et al. 2012).

Glarus field samples

The Glarus Alps (Switzerland) belongs to the Helvetic zone of the northern margin of the Central Alps, and was affected by low-grade metamorphism. Details about the location and composition of the samples analyzed by STXM in the present study are provided in Lahfid et al. (2010). The selected rock samples are metapelites, more or less clayey or sandy marls, with various proportions of quartz, calcite, and clay minerals. Three samples (noted Glarus GL07 13, 16, and 20, as in Lahfid et al. 2010), containing chlorites and K-deficient micas, were milled by FIB. The compositions of the chlorites and micas were obtained on the FIB foils by analytical electron microscopy analyses described elsewhere (Bourdelle et al. 2012).

XANES spectroscopy

Part of the STXM analyses were performed at the Advanced Light Source (ALS) (Lawrence Berkeley National Laboratory) on branch line 11.0.2.2 following the procedures described in Miot et al. (2009). The ALS storage ring was operated at 1.9 GeV and 500 mA current in a top-up mode. More details on the branch line 11.0.2.2 and beam characteristics are given by Bluhm et al. (2009). Stacks of images were obtained by scanning the sample in

the x–y directions of selected sample areas over the 690–730 eV energy range (Fe $L_{2,3}$-edge) using an energy increment of 0.789 eV between 690 and 705 eV, 0.10 eV in the 705–713 eV energy range, 0.19 eV in the 713–719 eV energy range, 0.155 eV in the 719-726 eV energy range, and 0.475 eV in the 726-730 eV energy range. The dwell time per pixel and energy point was 1.3 ms.

Some data (chlorite Ch1, chlorite 'prochlorite' and Ti-mica) were acquired on the Pollux beamline at the Swiss Light Source (SLS, Villigen, Switzerland). The SLS synchrotron storage ring was operated at 2.4 GeV and 300 mA current in a top-up mode during data collection, and the characteristics of the beamline are detailed by Raabe et al. (2008). Stacks were obtained over the 690–730 eV energy range (Fe $L_{2,3}$-edge) using an energy increment of 0.667 eV between 690 and 700 eV, 0.15 eV in the 700–715 eV energy range, 0.40 eV in the 715–727 eV energy ranges, and 0.89 eV in the 727-730 eV energy range. The dwell time per pixel and energy point was 3.5 ms.

At both the ALS and the SLS, focus was achieved systematically for each sample, and precision in the determination of the focus position was better than the focus depth. Image stacks were aligned and XANES spectra were derived from areas of interest using the aXis2000 software (Hitchcock, 2012). Potential beam damage caused by the incident photon beam was assessed by monitoring spectral changes at the Fe $L_{2,3}$-edge with increasing dwell times up to a hundred milliseconds (10, 50, and 100 ms).

Spectra processing

Energy calibration was performed using the gaseous neon $1s \rightarrow 3p$ electronic transition at 867.3 eV. As explained in Figure 1, the processing of spectra consisted of two steps. First, the contribution of lower energy absorption edges (i.e., background) was removed so that in the end, the pre-edge region is set to 0 optical density (noted OD) with a slope of zero. For that purpose, a "linear background" correction was applied to the spectrum. Second, the two edge steps resulting from transitions to unoccupied states in the continuum were subtracted using the following double arctan function (Chen et al. 1995; van Aken and Liebscher, 2002; Brotton et al. 2007):

$$f(\Delta E) = \frac{h_1}{\pi}\left(\arctan\left[\frac{\pi}{w_1}(\Delta E - E_1)\right] + \frac{\pi}{2}\right) + \frac{h_2}{\pi}\left(\arctan\left[\frac{\pi}{w_2}(\Delta E - E_2)\right] + \frac{\pi}{2}\right) \qquad (1)$$

where h_1 and h_2 are the step heights of the two arctan functions, w_1 and w_2 are fixed peak widths, and E_1 and E_2 are the positions of the inflection points resulting in an energy near the edge onset. Here, w_1 and w_2 are fixed to 1 eV (Fig. 1). Brotton et al. (2007) proposed setting the function slope w at 5 eV, to account for the slow onset of the continuum. They argued that a value smaller than 5 eV could induce spurious structures in the background-corrected spectrum. We observed that values of w = 1 eV or w = 5 eV provided similar results.

Results and Discussion

Evolution of Fe $L_{2,3}$–edge XANES spectra with changes in $Fe^{3+}/\Sigma Fe$

XANES spectra at the Fe $L_{2,3}$-edges of the reference phyllosilicates, fayalite and five Fe-bearing silicate glasses, corrected for continuum absorption, are shown in Figure 2, and the positions of major peaks are summarized in Table 1. These spectra are qualitatively similar to those described in several previous studies and were obtained using different analytical techniques (e.g., Crocombette et al. 1995; Heijboer et al. 2003; van Aken and Liebscher, 2002). Four major Fe $L_{2,3}$-edge XANES peaks are present in all samples. The two major peaks on the L_3 edge are noted as "L_3-a" and "L_3-b", and similarly, the major peaks on the L_2 edge are noted as "L_2-a" and "L_2-b". For all samples examined, the measured separations of the Fe L_3 and L_2 maxima, due to spin-orbit splitting (van Aken and Liebscher, 2002), are 12.9 ± 0.4 eV and 14.2 ± 1.4 eV for peaks a and b, respectively in agreement with previous EELS and XANES studies (e.g., de Smit et al. 2008; de Groot et al. 2010). However, although most of the spectra show a single asymmetrical L_3-a peak, some of them (i.e., VNI 92, VNI 114 and GAB 42, fayalite) display an "L_3-a" split into two peaks. In addition, these specific spectra show additional peaks on the L_3-a side at ~706.3 and ~706.8 eV. According to Wasinger et al. (2003), the presence of these minor peaks may be due to a specific atomic environment and/or orbital co-valency of iron in these mineral phases. Similarly, several minor peaks can be observed at around 719.8 eV on the L_2-edge for several samples (VNI 92, VNI 114, GAB 42, PyrNa 17R, fayalite).

The relative intensities of the different major peaks vary depending on the $Fe^{3+}/\Sigma Fe$ ratio (Figure 2). With increasing $Fe^{3+}/\Sigma Fe$ ratios, the relative intensity of the L_3-a peak decreases compared to that of the L_3-b peak; L_3-a is more intense than L_3-b in the XANES spectrum of the VNI 92 sample ($Fe^{3+}/\Sigma Fe$ = 0.35), whereas the opposite is observed for PyrNa

5R ($Fe^{3+}/\Sigma Fe$ = 0.61). Likewise, the relative intensity of L_2-a progressively decreases whereas that of L_2-b increases as $Fe^{3+}/\Sigma Fe$ increases. The energy position of L_2-b changes very little between the samples, whereas peaks L_3-a and L_2-a shift slightly towards higher energies when $Fe^{3+}/\Sigma Fe$ increases.

Quantification of $Fe^{3+}/\Sigma Fe$ from XANES Fe $L_{2,3}$-edge intensity ratios

As documented in Figure 2, the main variations in the XANES spectra of reference samples with varying $Fe^{3+}/\Sigma Fe$ ratios involve the L_3-b/L_3-a intensity ratio. More precisely, the L_3-b/L_3-a intensity ratio is linearly correlated with $Fe^{3+}/\Sigma Fe$ ratio with only a little scatter (R^2 = 0.96) for both the phyllosilicates and silicate glasses (Figure 3). The correlation is described by equation (2):

$$\frac{Fe^{3+}}{\sum Fe} = \frac{R_{L_3} - 0.1867}{0.01991} \qquad \text{with} \qquad R_{L_3} = \frac{I(L_3 - b)}{I(L_3 - a)} \qquad (2)$$

This approach requires only five XANES images to map $Fe^{3+}/\Sigma Fe$ (see Figure 4 and below): two images in the pre-edge (needed to apply the "linear background correction" at each pixel of the image), one at 708.7 eV to quantify the L_3-a peak, one at 710.25 eV to quantify the L_3-b peak, and one at 718 eV, to remove the edge step of the arctan function. Finally, the ratio of the resulting 708.7 and 710.25 eV images can be used to determine the R_{L3} parameter at each pixel of the image.

This calibration is useful but has some limitations. The L_3 peaks, which are much more intense than the L_2 peaks, are more susceptible to absorption saturation (see Hanhan et al. 2009, where saturation effects are described for Ca 2p edge spectra). This phenomenon occurs when the sample is too thick and/or highly concentrated in Fe implying that few photons are transmitted. This may trigger a non-linear response of the detection and an artifactual modification of the relative peak heights. The use of a spectral parameter correlated with $Fe^{3+}/\Sigma Fe$ based on the less absorbing L_2-edge may provide in this case an interesting way of circumventing absorption saturation issues encountered with the L_3-edge.

Figure 2 shows that an increase of $Fe^{3+}/\Sigma Fe$ is associated with a decrease of the intensity of L_2-a. Figure 5 shows the correlation between $Fe^{3+}/\Sigma Fe$ and R_{L2}, a ratio that reflects the importance of L_2-b relative to the total L_2. Similar to the modified intensity defined by van Aken and Liebscher (2002), the L_2-b contribution is computed as an integration window of 2

eV width centred around the maximum L_2-b intensity; the ratio R_{L2} is calculated from this modified L_2-b intensity and the total integral intensity of L_2-edge. The correlation is high (R^2 = 0.97), and is described by equation (3):

$$\frac{Fe^{3+}}{\sum Fe} = \frac{R_{L_2} - 0.1476}{0.00297} \qquad \text{with} \qquad R_{L_2} = \frac{I(L_2 - b)_{\text{mod}}}{A(L_2)_{total}} \qquad (3)$$

This approach requires the acquisition of a complete stack of images (i.e., as many images as energy points are required to obtain a complete spectrum with a given spectral resolution) between 715 and 730 eV, to cover the entire L_2-edge, and to calculate the double arctan function (equation (1)). As a consequence, the acquisition time required for this method is longer than for the L_3-b/L_3-a intensity ratio method (e.g. 30-40 min *versus* 5-10 min for an area of 150 by 150 pixels). However, this second method seems to be more accurate, especially because (1) the calibration data are less scattered (Fig. 5 *versus* Fig. 3) and (2) the intensity integration improves the signal-to-noise ratio. Several other methods of calibration have been tested, sometimes giving equation with a high correlation (with a R^2 up to 0.95), but the two methods proposed here seem to be a good trade-off between $Fe^{3+}/\Sigma Fe$ estimation accuracy, acquisition time and ease of use.

Assessment of saturation and beam damage effects

When particles are sufficiently thin, the intensity of each spectral feature changes linearly with thickness. However, Hanhan et al. (2009) showed that in the case of samples that are too thick, one can observe distortions of the Ca 2p spectrum due to a saturation effects. These observations led the authors of that study to set a maximum peak intensity, which should not be exceeded to avoid saturation phenomena.

Similarly, we determined the maximum peak intensity below which the Fe L_{23} spectra are undistorted and vary linearly. For this purpose, a powder of the smectite Nau-2 sample with grains of various sizes was analyzed by STXM. Figure 6 plots the difference between the intensities at 710.35 (L_3-b) and 723.54 eV (L_2-b) (corrected from the pre-edge intensity) vs. the intensity at 710.35 eV (L_3-b, i.e., the peak of maximum intensity for Nau-2, hence the most susceptible to saturation) for each pixel of the stack of images (i.e., a total of 6336 pixels). The difference between L_3-b and L_2-b intensities increases linearly when L_3-b intensity is lower than ~1.5 OD. Once the L_3-b intensity exceeds 1.5 OD, the L_3-b – L_2-b

difference increases more slowly than L_3-b, underlining (i) the distortion of the spectra for the considered pixels, and (ii) the faster increase of L_2-b intensity compared to that of L_3-b with increasing sample thickness. All the data presented in this study were therefore collected from areas presenting a L_3 peak intensity lower than 1.5 OD.

The spectrum may be also influenced by the crystal orientation relative to the direction of polarization of the x-ray beam, a process called linear dichroism. Therefore several XANES spectra were measured on the same part of a FIB foil after sequential rotation of the linear polarisation (see Benzerara et al., 2011 for details on the procedure). The variation of resulting $Fe^{3+}/\Sigma Fe$ estimates is negligible, showing that sample orientation does not affect the Fe^{3+} quantification.

Beam damage can also potentially alter assessment of the $Fe^{3+}/\Sigma Fe$ ratio. Here, beam damage was evaluated by monitoring spectral changes at the Fe $L_{2,3}$-edge with increasing dwell times from 10 up to 100 milliseconds. Figure 7 shows that $Fe^{3+}/\Sigma Fe$ ratios derived from XANES spectra are only slightly affected by increasing dwell time. In particular, no significant change was observed for typical dwell times used during routine analyses of the samples (i.e., ~1.3 and 3.5 ms per energy- and image-point for ALS and SLS synchrotrons, respectively).

Application to a geological case: chlorites and micas from Glarus (Central Alps, Switzerland)

To go further, we have applied the methods proposed here on micrometer-sized chlorite and mica/illite-like grains sampled in the Glarus area of Switzerland and cut by FIB-milling. The temperatures of chlorite formation were calculated from analytical electron microscopy (AEM) chemical analyses, based on the thermometer by Bourdelle et al. (2013b), which does not require $Fe^{3+}/\Sigma Fe$ input, and the thermometer by Inoue et al. (2009), which needs a previous estimation of Fe^{3+} content. The $Fe^{3+}/\Sigma Fe$ ratios were estimated for each FIB foil by XANES from equations (2) and (3). The results are given in Figure 8 and Table 2.

From images converted to optical density units taken at 708.7 eV, we can easily distinguish Fe-rich and Fe-poor minerals: chlorites appear as light grey and represent the Fe-rich phase, whereas micas are dark, i.e., Fe-poor. XANES spectra, acquired along the mica-chlorite contacts show that the $Fe^{3+}/\Sigma Fe$ ratio is higher in illite than in chlorite: the $Fe^{3+}/\Sigma Fe$ ratios estimated by equation (3) range from 22.3% to 27.9% in chlorite, whereas these ratios vary between 30 and 65.5% in illite-like phase. Equation (2) provides consistent estimations, suggesting that both calibrations are reliable. This analysis shows that K-deficient micas can

contain a high proportion of ferric iron (e.g. samples 13 and 20). Despite the relatively high $Fe^{3+}/\Sigma Fe$ ratio in some illite-like crystals, the total Fe^{3+} content remains higher in chlorite.

Figure 8 also shows the variations of $Fe^{3+}/\Sigma Fe$ ratios vs. the temperature of formation, which was estimated by chlorite thermometry (Table 2). In this respect, $Fe^{3+}/\Sigma Fe$ ratio increases slightly in chlorites with increasing temperature, whereas this ratio decreases in K-deficient micas. It should be noted that, contrarily to the Bourdelle et al. (2013b) model, some geothermometers based on thermodynamic models for chlorite (e.g., Inoue et al. 2009), require prior determination of the $Fe^{3+}/\Sigma Fe$ ratio. When this value in not known, it is set to zero as the default in these types of models. Interestingly, the comparison of results provided by different thermometers in Table 2 shows that the Inoue and Bourdelle geothermometers yield very different temperature results (differences of up to 76°C) when $Fe^{3+}/\Sigma Fe$ ratio is not known. In contrast, taking into account the $Fe^{3+}/\Sigma Fe$, the two thermometers provide more similar temperatures estimates (a maximum difference of less than 28°C, i.e., within the uncertainty of the thermometers), showing the cross-check validity of the $Fe^{3+}/\Sigma Fe$ estimation. A variation of the $Fe^{3+}/\Sigma Fe$ ratio from 0 to ~23% in chlorites implies a decrease in the temperatures calculated by the Inoue model of 20, 40, and 46°C depending on the sample.

Figure 9 displays an example of $Fe^{3+}/\Sigma Fe$ mapping at the nanometer-scale derived from images at 706, 708.7, 710.25, and 718 eV using Eq. (2) (see Figure 4). The analysis was carried out on the Glarus GL07 20 FIB foil. The scanned area measures 3.3 x 3.5 micrometers with a pixel size of 88 nm x 88 nm. The analysis of the illite-chlorite contacts by AEM showed that they are approximately perpendicular to the FIB foil surface, i.e., there is only a little overlap between the two minerals at their contact. The spatial averaging effect of the x-ray beam over the pixel size (i.e., 88 nm) sets the limit of the minimum distance over which illite-chlorite contacts can be discriminated. Beyond this distance, the intracrystalline variation of the $Fe^{3+}/\Sigma Fe$ ratio in the illite-like phase can be interpreted as an authentic zonation, from ~55% in crystal rims (conforming to the spectra presented in Figure 8) to ~85% in several crystal core clusters. In the same way, the $Fe^{3+}/\Sigma Fe$ ratio distribution draws a subtle zonation in the chlorite, with a $Fe^{3+}/\Sigma Fe$ ratio ranging from 18 to ~23% on the crystal rim, in accordance with the spectra shown in Figure 8. Such variations of the $Fe^{3+}/\Sigma Fe$ ratio within the crystals are equivalent to several degrees or tens of degrees in the temperature estimation, especially when this variation is associated with a variation in composition. One can expect that this zonation is a crucial issue in application of geothermometers (de Andrade et al. 2006; Bourdelle et al. 2013a), and the redox gap between illite and chlorite raises the issue of the crystallisation processes.

In summary, the STXM-based XANES study of FIB foils from the Glarus, Switzerland samples enables (i) estimation of the $Fe^{3+}/\Sigma Fe$ ratio in each phase preserving the mineral texture, and (ii) establishment of iron redox mapping with high spatial resolution.

Conclusion

In this study, we have demonstrated the reliability of two methods that allow quantitative determination of $Fe^{3+}/\Sigma Fe$ ratios in silicate phases using STXM coupled with XANES spectroscopy at the Fe $L_{2,3}$-edges. These approaches are similar to those proposed by van Aken and Liebscher (2002) for EELS measurements but are here calibrated for STXM. The two calibrations are based on reference samples with variable but known $Fe^{3+}/\Sigma Fe$ ratios, which were prepared as powders or as FIB foils. We tested these calibrations on three FIB foils extracted from field samples of phyllosilicates (Glarus, Switzerland chlorite and illite samples from metapelites), demonstrating the potential of these methods for quantifying $Fe^{3+}/\Sigma Fe$ ratios at the submicrometer-scale. This approach will allow more quantitative mineralogical or geomicrobiological studies requiring estimation of the iron redox state at the nanoscale for terrestrial or extraterrestrial Fe-rich samples.

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550 **Tables**

551

552 **Table 1** Reference samples used for XANES - $Fe^{3+}/\Sigma Fe$ ratio quantification

553

Type	Sample	Source	FeO %wt.*	$Fe^{3+}/\Sigma Fe$ Redox ratio	Position of maximum peak intensity (eV)			
					L_3-a	L_3-b	L_2-a	L_2-b
Silicate glass	PyrNa	Magnien et al. (2004)	12.83	0.73 ±0.03	708.57	710.04	721.68	723.54
Silicate glass	PyrNa17R	Magnien et al. (2004)	12.75	0.09 ±0.01	708.36	710.77	721.37	723.54
Silicate glass	PyrNa5R	Magnien et al. (2004)	12.92	0.61 ±0.01	708.57	710.04	721.68	723.54
Silicate glass	PyrNa750	Magnien et al. (2004)	12.68	0.94 ±0.03	708.67	710.25	721.68	723.54
Silicate glass	PyrNa1200	Magnien et al. (2004)	13.52	0.89 ±0.05	708.67	708.93	721.84	724.54
Nesosilicate	Fayalite	Neuville D.	70.50	0.00 ±0.00‡	708.09	710.71	720.75	723.54
Phyllosilicate	Smectite Nau-2	Keeling et al. (2000)	34.10	1.00 ±0.00‡	708.57	710.35	721.84	723.54
Phyllosilicate	Clintonite	Joswig et al. (1986)	3.01	0.69 ±0.03‡	708.25	710.08	721.49	723.53
Phyllosilicate	Ti-mica	Shingaro et al. (2005)	19.38	0.03 ±0.03‡	708.04	710.04	721.00	723.54
Phyllosilicate	Chlorite 'Prochlorite'	MNHN†	14.50	0.30 ±0.10‡	708.26	710.21	721.51	723.51
Phyllosilicate	Chlorite Ch1	This study	40.10	0.17 ±0.05	708.04	710.19	721.00	723.54
Phyllosilicate	Chlorite GAB 42	Rigault (2010)	28.50	0.14 ±0.03	708.15	710.46	720.75	723.54
Phyllosilicate	Chlorite VNI 92	Rigault (2010)	20.73	0.35 ±0.03	708.15	710.25	720.75	723.54
Phyllosilicate	Chlorite VNI 114	Rigault (2010)	20.82	0.20 ±0.03	708.04	710.25	720.75	723.54

554 * ΣOxides wt% = 100 as basis and all iron reported as ferrous. † MNHN: Collection of Muséum National
555 d'Histoire Naturelle, France. ‡ assumed error deviation.

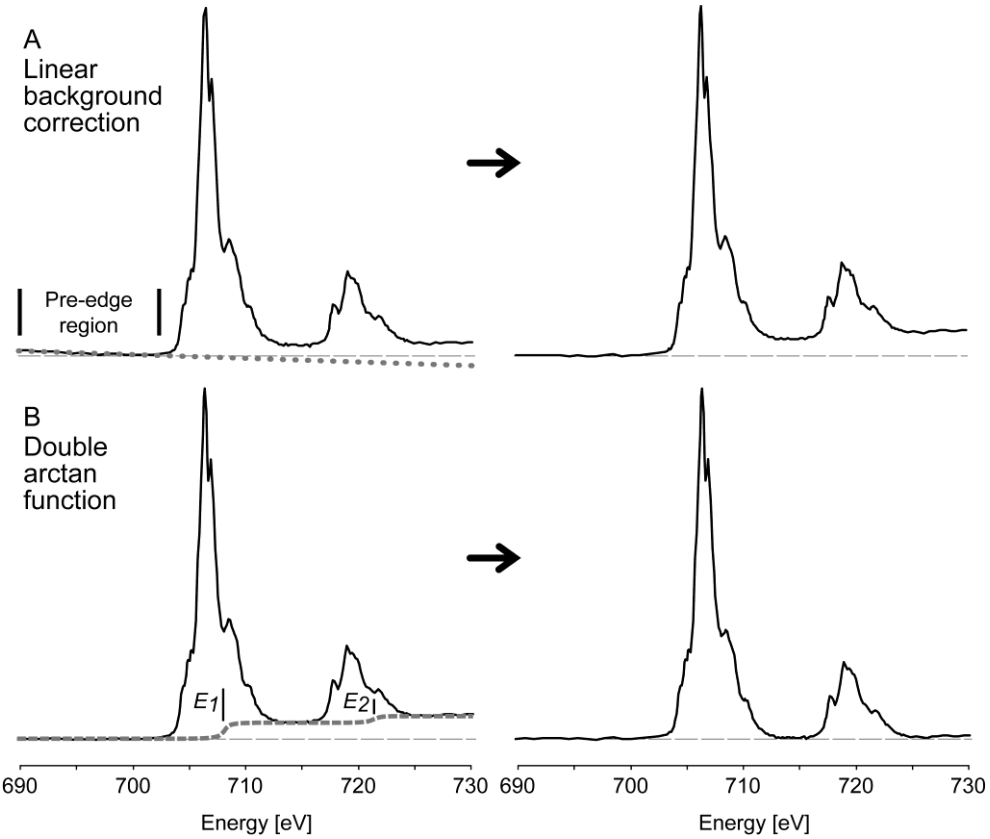
556

Table 2 AEM chemical representative analyses of Glarus chlorites and comparison of thermometers results taking into account $Fe^{3+}/\Sigma Fe$ ratios (Eq. (3)); regarding the scatter of data on Figure 3, we infer a precision of ±5% on the $Fe^{3+}/\Sigma Fe$. Analyses were carried out on crystal rims, along the illite-chlorite contact, according to Bourdelle et al. (2013a). Temperature estimations were performed with Bourdelle et al. (2013b) and Inoue et al. (2009) thermometers, with and without consideration of Fe^{3+}, in accordance with the recommendations made by each authors. Taking into account the $Fe^{3+}/\Sigma Fe$ ratios in the Inoue model allows to obtain similar results to those calculated with the Bourdelle model (which is a pure Fe^{2+} model), i.e., with a difference less than 30°C (equivalent to the uncertainty of each model)

Chlorite [%wt]	GL07 13	GL07 20	GL07 16
SiO_2	32.70	31.40	31.57
TiO_2	0.07	0.06	0.04
Al_2O_3	26.41	25.56	26.02
FeO	26.63	28.53	25.13
MnO	0.00	0.00	0.00
MgO	13.75	13.97	16.70
CaO	0.06	0.10	0.29
Na_2O	0.00	0.04	0.00
K_2O	0.39	0.35	0.25
$Fe^{3+}/\Sigma Fe$ [%]	22.30	23.80	27.90
T_{BOU} [°C]	135	170	182
$T_{INO\text{-}Fe2+}$ [°C]	165 (*+30*)	236 (*+66*)	258 (*+76*)
$T_{INO\text{-}Fe3+}$ [°C]	145 (*+10*)	196 (*+26*)	210 (*+28*)

T_{BOU}: temperatures calculated with the Fe^{2+}-pure model of Bourdelle et al. (2013b), considering $Fe_{tot} = Fe^{2+}$. $T_{INO\text{-}Fe2+}$: temperatures calculated with the Inoue et al. (2009) model, considering $Fe_{tot} = Fe^{2+}$. $T_{INO\text{-}Fe3+}$: temperatures calculated with the Inoue et al. (2009) model, using the estimated $Fe^{3+}/\Sigma Fe$.

Figures



574

575 **Fig. 1** Subtraction of background from XANES spectra at Fe L-edge, using linear and double

576 arctan functions ($w_1 = w_2 = 1$ eV), for chlorite GAB 42

577

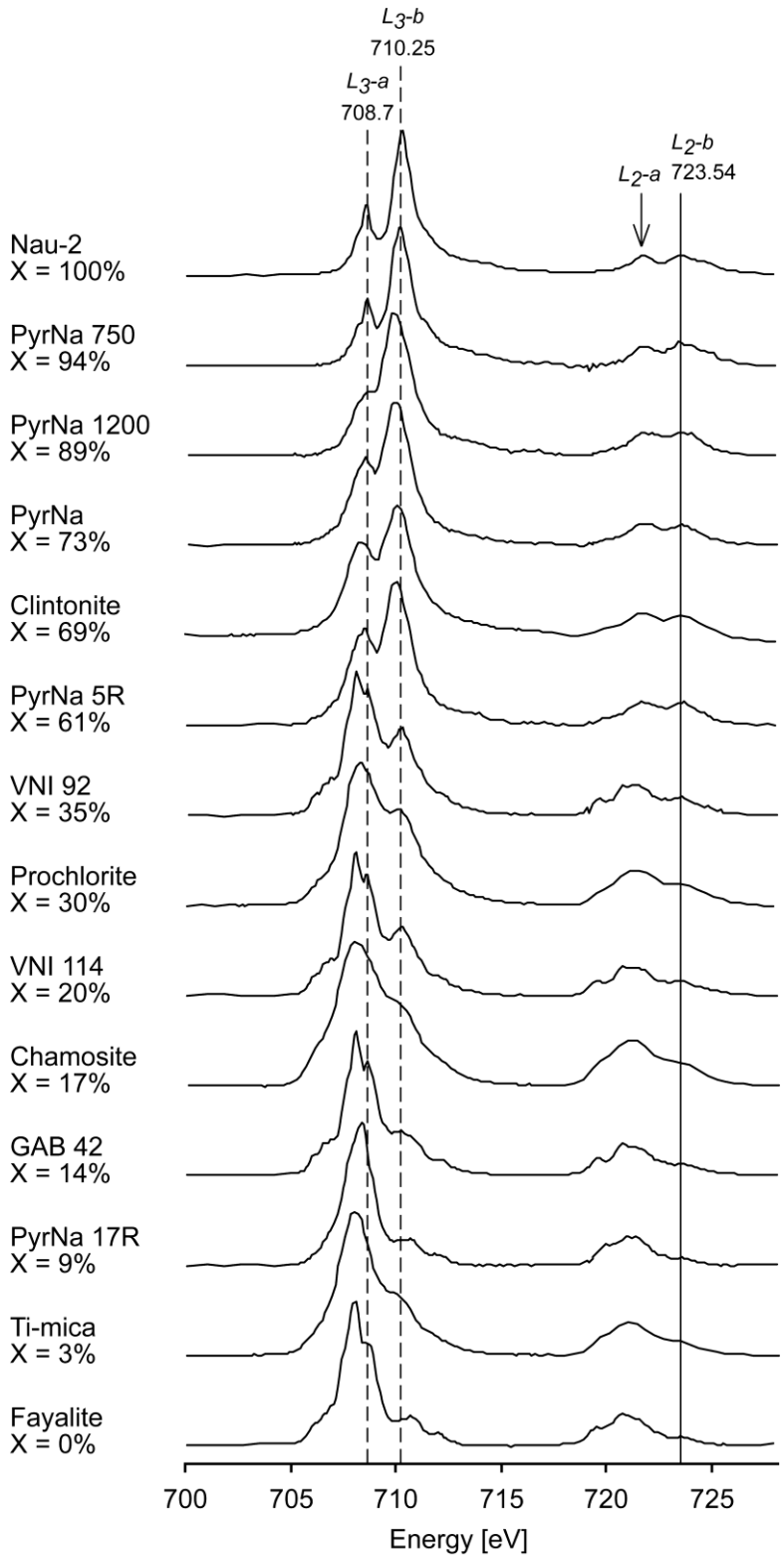


Fig. 2 Representative XANES spectra at the Fe $L_{2,3}$-edges for the reference silicates. The spectra have been normalised to the integral Fe L_3-edge intensity, and some of the spectra have been shifted vertically for clarity (normalized intensity with arbitrary units). The dotted lines represent the energies fixed to determine the Fe^{3+} concentration from the Fe L_3-peaks'

583 intensity ratio. The solid line underlines the position of L_2-b maximum intensity, which is
584 identical for all spectra. X = $Fe^{3+}/\Sigma Fe$ ratios of Table 1
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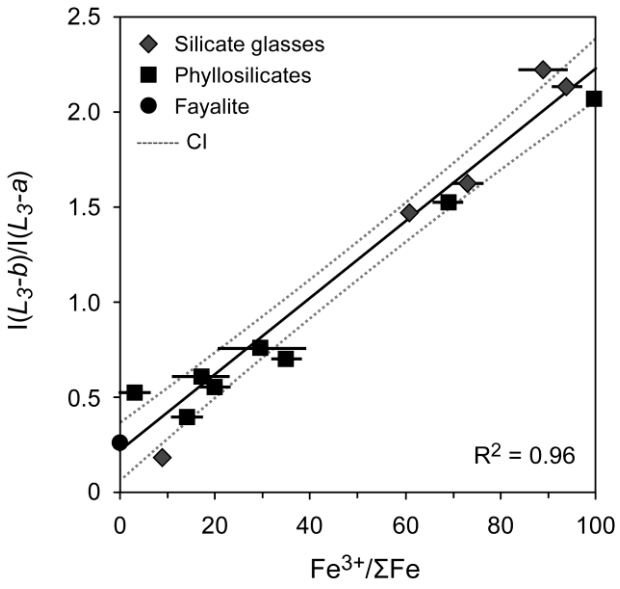
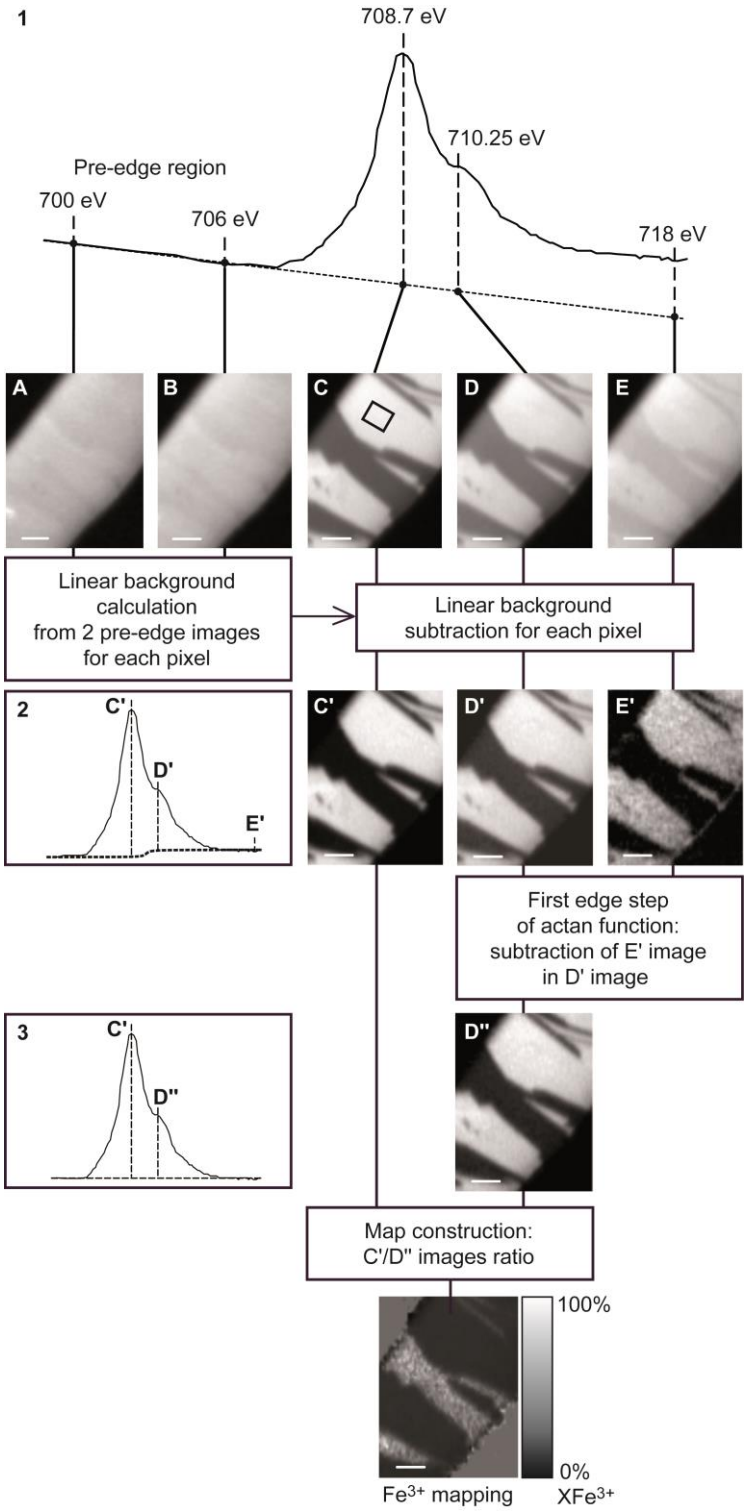


586

587 **Fig. 3** L_3-edge intensity ratio *I(L_3-b)/I(L_3-a)* from XANES spectra versus ferric iron
588 concentration $Fe^{3+}/\Sigma Fe$ quantified by independent methods for the selected silicates. CI:
589 confidence interval (95%)
590



591

592 **Fig. 4** Determination of the $Fe^{3+}/\Sigma Fe$ ratio from 5 selected energy images: two images in the
593 pre-edge (to apply the "linear background correction" at each pixel of the image), one at 708.7
594 eV to quantify the L_3-a peak, one at 710.25 eV to quantify the L_3-b peak, and one at 718 eV,
595 to remove the edge step of the arctan function. Finally, the ratio of the resulting 708.7 and
596 710.25 eV images can be used to determine the R_{L3} parameter at each pixel of the image, and
597 obtain iron redox mapping. All images are OD images (70 x 90 pixels), where the illite and

chlorite are the dark- and light- grey phases, respectively. As an illustration, spectrum #1 was retrieved from 110 images (i.e., 110 energy points) on a chlorite area (dark rectangle on image C); spectrum #2 was obtained after the linear function subtraction from spectrum #1 and spectrum #3 after the actan function subtraction from spectrum #2. Case of FIB foil of Glarus GL07 20 sample.

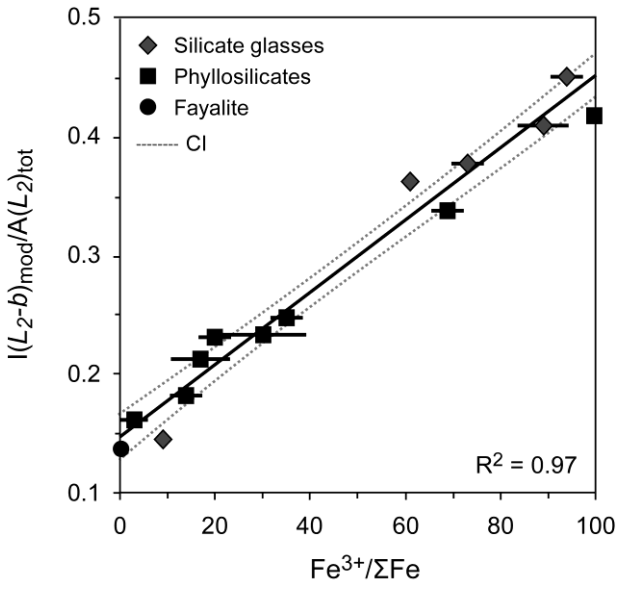
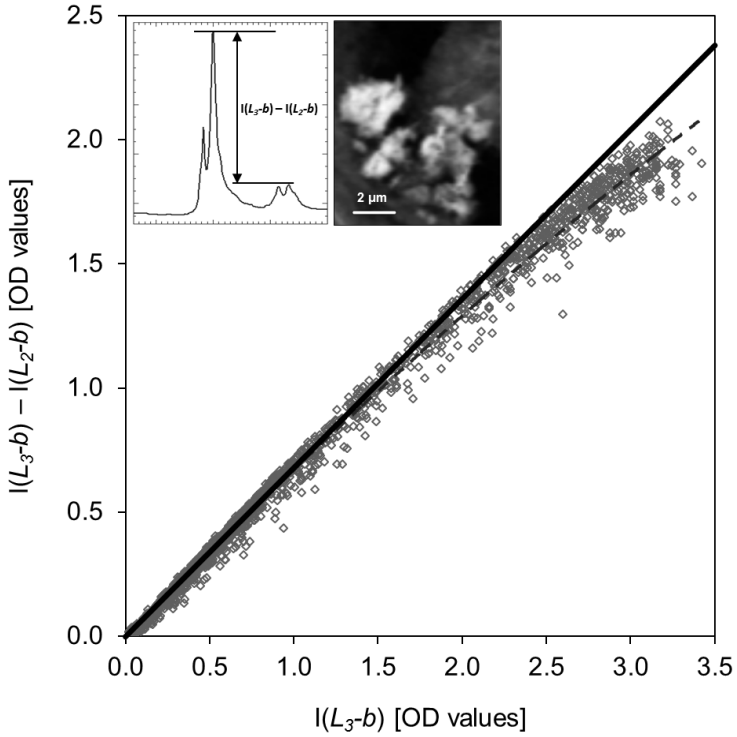


Fig. 5 L_2-edge integral intensity ratio (i.e., integral intensity of maximum L_2-b ±0.1 eV over total integral intensity (area) of L_2-edge) from XANES spectra versus ferric iron concentration $Fe^{3+}/\Sigma Fe$ quantified by independent methods for the reference silicates. CI: confidence interval (95%)



610

611 **Fig. 6** Difference, pixel by pixel, of intensity detected between the 710.25 and the 723.54 eV
612 images (in which a pre-edge image was not subtracted) versus the intensity of the 710.25 eV
613 image of a Smectite Nau-2 STXM-map (Nau-2, 72 x 88 pixels = 6336 points), i.e., the L_3-b –
614 L_2-b intensity difference versus the L_3-b intensity for each pixel. The dashed line was
615 calculated from a quadratic equation. Insets: representative spectra and optical density image
616 (710.25 eV) for Nau-2 sample

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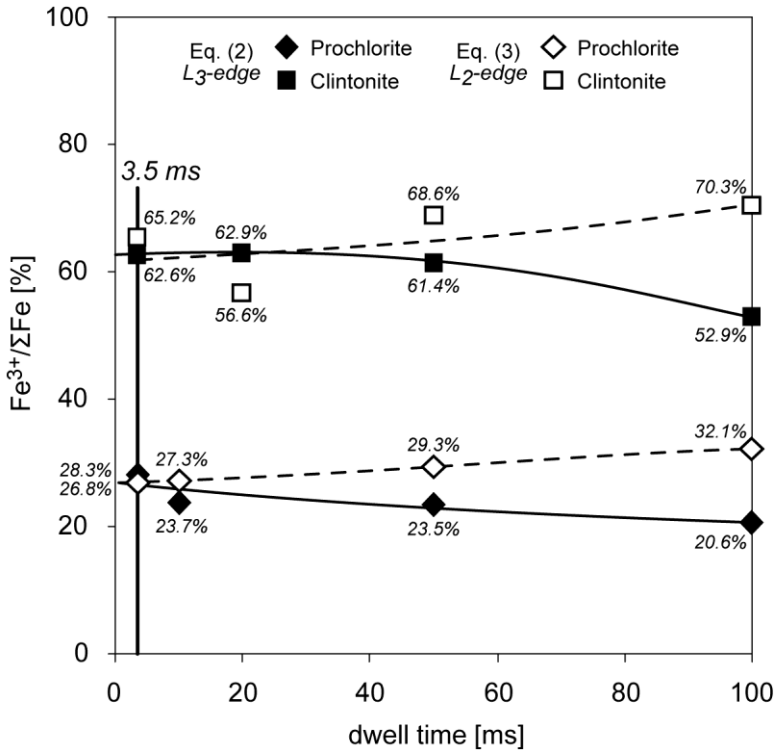

Fig. 7 Beam-induced radiation damage during STXM analyses of chlorite 'prochlorite' (XFe^{3+} = 30%) and clintonite (XFe^{3+} = 69%). Evolution of the $Fe^{3+}/\Sigma Fe$ ratios as a function of dwell time, estimated by Eq. (2) and (3) from XANES spectra. Data were fit by a quadratic function. The beam radiations (increasing dwell) involve (1) a decrease of XFe^{3+} calculated from L_3-edge (Eq. 2) and (2) an increase of XFe^{3+} calculated from L_2-edge (Eq. 3). Spectra of reference samples and Glarus samples (see text) were recorded with a dwell time of 1.3-3.5 ms per point and energy: the beam radiation damage is thus negligible with our analytical conditions for data collection

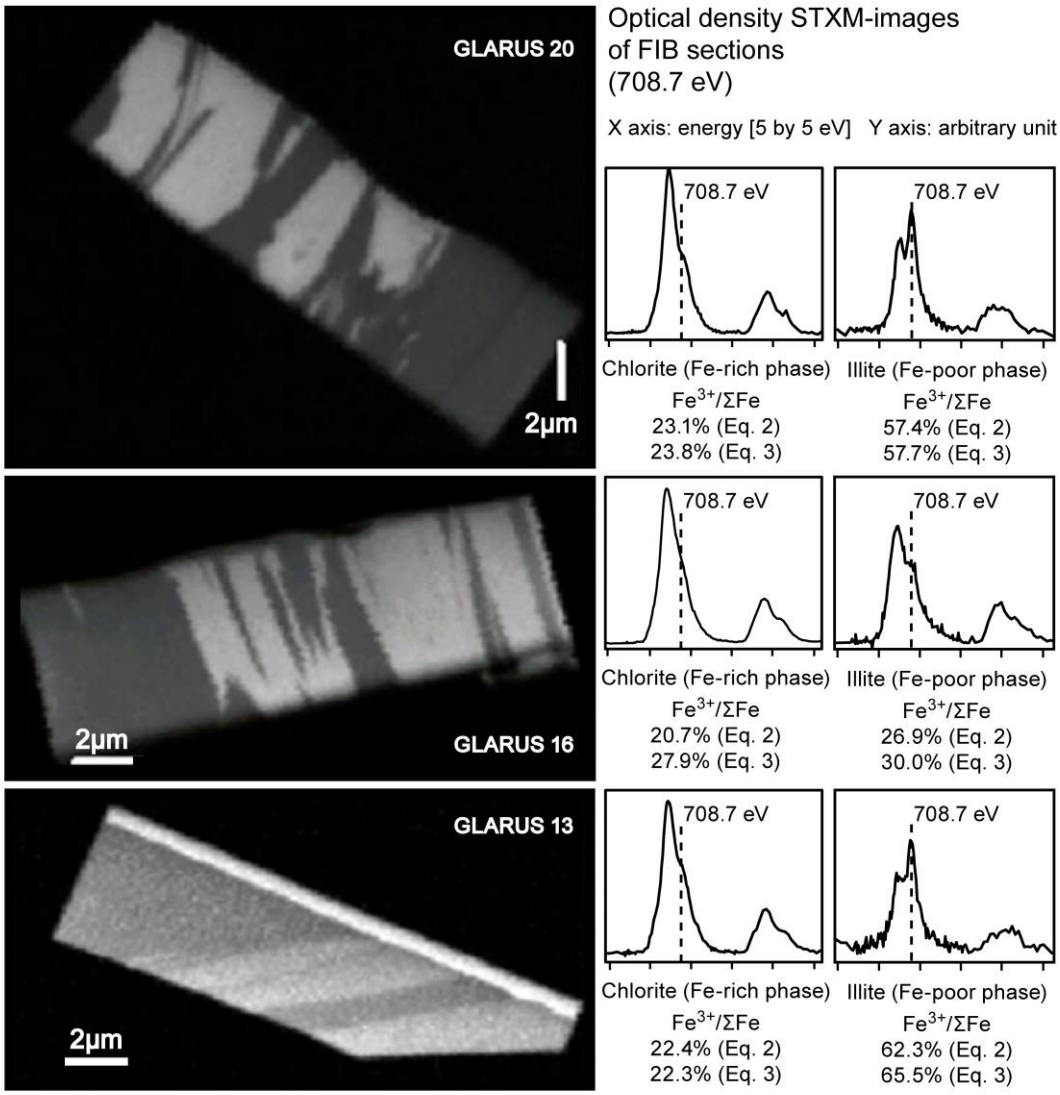

Fig. 8 Scanning transmission x-ray microscopy (STXM) and XANES analysis and $Fe^{3+}/\Sigma Fe$ estimations for FIB foils of Glarus samples (chlorite and illite). [Left] Optical density images of FIB foils at 708.7 eV. The illite and chlorite are the dark- and light- grey phases, respectively. [Right] XANES spectra of areas of interest and calculated Fe^{3+} concentrations associated (crystals rims)

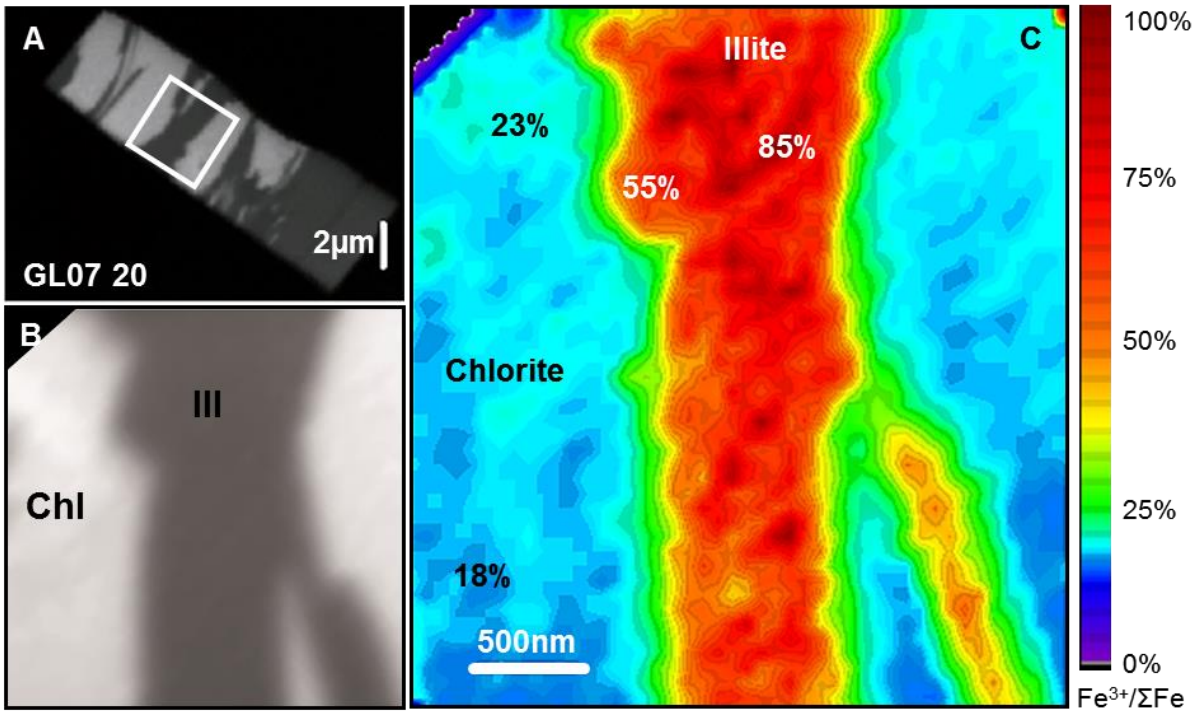

Fig. 9 Quantitative Fe redox nanomapping on FIB foil of Glarus GL07 20 sample. (a) Optical density image at 708.7 eV of Glarus GL07 20 FIB section, where the illite and chlorite are the dark- and light- grey phases, respectively. (b) Optical density image at 708.7 eV of the area of interest. (c) iron redox mapping, calculated from the 708.7 to 710.25 eV images ratio coupled with Eq. (2). The illite-chlorite contacts were analysed by AEM to check that they are approximately perpendicular to the FIB foil surface, i.e., there is only a small overlap between the two minerals at their contact. The spatial averaging effect of the x-ray beam over the pixel size (i.e., 88 nm) sets the limit of the minimum distance over which illite-chlorite contacts can be discriminated. Beyond this distance, the intracrystalline variation of $Fe^{3+}/\Sigma Fe$ ratio in the illite-like phase can be interpreted as an authentic zonation, from 55% to 85% in several crystal core clusters