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Dipolar couplings in solid polypeptides probed by ^{14}N NMR spectroscopy

Diego Carnevale¹, Benoit Grosjean² & Geoffrey Bodenhausen¹

The acquisition of ^{14}N NMR spectra in solid samples is challenging due to quadrupolar couplings with magnitudes up to several MHz. This nucleus is nonetheless important as it is involved in the formation of essential secondary structures in biological systems. Here we report the structural study of the atomic environment of amide functions in polypeptides using magic-angle spinning NMR spectroscopy of the ubiquitous ^{14}N isotope. The cyclic undecapeptide cyclosporin, in which only four hydrogen atoms are directly bound to nitrogen atoms, is chosen for illustration. Structural details of different environments can be revealed without resorting to isotopic enrichment. The network of inter- and intra-residue dipolar couplings between amide ^{14}N nuclei and nearby protons can be probed and mapped out up to a tunable cutoff distance. Density functional theory calculations of NMR quadrupolar interaction tensors agree well with the experimental evidence and allow the unambiguous assignment of all four non-methylated NH nitrogen sites and neighboring proton nuclei.

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Solid-state NMR is an invaluable tool to obtain structural details of arbitrary materials with atomic resolution. Inhomogeneous interactions such as anisotropic chemical shifts, dipolar and quadrupolar couplings reveal a wealth of information on the local environment on an atomic scale¹. In contrast to diffraction techniques that rely on long-range order, NMR can provide valuable information about amorphous systems. Furthermore, incoherent dynamic processes such as local motions and chemical exchange effects can be identified². Compositional and positional disorder in materials may also be investigated³. In contrast to NMR studies in solution, where the linewidths increase with the size of the dissolved molecule, solid-state NMR may be applied to proteins and assemblies without any restriction on mass, even in their native biological environments. Heterogeneous matrices such as bone tissue are also suitable for analysis with this powerful technique⁴.

In applications to biomolecules, it is common practice to enrich the ¹³C and ¹⁵N isotopes. The former nucleus is only 1.1% abundant, the remaining 98.9% ¹²C nuclei with spin quantum number $I = 0$ being NMR silent. The ¹⁵N isotope has an even lower natural abundance (0.4%), however, the highly abundant ¹⁴N isotope (99.6%) can be observed by NMR. Unfortunately, observation of this nucleus with $I = 1$ is particularly challenging due to (i) the absence of a favorable central transition $m_I = +1/2 \leftrightarrow m_I = -1/2$ that occurs in half-integer quadrupolar spins ($I > 1$) and (ii) the broadening due to quadrupolar couplings which affects the two observable single-quantum transitions $m_I = \pm 1 \leftrightarrow m_I = 0$ to both first and second order, resulting in inhomogeneously broadened lineshapes which typically span several MHz. Therefore, the possibility of acquiring ¹⁴N spectra routinely would clearly revolutionize the use of NMR spectroscopy in systems where isotopic enrichment of ¹⁵N nuclei is very challenging, exceedingly expensive, or indeed impossible.

Over the past decade or so, a number of experimental methods have been proposed to acquire ¹⁴N NMR spectra in the solid state^{5–8}. The intrinsic low sensitivity and efficiency of these approaches has limited most applications to relatively simple chemical systems such as single amino acids like glycine, histidine or alanine. More challenging systems that have been studied by ¹⁴N NMR techniques include di- or tripeptides such as AspAla, AlaAlaGly or GlyGlyGly^{9–11}. Other structural investigations have explored the different modes of self-assembly of nucleosides such as guanosine^{12–14}. Applications have also been reported on pharmaceutical polymorphs, cocrystals and amorphous dispersions^{15–18}. However, the acquisition of ¹⁴N spectra of large biomolecules such as the GB3 protein (6 kDa) has required the use of dynamic nuclear polarization (DNP) techniques¹⁹. We have recently introduced a double cross-polarization (DCP) transfer ¹H → ¹⁴N → ¹H for the indirect detection of ¹⁴N spectra²⁰. We have shown the method to be straightforward to implement and robust with respect to sample-specific parameters, and efficient for C_Q values up to about 3 MHz.

The calculation of NMR parameters based on density functional theory (DFT)²¹ has emerged as an invaluable tool for spectral interpretation and assignment in both solid-state and solution NMR^{22–28}. In particular, the combination of DFT methods and solid-state NMR techniques has become known as NMR crystallography^{29–34}. The possibility of computing both isotropic and anisotropic interactions expected for a molecular model can be used to assess the agreement between experimental evidence and different structural hypotheses. These computational tools are particularly useful for the prediction of ¹⁴N NMR spectra since the wide frequency dispersion of the resonances and their inhomogeneous lineshapes (so-called “powder patterns”) are determined by both anisotropic chemical shifts and second-order quadrupolar interactions.

In this paper, we present the structural study of a biologically-relevant polypeptide by means of ¹⁴N NMR spectroscopy at a very fast magic-angle spinning (MAS) rate $\nu_R = 100$ kHz^{35–39}. In particular, the non-methylated amide nitrogen sites in a powdered polycrystalline sample of cyclosporin, a well-known immunosuppressing drug^{40,41}, are investigated by means of DCP ¹H-¹⁴N-¹H through-space correlation experiments. This 1.2 kDa cyclic undecapeptide represents one of the largest molecules so far investigated by means of ¹⁴N NMR methods and should open the way for solid-state studies of proteins of several tens of kDa. Through-space dipolar couplings between ¹⁴NH nitrogen atoms and directly-bound ¹H^N protons can be readily identified. Spatial proximities between proton-carrying amide nitrogen atoms and remote backbone or side-chain protons of surrounding residues can also be determined by tuning the length of the polarization-transfer interval. The full network of heteronuclear ¹⁴N-¹H dipolar couplings can thus be characterized. DFT calculations of chemical shifts and quadrupolar couplings support the experimentally observed spectral features. Precious indications about the mean frequencies and inhomogeneous lineshapes of ¹⁴N resonances for different environments can be drawn from computations of NMR parameters. In particular, for a given ¹⁴N environment, the mean frequencies are predominantly determined by the magnitude of the quadrupolar constant C_Q , via the isotropic part of the second-order quadrupolar interaction, while isotropic chemical shifts are almost negligible.

Results

Solid-state NMR. The structure of cyclosporin reported in the literature⁴², initially obtained by X-ray diffraction techniques, was optimized with planewave-pseudopotential DFT methods by means of the VASP code^{43–45}. The resulting geometry is shown in Fig. 1(a). The sequence of residues in this undecapeptide is shown in (b). The four H-bonds (three transannular and one of 3 → 1

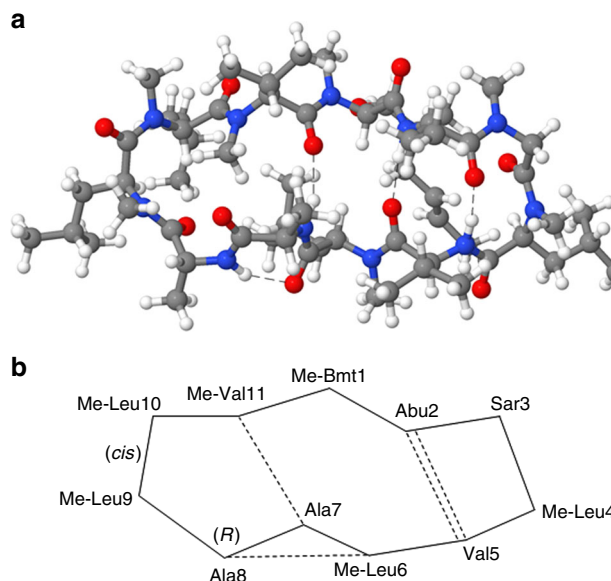


Fig. 1 Structure of cyclosporin and scheme depicting the sequence of residues. **a** DFT-optimized structure of cyclosporin. Nitrogen, oxygen, carbon and proton atoms are shown in blue, red, grey and white, respectively. Hydrogen bonds are represented by dashed lines. **b** Schematic representation of the structure of (a) where the relevant residues are shown. H-bonds are depicted by dashed lines. The two parallel dashed lines indicate an antiparallel β -pleated sheet structure. The D-Ala8 residue has an absolute (*R*) configuration, and there is a *cis*-amide linkage between Me-Leu9 and Me-Leu10

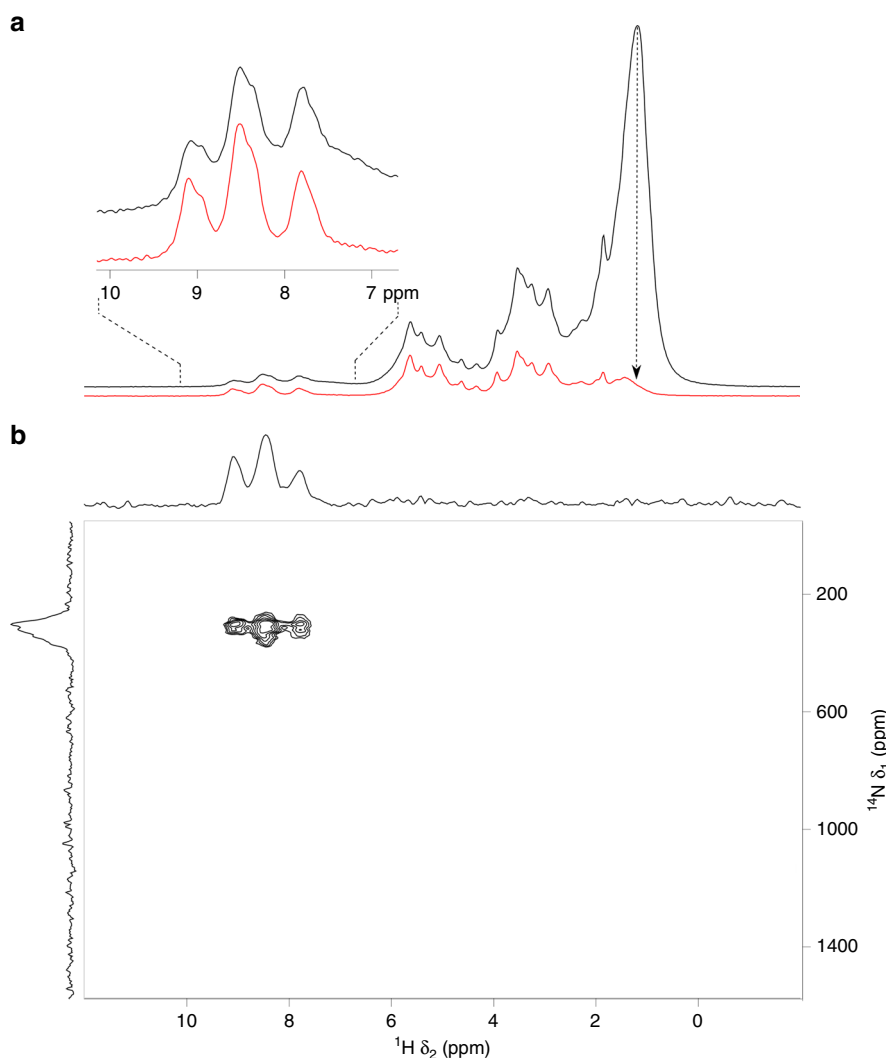


Fig. 2 ^1H 1D and ^1H - ^{14}N 2D NMR spectra of cyclosporin in the solid state. **a** Proton MAS NMR spectrum of cyclosporin acquired with a rotor-synchronized 90° - τ - 180° - τ -echo in a static field $B_0 = 20.0$ T (850 MHz for ^1H , 61.4 MHz for ^{14}N) at a spinning rate of $\nu_R = 100$ kHz (black). In red, proton spectrum obtained with selective presaturation of the methyl region by means of a weak rf -field $\nu_1 = 125$ Hz applied for a duration $\tau_p = 1$ s at the frequency indicated by a dashed vertical arrow. The black and red spectra are displaced vertically for clarity. The four amide protons H^{N} in the 7–10 ppm region are highlighted in the expansion. **b** Two-dimensional $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ double cross-polarization (DCP) spectrum of cyclosporin acquired with selective presaturation and short contact pulses of length $\tau_{\text{CP}} = 18\tau_R = 180$ μs . 400 transients were averaged for each of the 256 rotor-synchronized t_1 increments with a recovery delay of 0.2 s, resulting in a total experimental time of ca. 34 h. No line broadening was applied in **(a)**, whereas 100 and 500 Hz were applied to direct and indirect dimensions of **(b)**, respectively

type) involving the four non-methylated NH sites are indicated by dashed lines. These H-bonds result in a substantially rigid conformation in both crystalline state and in solution⁴². The proton spectrum of a powdered sample of cyclosporin acquired with a 90° - τ - 180° - τ -echo in a field $B_0 = 20.0$ T (ν_0 (^1H) = 850.0 MHz, ν_0 (^{14}N) = 61.4 MHz) and at a spinning rate $\nu_R = 100$ kHz (rotor period $\tau_R = 10$ μs) is shown in Fig. 2(a). High fields and ultra-fast spinning rates improve the resolution of proton spectra and allow the separation of distinct environments by virtue of the increased spread of the chemical shifts and more efficient averaging of homonuclear dipolar couplings. Figure 1(a) also shows, in red and vertically displaced for clarity, a spectrum obtained after initial selective presaturation⁴⁶ of the aliphatic region (dashed black arrow). The expansion shows how presaturation leaves the four amide H^{N} protons that appear in the range 7–10 ppm substantially unperturbed. Shoulders associated with all three peaks towards lower frequencies suggest that the sample may be inhomogeneous and possibly contain a minor

phase or polymorph. Attenuation of the intense signals of aliphatic protons that are not directly bound to any N atoms allows one to acquire the ^1H - ^{14}N correlation spectrum as shown in Fig. 2(b), which is substantially free of any t_1 noise artifacts that appear due to imperfect cancellation of intense signals if one does not use any presaturation. This spectrum was acquired with contact pulses for cross polarization of length $\tau_{\text{CP}} = 18\tau_R = 180$ μs , so that transfers over short distances in non-methylated NH amide groups are favored. The vertical ^{14}N (ω_1) dimension shows the full spectral width of $1/\tau_R = \nu_R = 100$ kHz (1629 ppm for ^{14}N at $B_0 = 20.0$ T). It is worth bearing in mind that the 4 amide H^{N} protons represent a fraction of only $4/111 = 3.6\%$ with respect to the full “bath” of 111 protons in this sample, without considering possible water molecules involved in hydration. This is much lower than in ^{14}N MAS studies previously reported which utilized samples such as glycine, alanine or histidine, where the analogous ratios are $3/5$, $3/7$ and $4/9$ (60, 43 and 44%), respectively. Nonetheless, a clear correlation between the amide

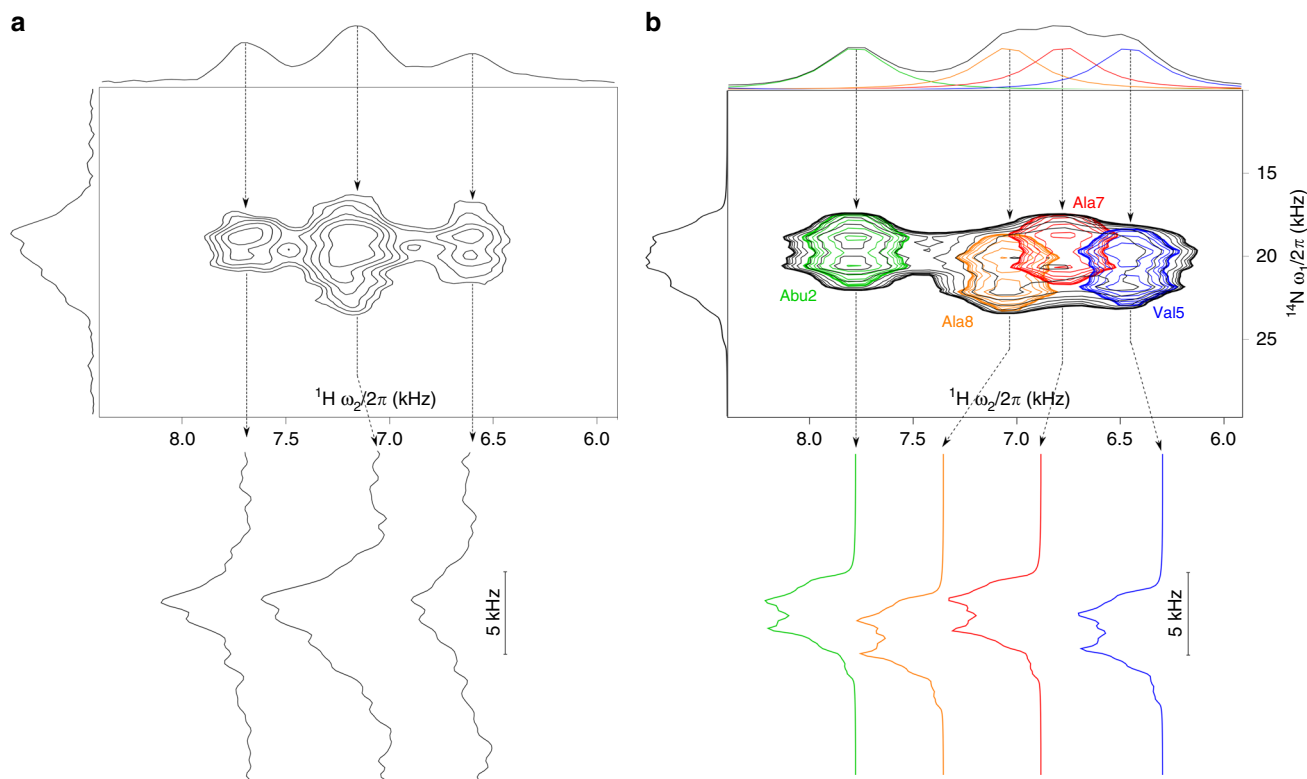


Fig. 3 Comparison between experimental and DFT-calculated ^1H - ^{14}N 2D correlation spectra. **a** Expansion of the amide region of the 2D spectrum of Fig. 2 (b). The arrows indicate the positions in the proton spectrum where the ^{14}N cross sections shown below were extracted. **b** Simulated spectrum analogous to that of (a) resulting from isotropic chemical shifts and quadrupolar parameters calculated by DFT methods. The 2D correlations and extracted ^{14}N second-order lineshapes of the four distinct NH sites are highlighted with different colors. The simulated numbers of points were 32 and 128, for direct and indirect dimensions, respectively. The corresponding line broadenings were 350 and 300 Hz

protons and nitrogen sites can be identified in Fig. 2(b) near ca. 300 ppm in the vertical ^{14}N dimension.

Short-range correlations. Figure 3(a) shows an expansion of the 2D spectrum where three correlation peaks can be easily resolved in the horizontal ^1H dimension, whereas the four NH^{H} signals appear substantially overlapped in the vertical ^{14}N dimension. In order to unravel the overlapping 2D correlations and achieve spectral assignment, DFT calculations of NMR parameters were performed using the CASTEP code⁴⁷. Figure 3(b) shows the resulting simulated 2D spectrum based on proton and nitrogen shifts as well as quadrupolar couplings predicted by CASTEP. The signals from separate sites are highlighted in different colors (Abu2 in green, Ala8 in orange, Ala7 in red, and Val5 in blue.) A few vertical cross sections were extracted from the 2D spectra in (a) and (b) as indicated by arrows. Considering that homogeneous signal decay does not allow subtle inhomogeneous features to be resolved, the simulated linewidths agree well with the experimental cross sections. More specifically, (i) the four NH correlations are indeed expected to overlap in the vertical ^{14}N dimension and (ii) the size of the second-order quadrupolar broadening is correctly estimated to ca. 5 kHz by DFT. The experimental vertical cross sections do not allow one to resolve any second-order features of the ^{14}N lineshapes. This is consistent with our earlier work²⁰ and can be attributed to (i) the non-uniform efficiency of the CP transfer over the inhomogeneous lineshapes, and (ii) homogeneous signal decay in the evolution interval.

Long-range correlations. Figure 4 shows a ^1H - ^{14}N correlation spectrum acquired with longer CP contact pulses of length $\tau_{\text{CP}} =$

$60\tau_{\text{R}} = 600 \mu\text{s}$. When compared to that of Fig. 2(a), this clearly shows long-distance correlations between the NH^{H} sites and remote protons. This ^1H - ^{14}N correlation spectrum, in analogy to routine ^1H - ^{15}N correlation techniques that can be used for $I = 1/2$ nuclei, supplies 2D contour ridges which can map interatomic proximities beyond one-bond connectivities. The molecular structure of Fig. 1, optimized by DFT methods, can be very helpful to assist the assignment of through-space correlations. By inspection of the internuclear distances in this structure, one finds that the ^{14}N nuclei of NH groups have intra-residue dipole-dipole couplings with H^{N} , H^{α} , H^{β} and H^{γ} side-chain protons of ca. 8, 1, 1 and 0.5 kHz, respectively. Other inter-residue couplings ranging from ca. 0.5 down to 0.15 kHz may also be identified. Relevant intra- and inter-residue distances shorter than 3.3 Å are given in Table 1. The assignment of long-range correlations based on the dipolar couplings in Table 1, supported by proton chemical shifts calculated by DFT, is highlighted by dashed magenta rectangles in the experimental spectrum of Fig. 4. The H^{α} protons are clearly distinguished near 5 ppm in the horizontal ^1H dimension. The strongest correlations are observed for the most shielded methyl protons of Val5, which are in proximity of both intra- (Val5) and inter-residue (Ala7) nitrogen sites. The $(\text{H}^{\beta})_3$ groups of Ala7 and Ala8 give rise to strong correlations with intra-residue NH^{H} nitrogen sites. Other clear correlations can be assigned to $(\text{H}^{\beta})_2$ of Abu2 and H^{γ} of Bmt1. The NCH_3 protons of Leu4 can transfer polarization to the NH^{H} group of Val 5, the NCH_3 protons of Val11 are in contact with the NH^{H} sites of both Ala7 and Ala8, and the NCH_3 protons of Sar3 are in spatial proximity with the NH^{H} site of Abu2. Also in agreement with the internuclear distances extracted from the structure optimized by DFT, one does not observe any correlations between the four NH^{H} sites and the side

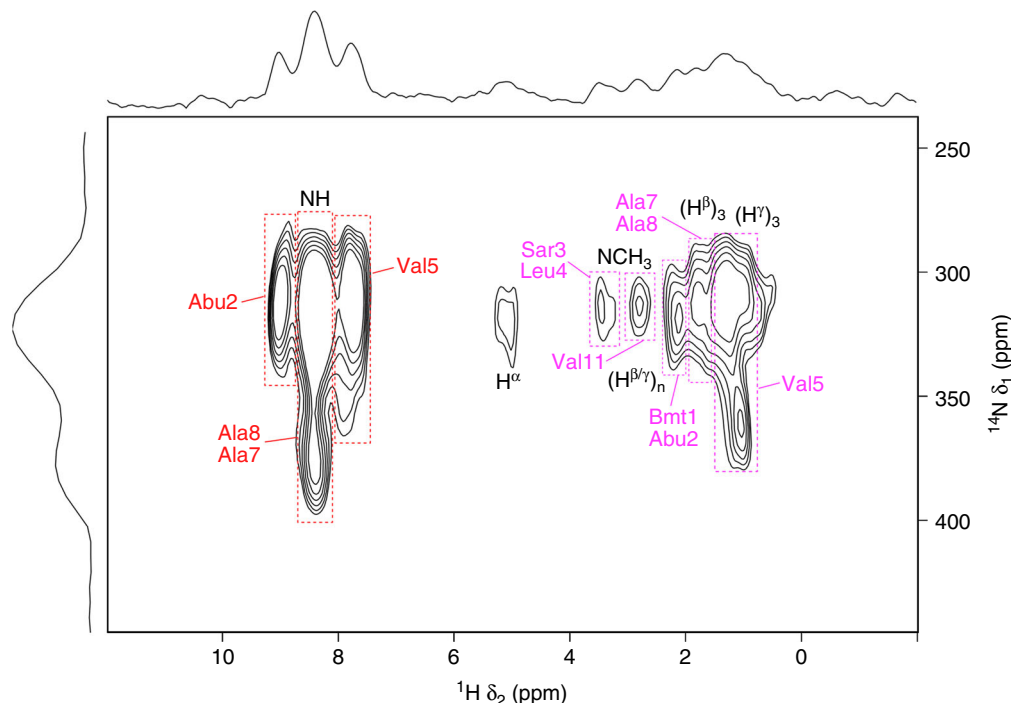


Fig. 4 Long-range correlation ^1H - ^{14}N 2D NMR spectrum. Two-dimensional $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ DCP spectrum analogous to that of Fig. 2(b) but acquired with longer contact pulses of length $\tau_{\text{CP}} = 60\tau_{\text{R}} = 600 \mu\text{s}$. One-bond distance correlations for the four NH sites are assigned according to the DFT calculations and highlighted in red. Long-range correlations between ^{14}N and nearby intra- and inter-residue protons are highlighted in magenta. These were assigned according to the dipolar contacts in Table 1. In all, 420 transients were averaged for each of the 64 rotor-synchronized t_1 increments with a recovery delay of 0.2 s, resulting in a total experimental time of ca. 9 h

Table 1 Internuclear distances calculated by DFT methods

				Bmt1	Bmt1	Sar3		
Abu2	H^{N}	H^{α}	$(\text{H}^{\beta})_2$	H^{α}	H^{γ}	NCH_3		
	1.03	2.06	2.63	2.51	2.90	2.91		
Val5	H^{N}	H^{α}	$(\text{H}^{\gamma})_3$	H^{β}	H^{α}	H^{ϵ}	NCH_3	
	1.02	2.05	2.66	2.75	2.96	3.29	3.17	
				Leu6	Val5	Ala8	Val11	
Ala7	H^{N}	H^{α}	$(\text{H}^{\beta})_3$	H^{α}	$(\text{H}^{\gamma})_3$	H^{N}	NCH_3	
	1.02	2.07	2.70	2.50	2.90	2.69	3.06	
Ala8	H^{N}	H^{α}	$(\text{H}^{\beta})_3$	H^{α}	NCH_3			
	1.03	2.06	2.68	2.65	3.23			

^{14}N - ^1H distances up to 3.3 Å between the four non-methylated N^{H} nitrogen atoms and various (intra- and inter-residue) protons in cyclosporin as extracted from the DFT-optimized structure. Amino acids involved in inter-residue contacts are listed in bold

chains of Leu4, Leu9 and Leu10. No evidence could be found for any methylated NCH_3 sites in the vertical ^{14}N dimension. We ascribe this finding to poor polarization transfer, possibly due to inefficient spin locking for ^{14}N sites with C_Q values that should be in excess of ca. 4 MHz according to DFT. The full set of relevant dipolar proton-nitrogen couplings utilized for the assignment of the spectrum of Fig. 4 is summarized in Table 1.

The atomic details of the four N^{H} sites are shown in Fig. 5. The correlation between the experimental ^1H shifts of Fig. 4, assigned on the basis of internuclear ^1H - ^{14}N distances of Table 1, and those calculated by CASTEP, is shown in Fig. 6. A correlation coefficient $R^2 = 0.995$ is obtained, with a root mean square error of 0.3 ppm, thus further supporting our assignment and the good agreement between our experimental and calculated data.

Quadrupolar parameters. Figure 7 presents simulated ^{14}N spectra based on the chemical shifts and quadrupolar parameters calculated by DFT. The NH and NCH_3 sites are shown in blue and red, respectively. Analogous spectra where only the isotropic shift interactions were taken into account appear on the right-hand side, displayed by dashed lines with the same color coding. The anisotropic part of the chemical shift was not included since it is fully averaged out by fast spinning ($\nu_{\text{R}} = 100 \text{ kHz}$). These simulations allow one to appreciate how the second-order quadrupolar interaction does not only produce inhomogeneous broadening, but also results in significant isotropic shifts of the centers of gravity (first moments) of the powder patterns. Importantly, the DFT calculations show that the NH and NCH_3 amide environments tend to fall in two distinct spectral regions. The NMR parameters for all 11 ^{14}N nuclei and for the relevant ^1H nuclei of the 4 NH sites are given in Table 2. All non-methylated NH sites feature very similar electric field gradients (EFGs), with a narrow dispersion of quadrupolar coupling constants $C_Q = -3.64 \pm 0.12 \text{ MHz}$. Similarly, the NCH_3 sites are associated with narrowly distributed quadrupolar couplings $C_Q = -4.27 \pm 0.09 \text{ MHz}$. Furthermore, the asymmetry parameters in the two environments are also distinct, with average values $\eta_Q = 0.38 \pm 0.04$ for NH sites and $\eta_Q = 0.13 \pm 0.02$ for NCH_3 sites.

The isotropic part of the second-order quadrupolar interaction is proportional to C_Q^2/ν_0 and results in shifts of the first moments of the inhomogeneous lineshapes that, for low- γ nuclei with large C_Q values such as ^{14}N , can significantly exceed the range of the isotropic chemical shifts. For the two single-quantum transitions $m_I = \pm 1 \leftrightarrow m_I = 0$ of a $I = 1$ spin, the isotropic quadrupolar shift is:⁴⁸

$$\nu_{\text{iso}}^{(Q)} = \frac{3}{40} \frac{C_Q^2}{\nu_0} \left(1 + \frac{1}{3} \eta_Q^2 \right). \quad (1)$$

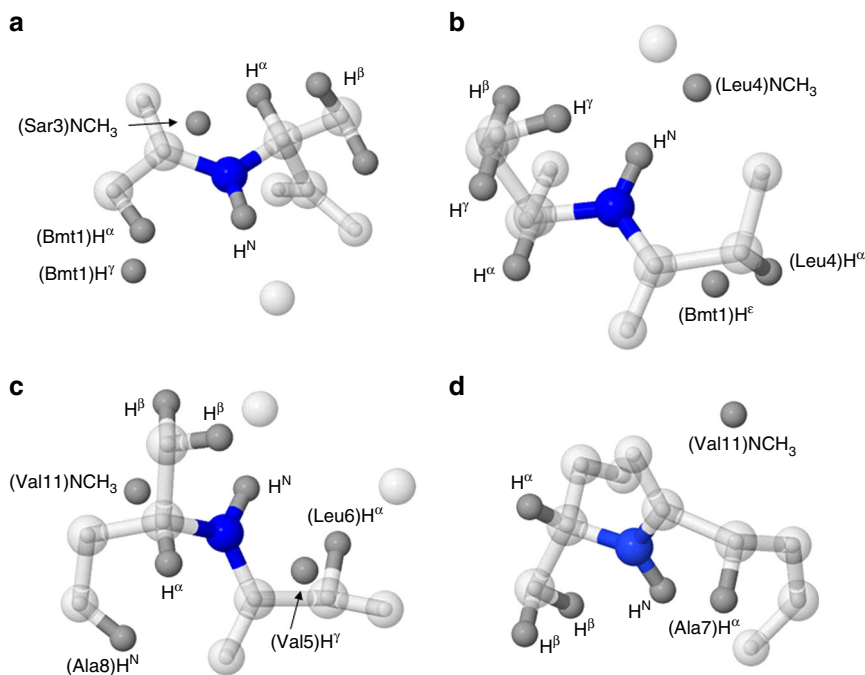


Fig. 5 ^1H - ^{14}N dipolar networks of non-methylated amide groups in cyclosporin. **a** Inter- and intra-residue proton nuclei within 3.3 Å of the NH nitrogen-14 site of Abu2 as extracted from the DFT optimized structure of cyclosporin. **b-d** Analogous environments for Val5 (**b**), Ala7 (**c**) and Ala8 (**d**) residues, respectively. Nitrogen and proton atoms are shown in blue and dark grey, respectively, whereas all other atoms are in transparent light grey for clarity

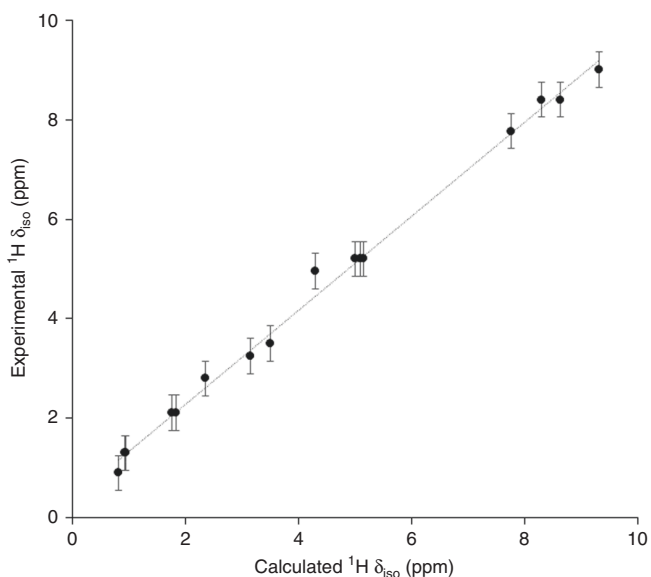


Fig. 6 Correlation between experimental and DFT-calculated ^1H isotropic chemical shifts. The experimental ^1H shifts of the 2D spectrum of Fig. 4, assigned on the basis of the ^1H - ^{14}N internuclear distances of Table 1, are plotted versus those calculated by CASTEP. A correlation coefficient $R^2 = 0.995$ is obtained, with root mean square error of 0.3 ppm. The error bars of the experimental data are set to $2 \times \text{lb} = 600$ Hz. Numerical values are given in Supplementary Table 1

For the average quadrupolar parameters calculated by DFT for cyclosporin in Table 2, Eq. (1) yields $\nu_{iso}^{(Q)} = 16.9$ and 22.4 kHz at 20.0 T for the N^{H} and NCH_3 environments, respectively. In other words, the isotropic quadrupolar interaction leads to a relative displacement of the ^{14}N signals of the N^{H} and NCH_3 sites by 5.5 kHz. In contrast, if one considers the isotropic part of the chemical shift interaction (which determines the dispersion of

^{15}N signals in solids and solutions), one finds that the average isotropic shifts of the NH and NCH_3 sites in cyclosporin are $\delta_{iso} = 6.6$ and 6.2 kHz, i.e., a difference of only ca. 0.4 kHz. The ranges of the average isotropic displacements for ^{14}N sites in cyclosporin are indicated by horizontal continuous and dashed bars in Fig. 7 for the shifts $\nu_{iso}^{(Q)}$ and δ_{iso} , respectively, with the same color coding. A continuous black bar indicates the overall difference in displacement between the average NH and NCH_3 sites due to both shift and quadrupolar interactions. These considerations imply that, since similar ^{14}N environments lead to narrowly distributed C_Q values, one can expect extensive overlap of lineshapes for chemically similar N environments (say, all NH sites), but a good separation between different chemical environments (say, between NH and NCH_3 sites) by virtue of their different quadrupolar interactions. This observation is promising for future applications, as the resolution required to distinguish similar chemical environments can be supplied, as in the case at hand, in the horizontal ^1H dimension of the 2D spectrum, where the resolution is expected to improve further at higher fields and faster spinning rates. Furthermore, these results suggest that the general lower accuracy of DFT-calculated isotropic shifts of ^{15}N atoms when compared to that of ^1H and ^{13}C species^{49,50}, may be reconsidered when dealing with ^{14}N isotopes in the solid state, as both isotropic displacement $\nu_{iso}^{(Q)}$ and anisotropic linewidth due to second-order quadrupolar interactions may largely exceed differences in isotropic chemical shifts.

Experimental ^{14}N linewidths. The 2D spectrum shown in Figs. 2 (b), 3(a) results from the acquisition of 256 rotor-synchronized t_1 time increments. The indirect acquisition time is therefore $t_1^{\text{max}} = 256 \times 10 \mu\text{s} = 2.56$ ms, which allows for an indirect digital resolution of ca. 400 Hz. However, the inhomogeneous ^{14}N lineshapes experimentally observed are on the order of 5 kHz, in full agreement with our DFT calculations. Considering that second-order features are difficult to identify due to non-uniform cross-polarization transfer for different crystallite

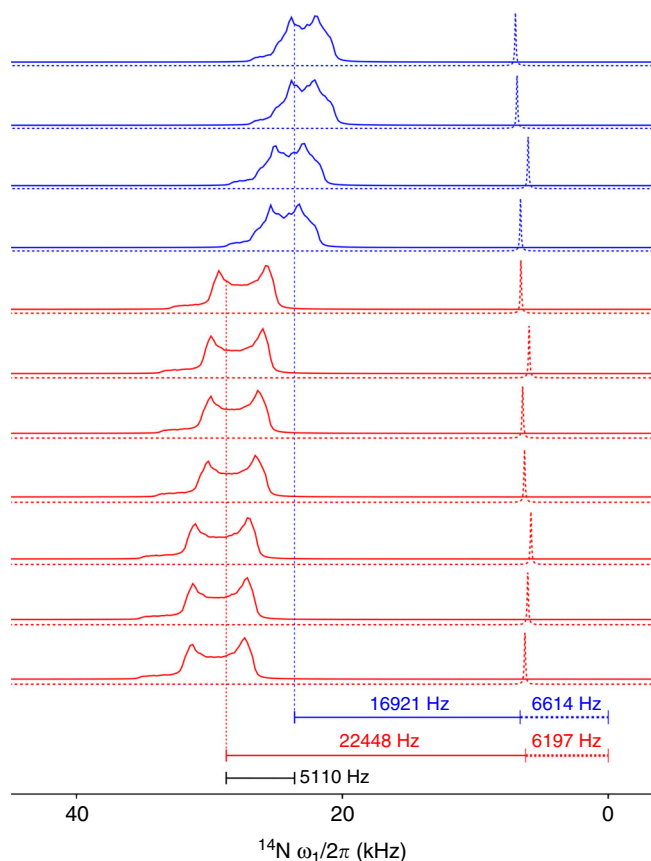


Fig. 7 Simulations of anisotropic ^{14}N lineshapes. One-dimensional ^{14}N NMR spectra of cyclosporin based on the chemical shift and quadrupolar parameters calculated by DFT. NH and NCH_3 sites are shown in blue and red, respectively. Analogous spectra which consider only the isotropic shifts for these latter spectra was reduced by a factor 20 to compensate for the narrow linewidths. The line broadening for anisotropic and isotropic simulations were 300 and 100 Hz, respectively. Horizontal bars in the lower part of the figure indicate the ranges of isotropic displacements due to the quadrupolar and chemical shift interactions, in continuous and dashed lines, respectively, assuming the averaged NMR parameters in Table 2. The overall difference in displacement between the average NH and NCH_3 sites due to both shift and quadrupolar interactions is shown by a black bar

orientations, homogeneous losses and the extensive overlap expected for similar chemical environments, one may conclude that the experimental times for the acquisition of spectra such as that of Fig. 2(b) may be substantially reduced without any loss of structural information. In order to prove this concept, the indirect ^{14}N projection of the spectrum of Fig. 2(b) is compared in Supplementary Figure 1 to that resulting from Fourier transformation of only 160 and 128 t_1 time increments. One can appreciate that the signal intensity is conserved, while the reduction of digital resolution has little effect on the observed projections. Moreover, the identical resolution in the direct ^1H dimension grants for substantially identical 2D correlation contours. Further experimental proof of this concept is given by the spectrum of Fig. 4, where only 64 t_1 time increments were acquired.

Discussion

This paper presents the structural study of amide functions in polypeptides like cyclosporin by means of solid-state ^{14}N NMR spectroscopy. Double cross polarization (DCP) $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$

Table 2 NMR parameters of all eleven ^{14}N sites of cyclosporin calculated with CASTEP

	Residue	^{14}N			$\nu_{\text{iso}}(^{\text{Q}})$ (ppm)	^1H δ_{iso} (ppm)
		δ_{iso} (ppm)	C_{Q} (MHz)	η_{Q}		
NH	Abu2	111.7	-3.52	0.44	261.9	8.5
	Val5	98.1	-3.76	0.36	292.5	7.1
	Ala7	113.7	-3.54	0.36	259.7	7.4
	Ala8	107.4	-3.73	0.37	288.8	7.7
	Average	107.7	-3.64	0.38	275.7	7.7
	σ	6.9	0.12	0.04	17.3	0.6
NCH_3	Bmt1	104.8	-4.21	0.15	355.5	
	Sar3	102.5	-4.24	0.16	361.1	
	Leu4	94.7	-4.37	0.12	381.4	
	Leu6	107.0	-4.14	0.13	342.3	
	Leu9	96.9	-4.25	0.11	360.6	
	Leu10	98.5	-4.36	0.12	379.6	
	Val11	102.1	-4.35	0.13	378.6	
	Average	100.9	-4.27	0.13	365.6	
	σ	4.4	0.09	0.02	14.7	

The magnetic shieldings were arbitrarily referenced to 200 ppm for ^{14}N and 30 ppm for ^1H nuclei, with $\delta_{\text{iso}}^{(i)} = \sigma_{\text{ref}} - \sigma_{\text{iso}}^{(i)}$ for the i -th nucleus. The standard deviation σ provides a measure of the width of the distributions of all parameters

offers sufficient sensitivity to detect all four NH environments in a sample where protons bound to nitrogen atoms constitute less than 4% of the full proton bath. The use of long contact pulses allows one to monitor spatial proximities beyond immediate neighbors. Therefore, structural details similar to those derived from conventional solution-state NMR techniques on ^{15}N -enriched samples can be obtained in natural ^{14}N abundance in the solid state. Spurious signals rising from the large proton bath can be efficiently suppressed by weak selective presaturation of the aliphatic region. Depending on the desired resolution in the ^{14}N dimension, satisfactory spectra can be obtained in about 10 h when spinning at 100 kHz at 20 T. The structural details unveiled by these experiments are in agreement with interatomic distances obtained with DFT methods. The full set of dipolar ^{14}N - ^1H couplings involving amide nitrogen nuclei can be ascertained. Calculations of chemical shifts and quadrupolar couplings agree very well with the experimental evidence. Some overlap is expected and indeed observed for the four NH environments. The seven methylated NCH_3 sites are also predicted to overlap in part. The signals of the NH and NCH_3 sites, however, do not overlap because of distinct quadrupolar interactions that result in different second-order isotropic shifts. The use of very high fields above 23 T (above 1 GHz for ^1H , i.e., above 72 MHz for ^{14}N) would lead to a significant reduction of second-order quadrupole interactions, therefore lead to narrower signals, and possibly allow indirect detection of methylated NCH_3 sites by our DCP method. Furthermore, fast spinning beyond 100 kHz would further reduce broadening by proton-proton couplings and lead to line-narrowing in the ^1H dimension of the 2D correlation spectra. This study should open the way for solid-state ^{14}N NMR studies of proteins, nucleic acids and other biomolecules in natural abundance.

Methods

Solid-state NMR. The sample of cyclosporin A was purchased from Sigma-Aldrich and used without further purification. All DCP MAS spectra were recorded at the CEMHTI CNRS laboratory in Orleans, on a wide-bore Bruker 850 spectrometer (20.0 T) with an Avance-III console, using 0.7 mm ZrO_2 rotors containing ca. 1 mg cyclosporin spinning at $\nu_{\text{R}} = 100$ kHz in a double-resonance MAS probe. The proton rf -field strength was $\nu_1 = 179$ kHz (90° pulse width $\tau_p = 1.4$ μs .) The constant rf -field strengths during both $^1\text{H} \rightarrow ^{14}\text{N} \rightarrow ^1\text{H}$ cross-polarization intervals were $\nu_1(^1\text{H}) = 53$ and $\nu_1(^{14}\text{N}) = 83$ kHz, whereas a weak rf field $\nu_1(^1\text{H}) = 125$

Hz was utilized for selective presaturation during an interval $\tau_p = 1$ s. The lengths of the contact pulses were $\tau_{cp} = 180$ and $600 \mu\text{s}$, for the spectra of Figs. 2 and 4, respectively. Recovery delays of 0.2 and 1 s were used for experiments with and without presaturation, respectively. Phase cycling for coherence pathway selection was used on the ^{14}N channel only, with phases $\phi_1 = \{0\ 2\}$, $\phi_2 = \{0\ 0\ 2\ 2\}$ for the first and second contact pulses in the ^{14}N channel, respectively, with a receiver phase $\phi_{\text{rec}} = \{0\ 2\ 2\ 0\}$. The chemical shifts were referenced with respect to adamantane for ^1H and NH_4Cl for ^{14}N . The magic angle was adjusted using KBr as usual.

DFT calculations. All DFT calculations were performed at the gamma point using the PBE⁵¹ functional and a planewave energy cutoff of 50 Ry. The geometry optimization was performed with VASP 5.4.1^{43–45} on a single molecule in a $25 \times 25 \times 25 \text{ \AA}$ cubic cell, requiring forces acting on all atoms to be smaller than 0.02 eV/\AA . The electronic cores were described with the projector augmented wave method⁵². The resulting structure is superposed onto that obtained by X-ray diffraction in Supplementary Figure 2. NMR parameters were computed for the optimized geometries using CASTEP⁴⁷ as implemented in the Materials Studio suite, version 8.0, available on the national supercomputing facility CINES. The gauge-including projector augmented wave (GIPAW) algorithm was employed to reconstruct all-electron wave functions in the presence of a magnetic field^{53,54}. Cartesian coordinates of the optimized structure of cyclosporin utilized in this study and all calculated NMR parameters are given in Supplementary Tables 2 and 3.

NMR simulations. The simulations of the NMR spectra based on DFT parameters were carried out with Simpson⁵⁵ and involved rotor-synchronized detection of the FID assuming an initial state comprising only in-phase coherence $\rho(0) = I_x$. Since the choice of reference shieldings is arbitrary, the offset of the carrier frequency ν_{rf} in these simulations was adapted to match the experimental results. The simulated 2D spectrum of Fig. 3(b) represents pure adsorptive ideal correlations between ideal ^{14}N and ^1H lineshapes, and results from the sum of four individual 2D spectra associated with each of the four HN pairs. Each of these individual 2D spectra was constructed by forming the product $\text{Re}[S(\omega_1)] \text{Re}[S(\omega_2)]$, where the column vector $S(\omega_1)$ is the spectrum for an ideal ^{14}N second-order quadrupolar lineshape simulated with Simpson assuming NMR parameters calculated with CASTEP, and the row vector $S(\omega_2)$ is a Lorentzian lineshape simulated with Simpson and associated with a ^1H spin with an isotropic shift given by CASTEP calculations.

Data availability

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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Author contributions

D.C. designed the study, conducted the experiments, analyzed the data and wrote the manuscript, B.G. performed DFT calculations, G.B. wrote the manuscript.

Additional information

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