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# Cluster Assembled Silicon-Lithium Nanostructures: A Nanowire Confined Inside a Carbon Nanotube

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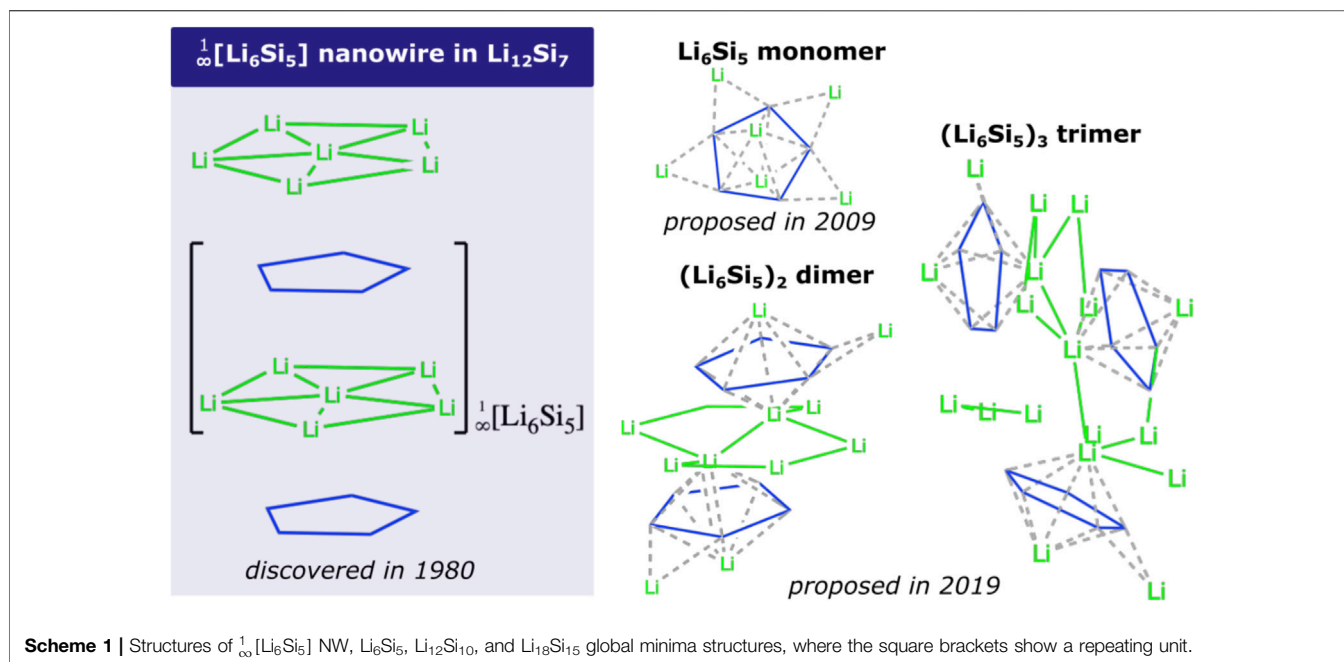
We computationally explore an alternative to stabilize one-dimensional (1D) silicon-lithium nanowires (NWs). The  $\text{Li}_{12}\text{Si}_9$  Zintl phase exhibits the NW  ${}_{\infty}[\text{Li}_6\text{Si}_5]$ , combined with Y-shaped  $\text{Si}_4$  structures. Interestingly, this NW could be assembled from the stacking of the  $\text{Li}_6\text{Si}_5$  aromatic cluster. The  ${}_{\infty}[\text{Li}_6\text{Si}_5]@\text{CNT}$  nanocomposite has been investigated with density functional theory (DFT), including molecular dynamics simulations and electronic structure calculations. We found that van der Waals interaction between Li's and CNT's walls is relevant for stabilizing this hybrid nanocomposite. This work suggests that nanostructured confinement (within CNTs) may be an alternative to stabilize this free NW, cleaning its properties regarding  $\text{Li}_{12}\text{Si}_9$  solid phase, i.e., metallic character, concerning the perturbation provided by their environment in the  $\text{Li}_{12}\text{Si}_7$  compound.

**Keywords:** nanowire, density functional theory, silicon-lithium clusters, carbon nanotube, metallic character

## INTRODUCTION

The insertion of inorganic materials into single-walled carbon nanotubes (SWCNTs), hereinafter identified simply as CNT, enables the encapsulation of extreme nanowires (NWs) with diameters comparable to a unit cell of the parent material (Green, 1998; Sloan et al., 2002; Spencer et al., 2014). Although NWs of similar diameter can be produced using several templates, such as zeolites (Derouane, 1998), mesoporous phases (Alba-Simionescu et al., 2006; Ke et al., 2009), and metal-organic framework (MOF) (Lu et al., 2012) type materials, CNTs present many advantages as templates; they are atomically smooth, electron transparent, readily available, and can be filled by bulk infiltration to create milligram quantities of encapsulated nanowires, at least on a laboratory scale. Thus, encapsulated NW-CNT are scientifically interesting not only on their own but also as precursors to a wide range of other extreme nanowire materials.

In 2016, Ivanov et al. published a theoretical prediction of helix-shaped lithium-phosphorus nanowires encapsulated into single-walled carbon nanotubes ( $\text{LiP}@\text{CNTs}$ ) (Ivanov et al., 2016). Note that helix-shaped  $\text{Li}_n\text{P}_n$  clusters ( $n = 5-9$ ) had previously been reported as global minimum structures (Ivanov et al., 2012). Some solid phases consist of structural motifs like atomic clusters, i.e., in Zintl phases. This connection brings consistency to the use of models based on



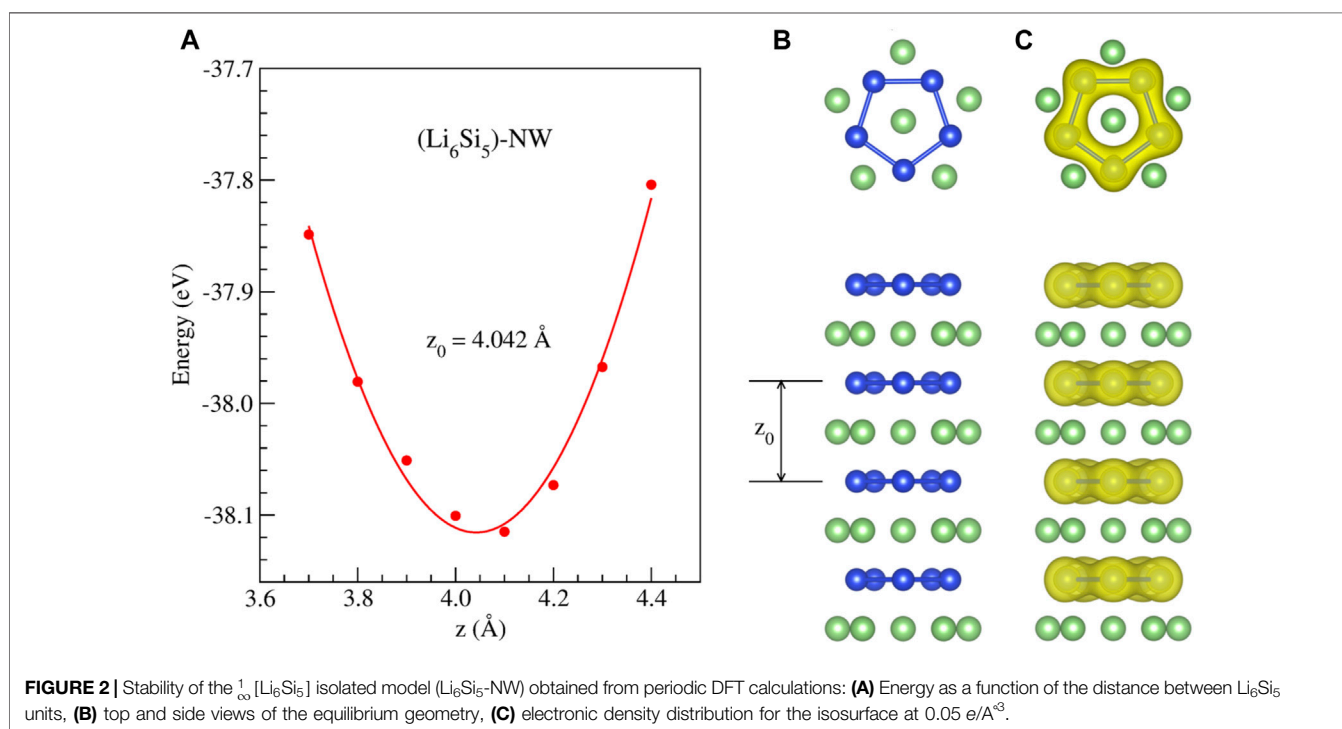
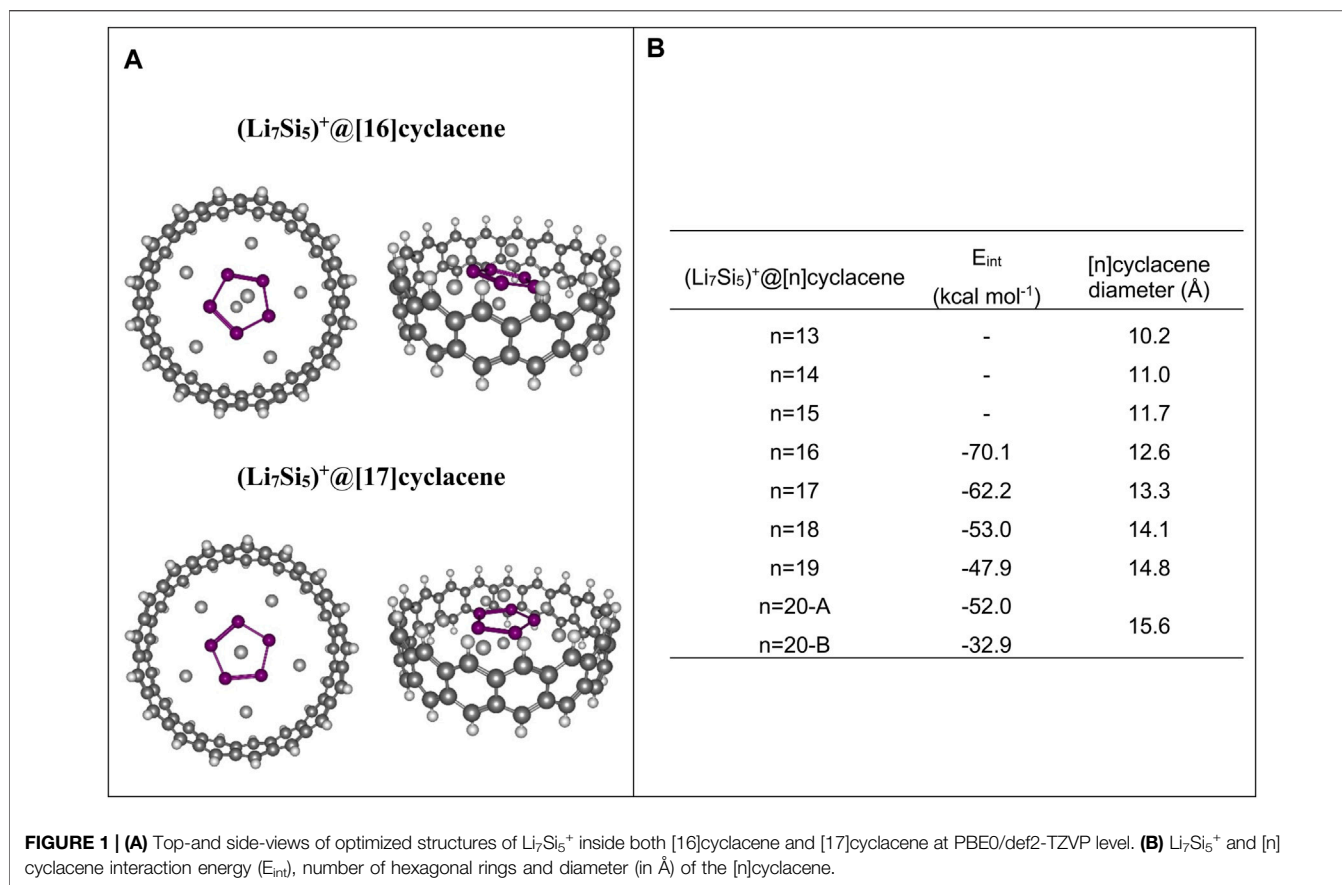
stable clusters to generate NWs inside nanotubes, as proposed in Ivanov's work (Ivanov et al., 2012). The study of these clusters inside CNTs can provide relevant information about these hybrid materials, for example, about their viability (stability analysis), their structural characteristics (geometry analysis), their physical and chemical properties (analysis of their electronic structure).

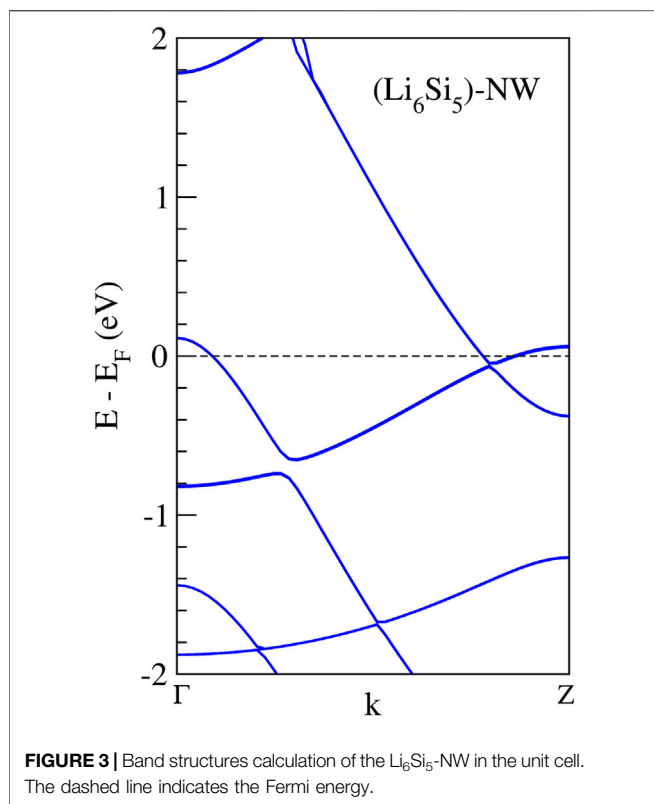
Due to its excellent energy storage capacity, Si has been extensively studied experimentally as a negative electrode material for Li-ion batteries (Gao et al., 2001; Ryu et al., 2004; Li et al., 2006, Li et al., 2008, Li et al., 2009; Obrovac et al., 2007; Song et al., 2014; Shin et al., 2020). Hence, Si lithifies at high temperature (415°C) in a LiCl-KCl melt, identifying potential plateaus evidencing the crystalline phases  $\text{Li}_{12}\text{Si}_7$ ,  $\text{Li}_7\text{Si}_3$ ,  $\text{Li}_{13}\text{Si}_4$ , and  $\text{Li}_{22}\text{Si}_5$  (Wen and Huggins, 1981). In particular, the binary (non-paramagnetic) Zintl-type  $\text{Li}_{12}\text{Si}_7$  silicide contains semi-infinite sandwich-like  ${}^1_{\infty}[\text{Li}_6\text{Si}_5]$  linear chains, consisting of  $\text{Si}_5$  pentagons intercalated with Li atoms (see **Scheme 1**). Note that the unit cell of the Zintl  $\text{Li}_{12}\text{Si}_7$  phase has been rationalized (Nesper, 1990; Chevrier et al., 2010; Köster et al., 2011; Kuhn et al., 2011a, Kuhn et al., 2011b) as  $(\text{Li}_6^{6+}[\text{Si}_5]_6)_2$  ( $\text{Li}_{12}^{10+}[\text{Si}_4]_{10}$ )<sub>2</sub>, with two well-defined silicon moieties: planar  $\text{Si}_5$  rings and the Y-shaped  $\text{Si}_4$  moiety. Such a structural pattern is justified by assigning 26 electrons (20 from 6Si + 6 from 6Li) to the  $\text{Si}_5^{6-}$  ring, favoring Hückel's aromaticity (Hückel, 1930, Hückel, 1931a, Hückel, 1931b; Zhao et al., 2017). This aromatic character is supported by experimental evidence of an upfield shift (to -17.2 ppm) of Li (at the center of the  $\text{Li}_6$  fragment in **Scheme 1**) in the corresponding magic angle NMR (MAS) spectrum (Kuhn et al., 2011b; Köster et al., 2011). It is noteworthy that the  $\text{Si}_5^{6-}$  structural motif is also present in the ternary compound  $\text{Li}_8\text{MgSi}_6$  (Nesper et al., 1986a). On the other hand, at the cluster level, our group has identified that the global minimum (GM) of the  $\text{Li}_6\text{Si}_5$  cluster, consists of an aromatic  $\text{Si}_5^{6-}$  pentagon surrounded by 6  $\text{Li}^+$  counterions (Tiznado et al., 2009; Perez-Peralta et al., 2011; Contreras et al., 2013; Vásquez-Espinal et al., 2018). More

recently, we have identified the GM structures of the oligomers  $(\text{Li}_6\text{Si}_5)_2$  and  $(\text{Li}_6\text{Si}_5)_3$  (Yañez et al., 2019b; Manrique-de-la-Cuba et al., 2021), which also consist of aromatic  $\text{Si}_5$  rings surrounded by Li's (see **Scheme 1**). However, the stacking of  $\text{Li}_6\text{Si}_5$  units does not tend to form the nanowire identified in  $\text{Li}_{12}\text{Si}_7$ , suggesting that  $\text{Li}_{12}^{10+}[\text{Si}_4]_{10}^{10-}$  component (with the Y-shaped  $\text{Si}_4$  moiety) contributes decisively to the stabilization of this NW. In mentioned cluster studies, explorations of the potential energy surface have been performed by hybrid methods, including genetic algorithms (Yañez et al., 2019a, Yañez et al., 2020).

To build new class of materials with desirable properties, using atomic clusters instead of atoms as building blocks, is a remarkable possibility. However, it requires that atomic clusters must retain their identity when assembled, as Khanna and Jena first outlined when they coined the word "cluster-assembled materials" (CAMs) (Khanna and Jena, 1992). These authors argued that the clusters' coupling would have a unique effect on both the assembled material's electronic structure and mechanical properties, which is not possible when the assembly blocks are atoms (Khanna and Jena, 1995; Jena et al., 1996; Jena and Khanna, 1996; Claridge et al., 2009; Jena and Sun, 2018). For a more detailed and timely overview of advances in the assembly of materials from clusters, please refer to the following reviews: (Chakraborty and Pradeep, 2017; Jena and Sun, 2018; Pinkard et al., 2018; Zheng et al., 2019; Doud et al., 2020).

Given the above background, here we evaluated, *in silico*, the stability of the isolated  ${}^1_{\infty}[\text{Li}_6\text{Si}_5]$  NW, as well as its electronic properties. In addition, we studied the hybrid material consisting of the NW confined in a CNT. The latter focused on identifying alternative ways to stabilize this conformation and to evaluate the effect of this association on NW electronic properties. Our density functional theory (DFT) calculations demonstrate that Li-Si@CNTs hybrid systems have excellent stability and thus have potential for experimental realization.





## COMPUTATIONAL METHODS

In the finite models (clusters), geometry optimizations and frequency calculations were performed at the PBE0 (Adamo and Barone, 1999)/Def2TZVP (Weigend and Ahlrichs, 2005) level with the Gaussian 16 program (M. J. Frisch, G. W. Trucks, H. B. Schlegel et al., 2016).

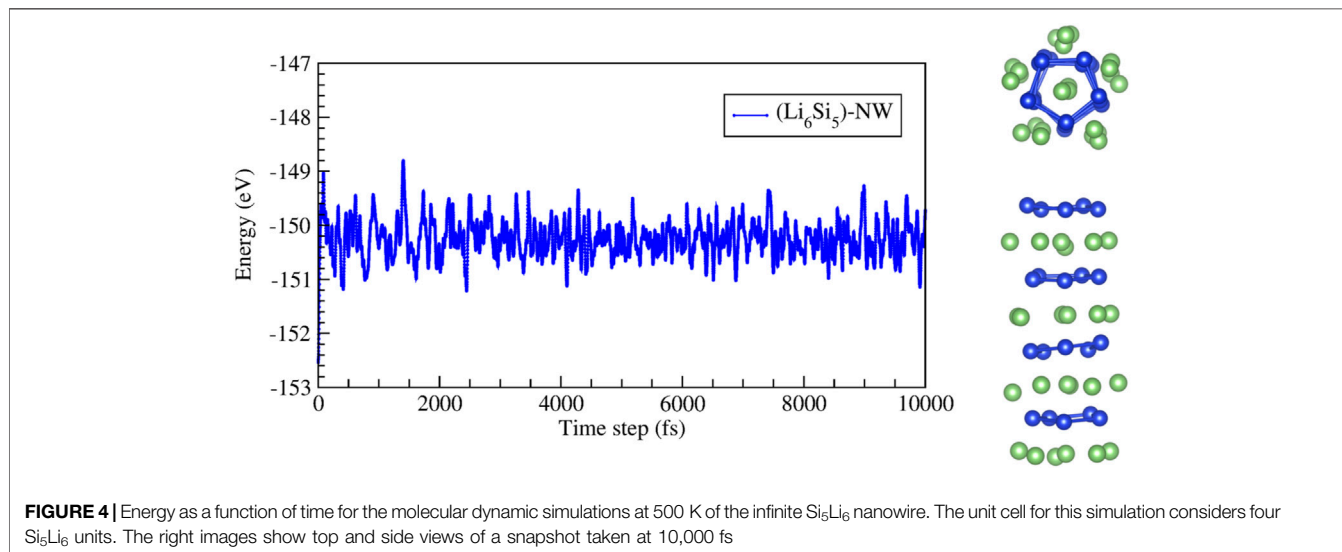
For the solid-state study, we performed first-principles calculations based on DFT (Sham and Kohn, 1966; Kohn

et al., 1996) as implemented in the Vienna Ab Initio Simulation Package (VASP) (Kresse and Furthmüller, 1996). The exchange-correlation energies were calculated at PBE-D3 level (Ernzerhof and Perdew, 1998; Grimme et al., 2010). Plane-wave basis set with a kinetic energy cutoff of 400 eV, and the projector augmented-wave method for the core-valence interaction was employed (Blöchl et al., 1994). The  ${}^1_{\infty}[\text{Li}_6\text{Si}_5]$  NWs were simulated in a large unit cell with volume  $(20 \times 20 \times z_0) \text{ \AA}^3$ , considering periodic boundary conditions along the NW direction, where  $z_0$  is the periodicity. Within this supercell, the lateral distance between NW images is set to 15 Å. We use a  $(1 \times 1 \times 10)$  Monkhorst-Pack k-point mesh (Monkhorst and Pack, 1976). We also study finite  $(\text{Li}_6\text{Si}_5)_4$  structures in the free space, also inside both an armchair and a zigzag single-walled carbon nanotube (SWCNTs). We consider armchair CNTs with chiral indexes (8,8), (9,9), and (10,10), which have diameters of 10.93, 12.27, and 13.63 Å, and zigzag CNTs with chiral indexes (14,0), (15,0), and (16,0) which have diameters of 11.04, 11.80, and 12.59 Å, respectively. For  $(\text{Li}_6\text{Si}_5)_4@$ CNT simulation,  $(22 \times 22 \times z_0) \text{ \AA}^3$  volume was used, where  $z_0$  is the periodicity chosen for the CNTs. All studied structures were allowed to freely relax without any constraint until forces on each atom were smaller than 25 meV/Å. To gain insights on the stability of the NW models in the free space and inside the SWCNTs, we performed Born-Oppenheimer ab initio molecular dynamics (BO-AIMD) simulations within the NVT ensemble at different temperatures, over a total simulation time of 10 ps, considering a time step of 1 fs.

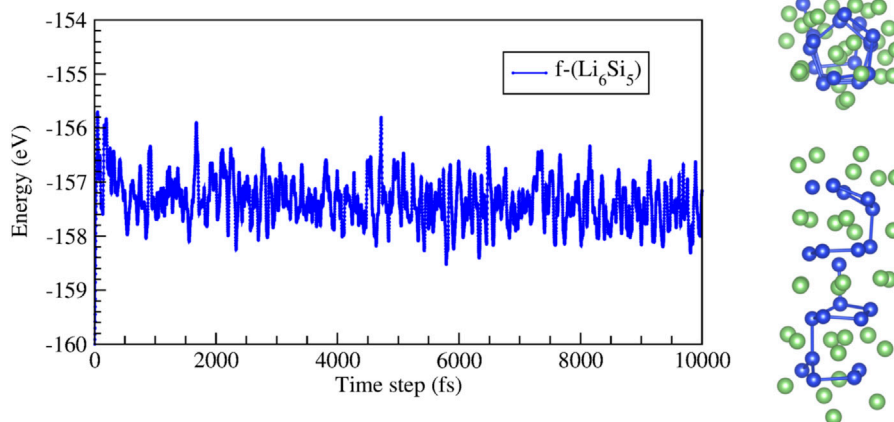
## RESULTS AND DISCUSSION

### Finite Model Tests to Estimate the Optimal Width of SWNTs

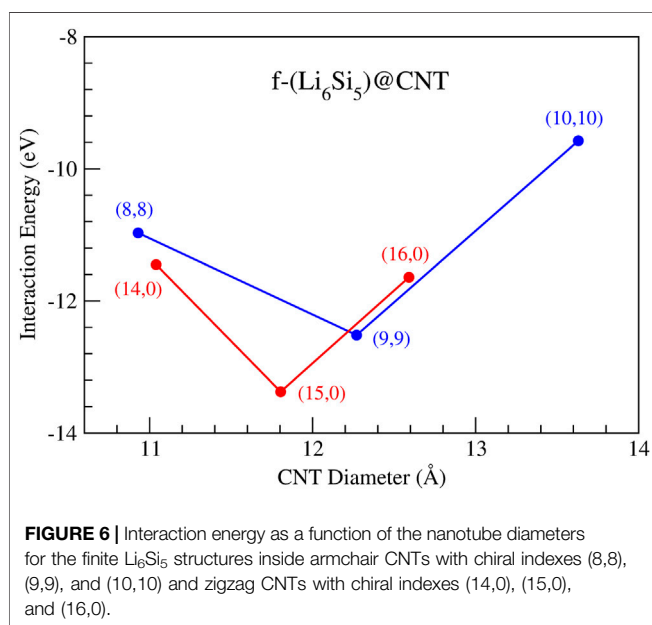
The first question that arises is which is the optimal SWCNT width to favor the  ${}^1_{\infty}[\text{Li}_6\text{Si}_5]$  NW grown? This is a relevant question, considering that the electronic structure of group 1 elements, such as Li, is particularly sensitive to confinement







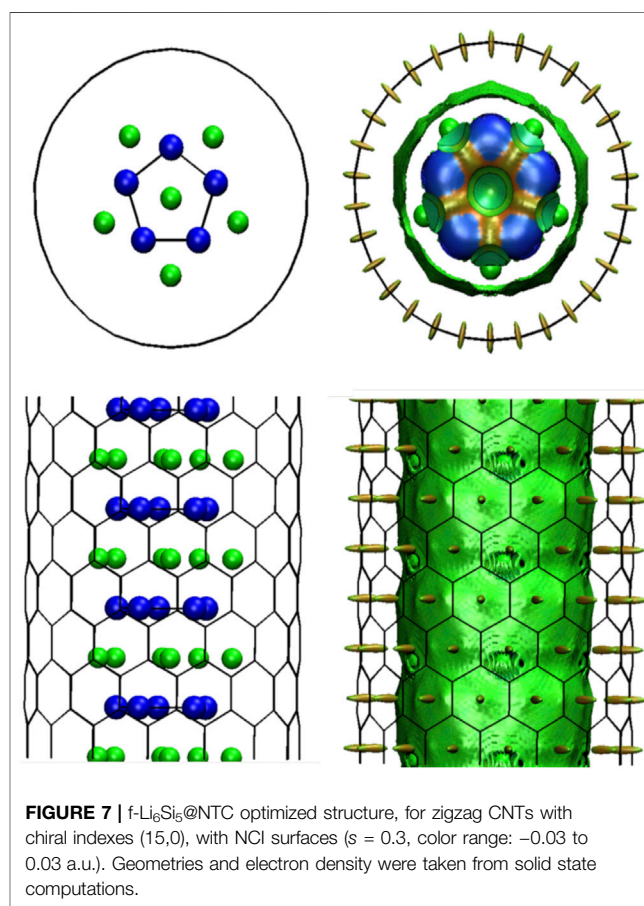
**FIGURE 5** | Energy as a function of time for molecular dynamic simulations at 500 K for the finite  $\text{Li}_6\text{Si}_5$  structure ( $f\text{-Li}_6\text{Si}_5$ ). The right images show top and side views of a snapshot taken at 10,000 fs



**FIGURE 6** | Interaction energy as a function of the nanotube diameters for the finite  $\text{Li}_6\text{Si}_5$  structures inside armchair CNTs with chiral indexes (8,8), (9,9), and (10,10) and zigzag CNTs with chiral indexes (14,0), (15,0), and (16,0).

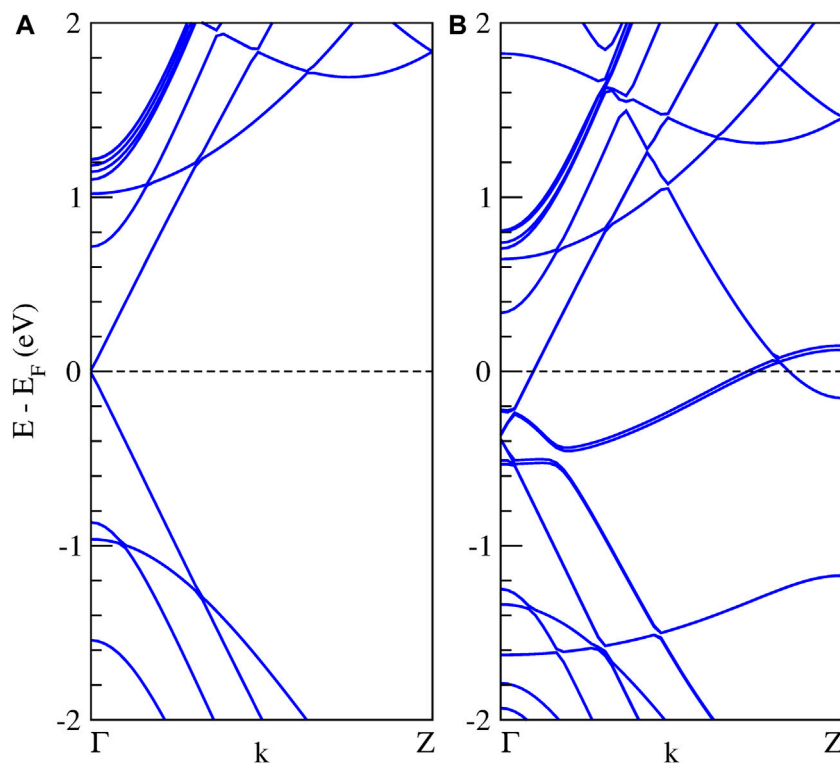
(Robles-Navarro et al., 2020). To get an idea of the nanotube widths to be considered in our study, we first performed a finite model analysis. This model consists of  $[n]$ cyclacenes ( $n = 13\text{--}20$ ) in their optimal structure (at the PBE0/Def2TZVP level), covering the diameter range from 10.2 to 15.6 Å. Then the  $\text{Li}_7\text{Si}_5^+$  cluster was placed, centering it on emulating the growth pattern towards the nanowire (see **Scheme 1**). We choose the star-shaped  $\text{D}_{5h}\text{-Li}_7\text{Si}_5^+$  cluster as a suitable model for projecting the nanowire inside the CNT due to its high symmetry and its analogy in electronic structure with the  $\text{Li}_6\text{Si}_5$  unit. In this study, we have kept the  $[n]$  cyclacene structure rigid, allowing only the optimization of the  $\text{Li}_7\text{Si}_5^+$  structure (at the PBE0/Def2TZVP level).

In the case of small  $[n]$ cyclacenes ( $n = 13\text{--}15$ ),  $\text{Li}_7\text{Si}_5^+$  cluster undergoes noticeable changes in the optimization process due to the



**FIGURE 7** |  $f\text{-Li}_6\text{Si}_5\text{@NTC}$  optimized structure, for zigzag CNTs with chiral indexes (15,0), with NCI surfaces ( $s = 0.3$ , color range:  $-0.03$  to  $0.03$  a.u.). Geometries and electron density were taken from solid state computations.

confinement effects. In contrast, when  $[n]$ cyclacenes with  $n = 16\text{--}20$  are used, the  $\text{Li}_7\text{Si}_5^+$  cluster maintains its structure at the end of the optimization process, leading to the best interaction energy,  $[E_{\text{int}} = E(\text{Li}_7\text{Si}_5^+@n[\text{cyclacene}]) - E(\text{Li}_7\text{Si}_5^+) - E(n[\text{cyclacene}])]$ , with [16] cyclacene ( $-70.1 \text{ kcal mol}^{-1}$  at PBE0/Def2TZVP level). Since this



**FIGURE 8** | Band structure calculations of the infinite  $\text{Li}_6\text{Si}_5$ -NW inside the zigzag carbon nanotubes, **(A)** the isolated (15,0) CNT, and **(B)** the  $\text{Li}_6\text{Si}_5$ -NW@CNT system. The dashed line indicates the Fermi energy.

analysis is only a reference for estimating the most suitable nanotube diameters to explore in the periodic calculations, we have not included basis set superposition error (BSSE) corrections. The most stable structures, as well as the  $E_{\text{int}}$ , are shown in **Figure 1**. The structures for the other complexes are shown in **Supplementary Figure S1** and their Cartesian coordinates in **Supplementary Table S1**. These results guided us to use CNTs with diameters in the range of 11–14 Å in next steps of our research.

### Insights on the stability of free $1/\infty [\text{Li}_6\text{Si}_5]$ NW.

We first studied the stability of periodically repeated  $\text{Li}_6\text{Si}_5$  units ( $\text{Li}_6\text{Si}_5$ -NW), which are stacked along the  $z$  direction, forming a one-dimensional structure, as shown in **Figure 2**.

We found a stable structure with an equilibrium distance between  $\text{Si}_5$  rings of 4.04 Å which is the periodicity of the  $\text{Li}_6\text{Si}_5$  unit cell. In the equilibrium geometry, the distance between Li atoms of the border is 3.33 Å while the distance with respect to the center one is 2.83 Å. The Si-Si distance between neighboring atoms is 2.37 Å, very close to the ones in  $\text{Li}_6\text{Si}_5$  monomer (between 2.30 and 2.35 Å). Our computations of the electronic band structure of the  $\text{Li}_6\text{Si}_5$ -NW in the primitive unit cell suggest a metallic character (see **Figure 3**). Note that the bandgaps of  $\text{Li}_{12}\text{Si}_7$  was reported from conductivity-temperature experimental measurements and found to be 0.6 eV (Nesper et al., 1986b).

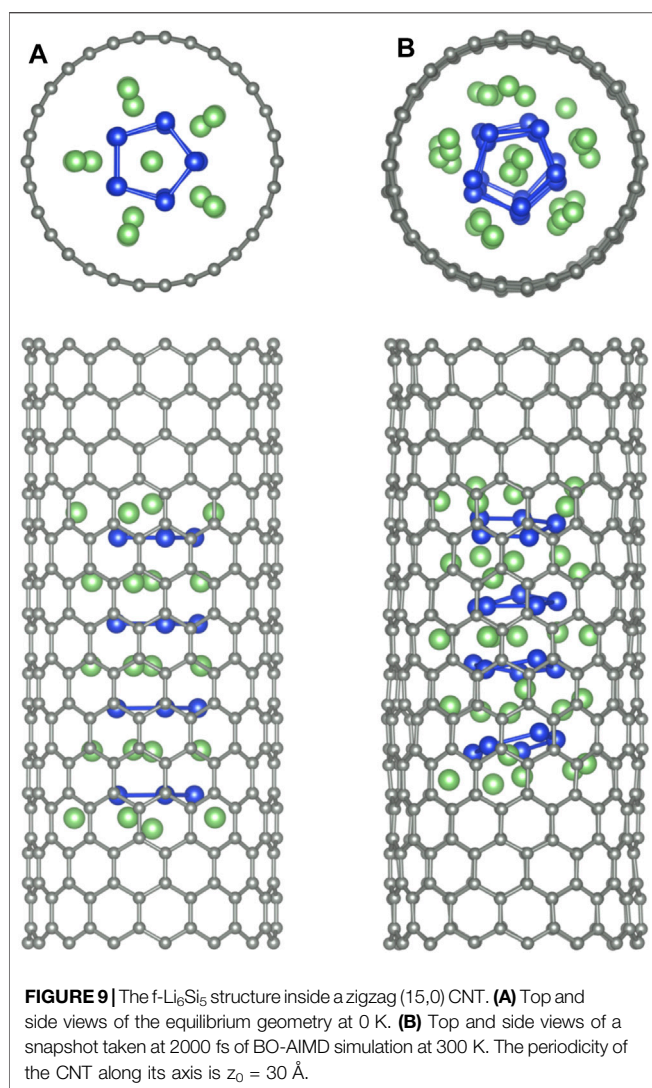
The stability of the  $\text{Li}_6\text{Si}_5$ -NW was also verified by BO-AIMD simulations at 300 K and 500K, during a simulation time of 10 ps.

The simulation was performed by considering four  $\text{Li}_6\text{Si}_5$  units in the periodic unit cell, as shown in **Figure 4**. We observe that at 500 K the  $\text{Si}_5\text{Li}_6$  NW preserves its stability, showing energy fluctuations of around 2 eV. Supporting information contains short movies extracted from the BO-AIMD simulations.

### Stability of the $\text{Li}_6\text{Si}_5$ -NW Inside the CNTs

Next, we studied a finite  $\text{Li}_6\text{Si}_5$  structure ( $f\text{-Li}_6\text{Si}_5$ ) in the free space and encapsulated it inside both armchair and zigzag carbon nanotubes ( $f\text{-Li}_6\text{Si}_5\text{@CNT}$ ). For the  $f\text{-Li}_6\text{Si}_5$  structure, we consider four  $\text{Si}_5$  rings surrounded by five  $\text{Li}_6$  moieties. BO-AIMD simulations provide insights on the stability of the  $f\text{-Li}_6\text{Si}_5$  system in the free space at 300 K and 500 K. We find that at 300 K, the  $f\text{-Li}_6\text{Si}_5$  structure preserves its stability. Still, at 500 K, it tends to form Si-Si bonds between adjacent Si rings without losing its one-dimensional array, as shown in **Figure 5**. However, it is important to note that this model does not have the exact stoichiometry of NW because, to maintain symmetry, an extra  $\text{Li}_6$  unit group is added, i.e.,  $[(\text{Li}_6)_5(\text{Si}_5)_4]$ .

To study the  $f\text{-Li}_6\text{Si}_5$  structure inside the CNTs, we consider three armchair CNTs with chiral indexes of (8,8), (9,9), and (10,10), and three zigzag CNTs with chiral indexes of (14,0), (15,0), and (16,0). With this choice, we seek to find the CNT size that best accommodates the  $\text{Li}_6\text{Si}_5$ -NW inside. Note that we selected these CNTs according to our preliminary findings from the  $\text{Li}_7\text{Si}_5^+\text{@}[n]\text{cyclacene}$  model, suggesting diameters between 12 and 15 Å. The CNTs were simulated with periodic boundary conditions along its axis with a periodicity of  $z_0 = 30$  Å.



The latter allows a vacuum region for the encapsulated  $f\text{-Li}_6\text{Si}_5$  structure of  $14 \text{ \AA}$ , allowing the atomic movement inside the CNT. Next, we calculate the  $E_{\text{int}}$  between encapsulated  $f\text{-Li}_6\text{Si}_5$  and the CNTs by the equation:

$$E_{\text{int}} = E_{\text{tot}}(f\text{-Li}_6\text{Si}_5@CNT) - E_{\text{tot}}(f\text{-Li}_6\text{Si}_5) - E_{\text{tot}}(CNT),$$

where  $E_{\text{tot}}(f\text{-Li}_6\text{Si}_5)$  and  $E_{\text{tot}}(CNT)$  are the total energies of the isolated  $f\text{-Li}_6\text{Si}_5$  and the CNT, respectively, and  $E_{\text{tot}}(f\text{-Li}_6\text{Si}_5@CNT)$  is the total energy of  $f\text{-Li}_6\text{Si}_5$  unit inside the CNT. Our results for the  $E_{\text{int}}$  are shown in **Figure 6**. The armchair (9,9) and the zigzag (15,0) CNTs of  $12.3$  and  $11.8 \text{ \AA}$  in diameter, respectively, exhibit the more stabilizing  $E_{\text{int}}$ , being the best candidates to accommodate the  $f\text{-Li}_6\text{Si}_5$  inside. In addition, the  $f\text{-Li}_6\text{Si}_5$  is better stabilized inside the zigzag (15,0) CNT than inside the armchair (9,9) CNT by  $0.9 \text{ eV}$ . Noteworthy, the larger-diameter CNTs are energetically less favorable to encapsulate the  $f\text{-Li}_6\text{Si}_5$ , as shown in **Figure 6**, in agreement with the  $\text{Li}_7\text{Si}_5^+@[n]$  cyclacene model. This behavior is presumable due to the van der Waals (vdW) interaction between the  $f\text{-Li}_6\text{Si}_5$  and the CNT internal walls, stabilizing the system. The non-covalent

interaction index (NCI) plots confirm the non-covalent character of  $f\text{-Li}_6\text{Si}_5$  with the CNT [ $f\text{-Li}_6\text{Si}_5$  inside the zigzag (15,0) CNT]. In the NCI method (Johnson et al., 2010; Contreras-García et al., 2011), an isosurface of the reduced density gradient ( $s$ ) is colored with the density times the sign of the second eigenvalue of the electron density Hessian matrix,  $\lambda_2$ , to distinguish between attractive and repulsive interactions. The following color code is used: blue for attractive such as hydrogen bonds, green for very weak interactions such as vdW and red for steric repulsion. **Figure 7** depicts the second one (vdW) between  $f\text{-Li}_6\text{Si}_5$  and the walls of the CNT.

We also calculate the band structure of the  $\text{Li}_6\text{Si}_5\text{-NW}$  inside the zigzag (15,0) CNT. It is important to note that the unit cells of the  $\text{Li}_6\text{Si}_5\text{-NW}$  and the CNT have a mismatch of  $5.7\%$ , which means that the  $\text{Li}_6\text{Si}_5\text{-NW}$  is not in its equilibrium geometry in the  $\text{Li}_6\text{Si}_5\text{-NW}@CNT$  unit cell, where the distance between the  $\text{Si}_5$  rings increases by  $0.23 \text{ \AA}$ . However, this mismatch is relatively small and should not affect the electronic properties of the system. For the isolated CNT we find a small bandgap of  $0.02 \text{ eV}$  as shown in **Figure 8A**, which is in good agreement with the measured value of  $0.029 \pm 0.004 \text{ eV}$  (Ouyang et al., 2001). Whereas the  $\text{Li}_6\text{Si}_5\text{-NW}@CNT$  system exhibits metallic properties as shown **Figure 8B**, suggesting that the  $\text{Li}_6\text{Si}_5\text{-NW}$  would preserve its electronic properties inside the CNT as can be compared with **Figure 3**.

Finally, we study the stability of the  $f\text{-Li}_6\text{Si}_5$  structure inside both zigzag (15,0) and armchair (9,9) CNTs by performing BO-AIMD simulations. **Figure 9A** shows the equilibrium geometry of the  $f\text{-Si}_5\text{Li}_6$  structure inside the (15,0) CNT. We find that the structure remains almost unchanged with respect to  $f\text{-Si}_5\text{Li}_6$  in the free space, showing that the CNT would have a small influence in the  $\text{Li}_6\text{Si}_5$  NW stability. We only note a small displacement of the Li ions at the extreme of the  $f\text{-Li}_6\text{Si}_5$  structure which move toward the CNT wall. The integrity of the  $f\text{-Li}_6\text{Si}_5$  structure inside the (15,0) and (9,9) CNTs was investigated by BO-AIMD simulations at  $300 \text{ K}$ . We find that the  $f\text{-Si}_5\text{Li}_6$  structure preserves its stability where the Li ions move around the  $\text{Si}_5$  ring without detaching. Similar results are found for the  $f\text{-Li}_6\text{Si}_5$  structure inside the armchair (9,9) CNT, indicating that the formation and stability of the  $\text{Li}_6\text{Si}_5$  NW inside the CNTs is independent of its chirality. This result suggests that  $\text{Li}_6\text{Si}_5\text{-NW}$  are likely to form inside CNTs in a compact form, which would allow efficient storage of Li ions into CNTs mediated by  $\text{Si}_5$  rings. **Supplementary Figure S2** shows the variation of the total energy for the BO-AIMD simulation of  $f\text{-Li}_6\text{Si}_5$  inside the (15,0) and (9,9) CNTs at  $300 \text{ K}$ . We observe energy fluctuation of around  $5 \text{ eV}$  in both CNTs, preserving the stability of the  $f\text{-Li}_6\text{Si}_5$  structure.

## CONCLUSION

Using periodic DFT calculations and Born-Oppenheimer *ab initio* molecular dynamic simulations, we have shown that  $\text{Li}_6\text{Si}_5$  units can be stacked one above the other, forming a one-dimensional structure linked together by Coulomb interactions. This study complements previous findings, where we



demonstrated that  $\text{Li}_6\text{Si}_5$ ,  $(\text{Li}_6\text{Si}_5)_2$ , and  $(\text{Li}_6\text{Si}_5)_3$  lowest energy structures contain one, two, and three  $\text{Si}_5^{6-}$  aromatic rings stabilized by  $\text{Li}^+$  counterions. Additionally, the  ${}^1_{\infty}[\text{Li}_6\text{Si}_5]$  nanowire was identified in the Zintl  $\text{Li}_{12}\text{Si}_7$  compound but coexisting with Y-shaped  $\text{Si}_4$  moieties. In this case, we support the stability of the isolated Si-Li-nanowire—additionally, the relaxed structure (at room temperature) exhibits metallic characteristics.

We also found that finite  $(\text{Li}_6\text{Si}_5)_4$  systems are stable inside both armchair and zigzag carbon nanotubes of around 12 Å in diameter, preserving its stability at room temperature, supporting the viable formation of  $\text{Li}_6\text{Si}_5$ -NW inside the CNTs. Interestingly, the  $\text{Li}_6\text{Si}_5$ -NW@CNTs hybrid nanocomposite maintains the metallic character. Finally, in the  $\text{Li}_6\text{Si}_5$ -NW, the  $\text{Li}_6\text{Si}_5$  units are connected by strong electrostatic interactions ( $\text{Si}_5^{6-}$  aromatic pentagons intercalated with the  $\text{Li}_6^{6+}$  moiety) in agreement with the Zintl ion concept. In the  $[\text{Li}_6\text{Si}_5\text{-NW}]@$ CNTs, NCI predicts that  $\text{Li}_6\text{Si}_5$ -NW interacts with the CNT walls by van der Waals interactions Frisch et al., 2016.

## DATA AVAILABILITY STATEMENT

The original contributions presented in the study are included in the article/**Supplementary Material**, further inquiries can be directed to the corresponding author.

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## AUTHOR CONTRIBUTIONS

All authors listed have made a substantial, direct, and intellectual contribution to the work and approved it for publication.

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## SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: <https://www.frontiersin.org/articles/10.3389/fchem.2021.767421/full#supplementary-material>

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